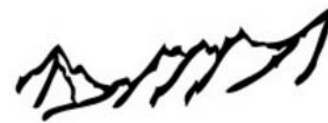


Space plasmas measurement techniques

Milan Maksimovic

CNRS & LESIA, Paris Observatory, France

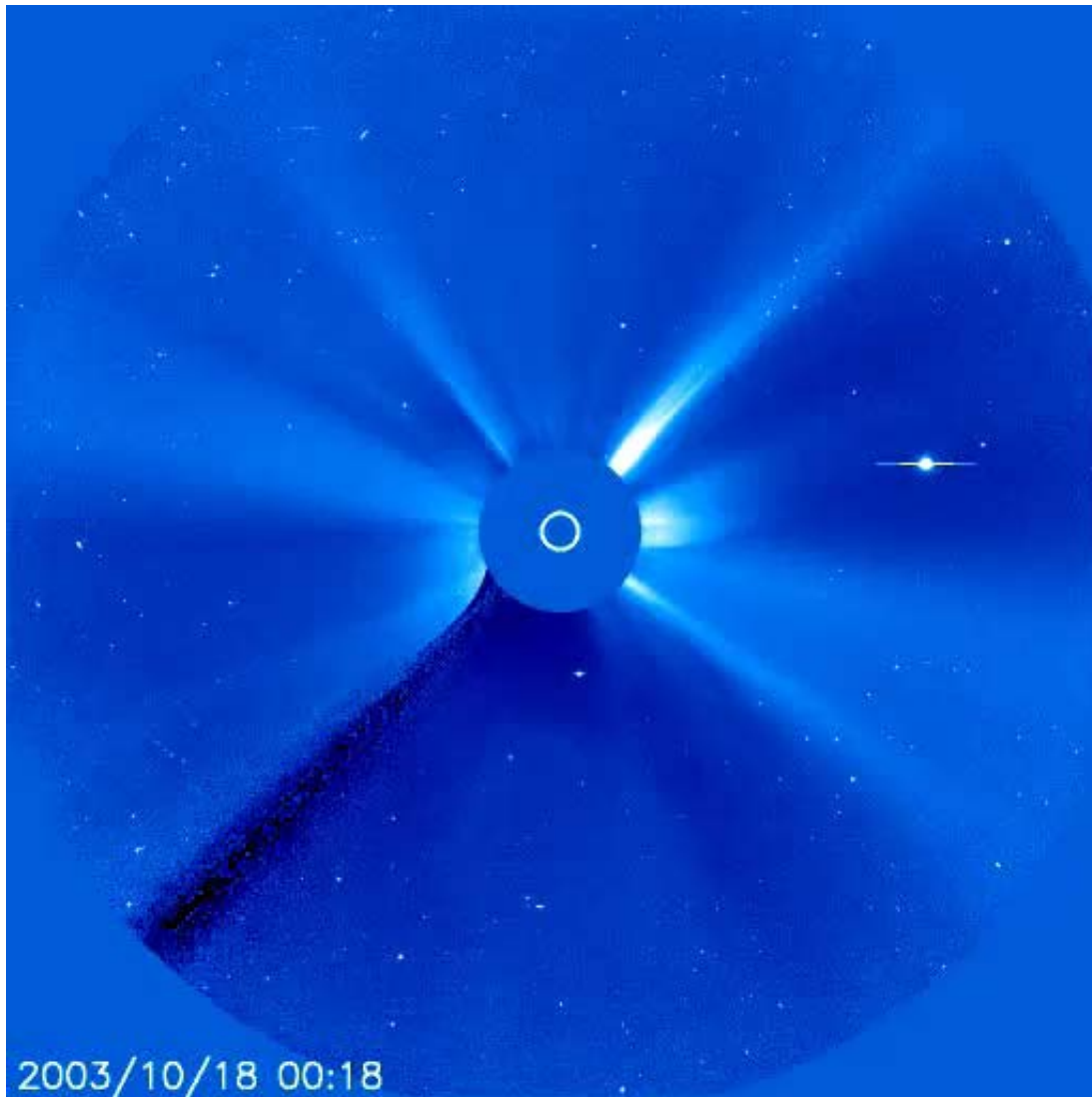
ÉCOLE DE PHYSIQUE
des HOUCHES



UNIVERSITÉ
Grenoble
Alpes

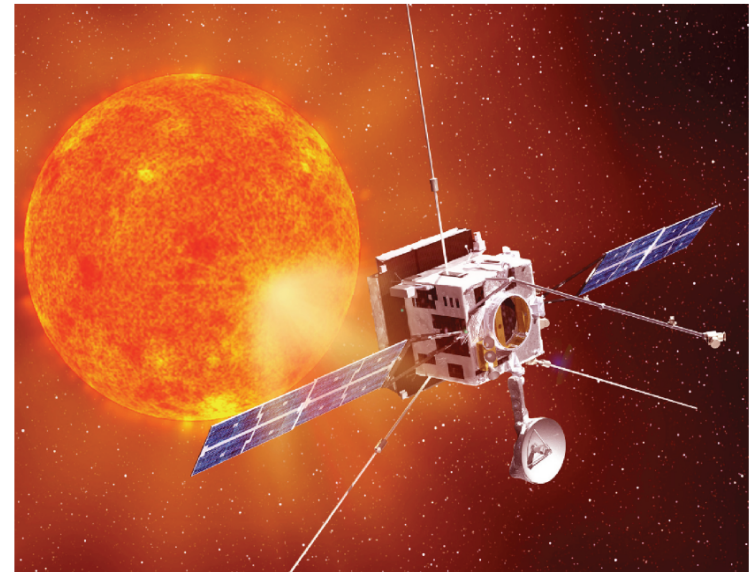


FROM LABORATORIES TO ASTROPHYSICS: THE EXPANDING UNIVERSE OF PLASMA PHYSICS



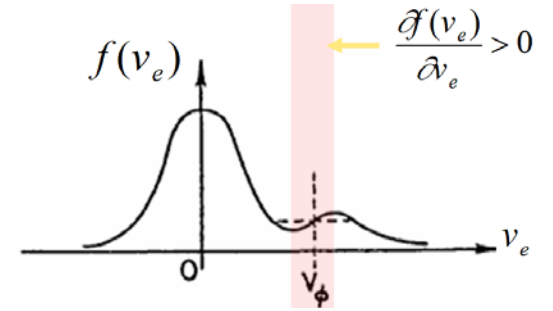
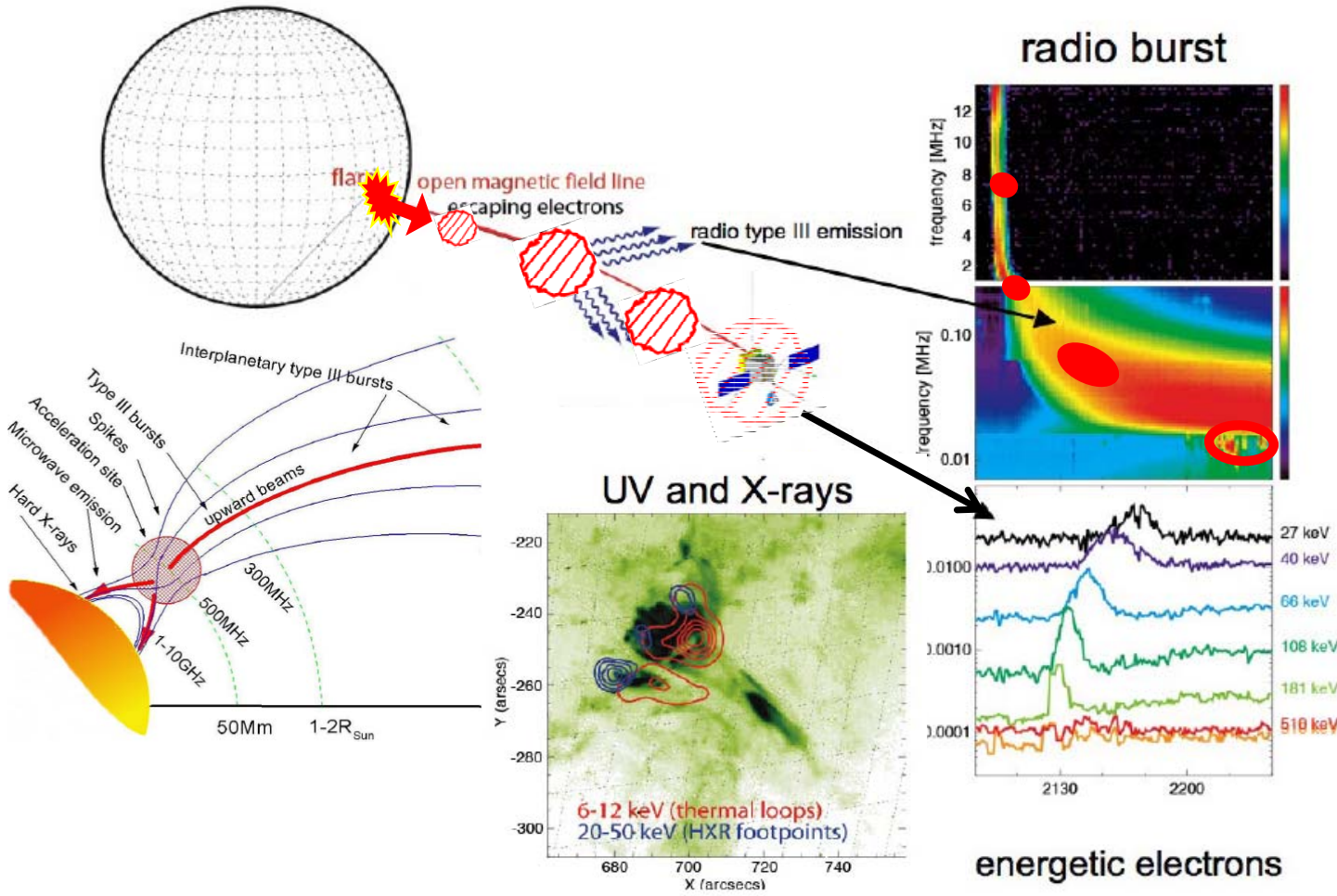
Solar Orbiter

Exploring the Sun-heliosphere connection



- **ESA mission with NASA participation (launcher + two instruments)**
- **Up to 0.28 AU with dedicated in-situ & remote sensing instrumentation (closest approach possible)**
- **Out of ecliptic observations**
- **1st M-class ESA Cosmic Vision (Launch ~Feb. 2019)**

Interplanetary high energy electron & Solar Radio Bursts

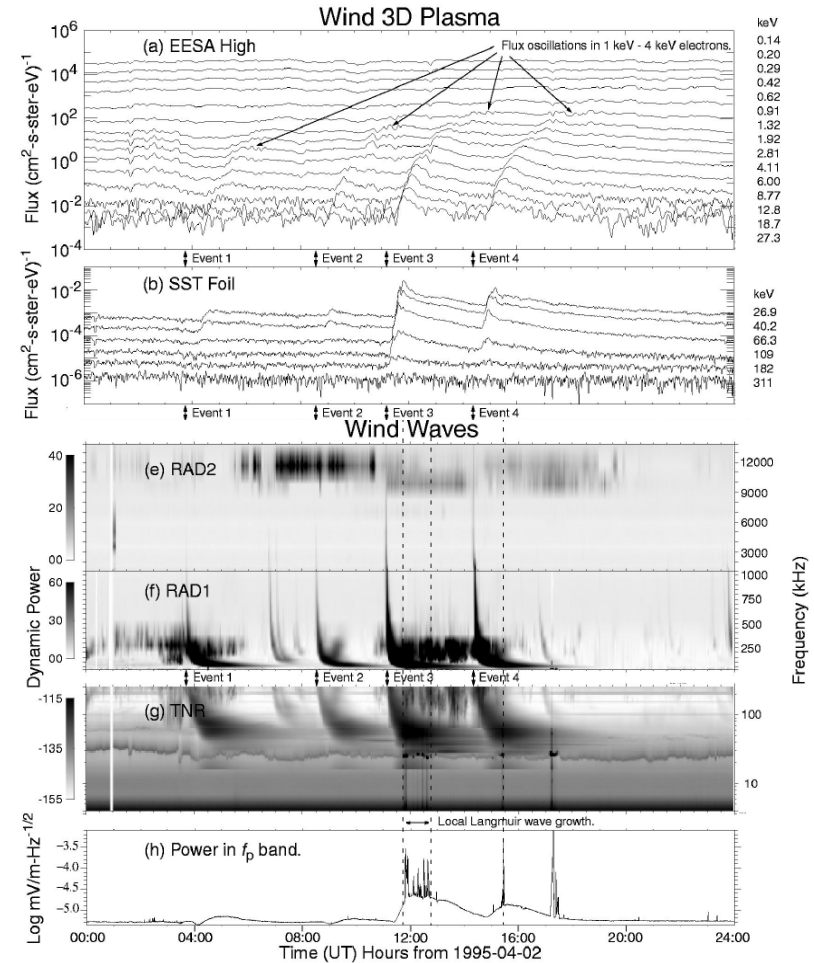
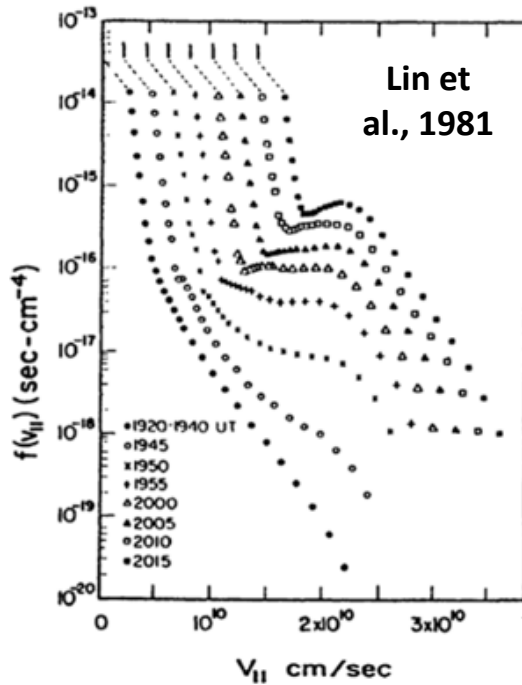
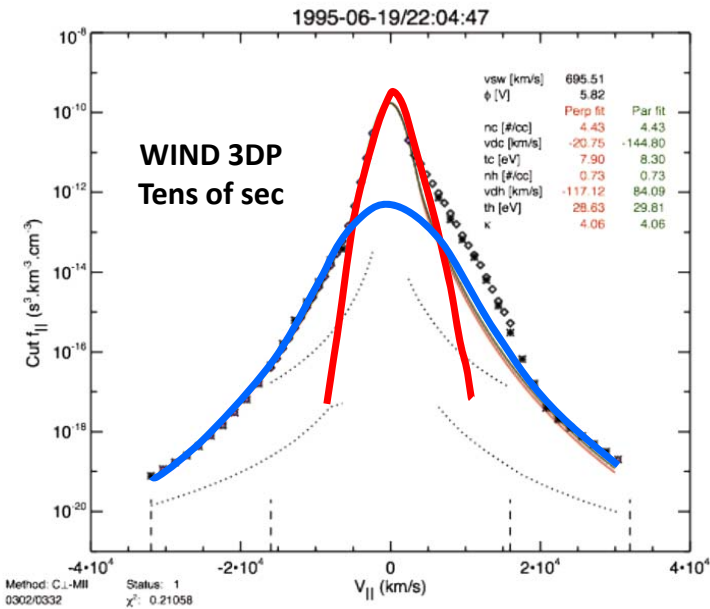


**Electrostatic
Langmuir waves
→ radio emission**

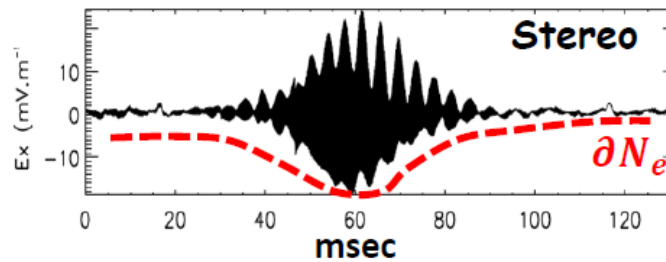
$$\left. \begin{aligned} F_p (\text{kHz}) &\propto \sqrt{N_e (\text{cm}^{-3})} \\ N_e &\propto 1/R^2 (\text{au}) \end{aligned} \right\}$$

$$\rightarrow F_p \propto \frac{1}{R}$$

All necessary measurements will be available on Solar Orbiter



- + plasmas measurements
 - Np, Tp, Vp (16Hz)
 - High cadence Ne
 - Bx, By, Bz DC & AC (up to the plasma frequency)



Adapted from [Ergun et al., 1998]

Outline

☐ **Particles measurements**

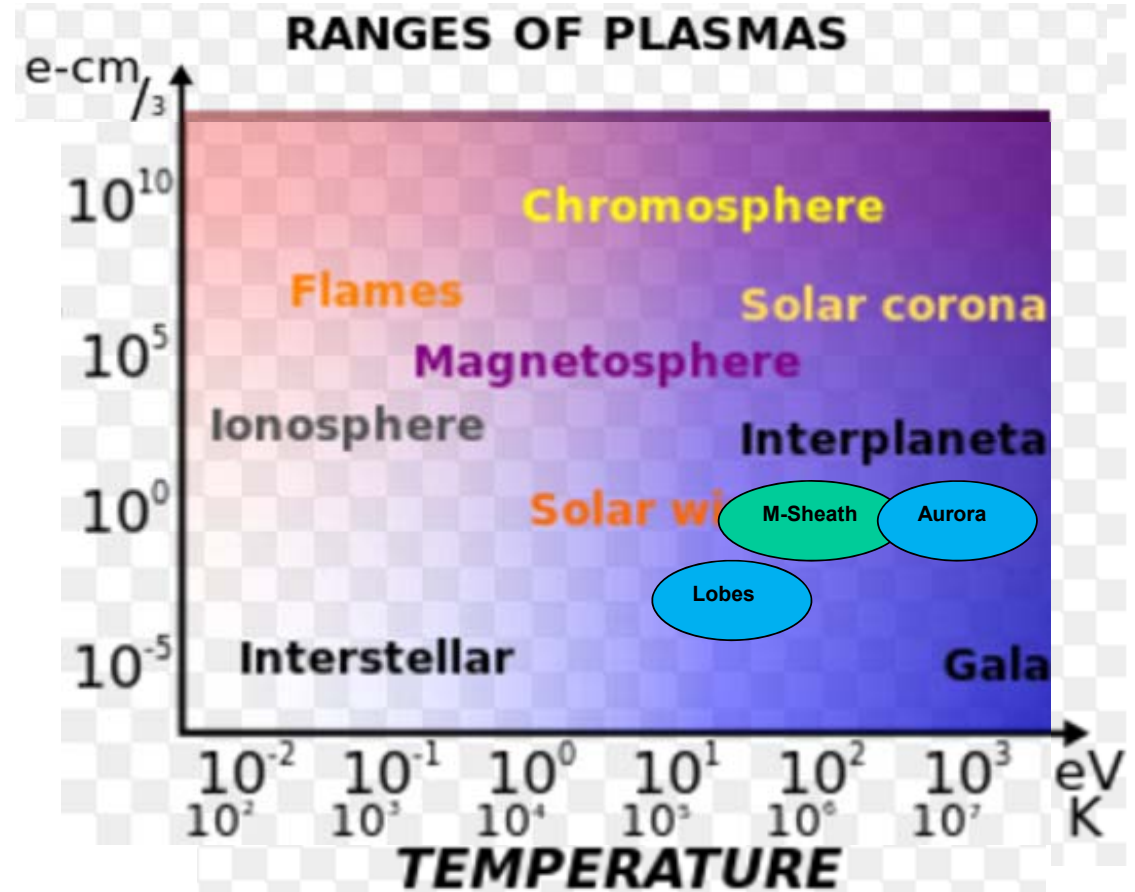
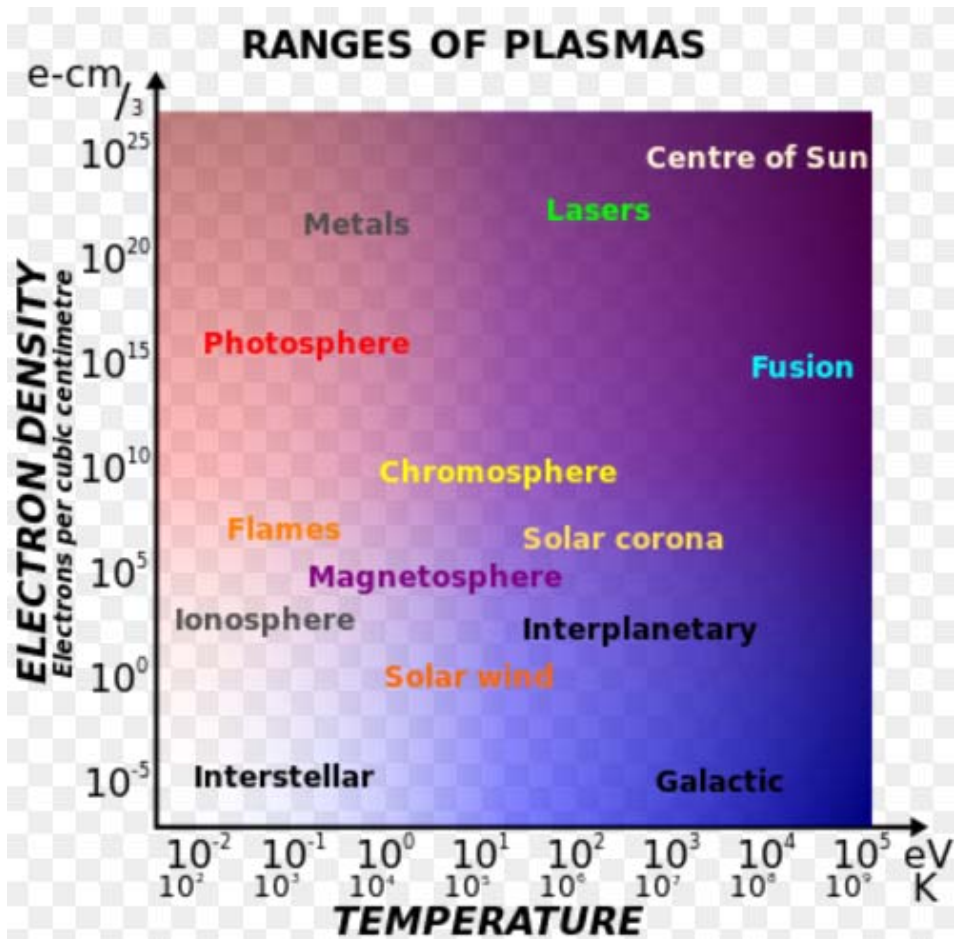
- **Electrons**
- **Ions & composition**
- **Energetic particles**

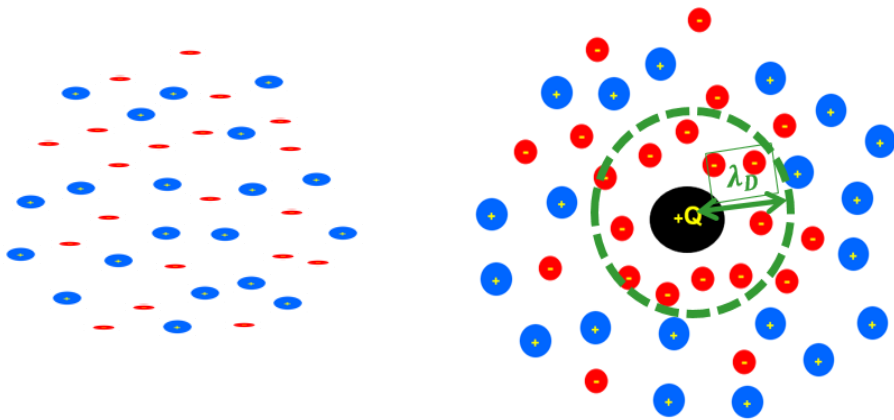
☐ **Waves measurements**

- **Magnetic DC & AC sensors**
- **Measuring DC-LF electric fields**
- **Measuring AC Electric Field: The Thermal Noise as an example**

☐ **Surprises, tricky data corrections and future challenges**

Astrophysical Plasmas accessible to in-situ measurements





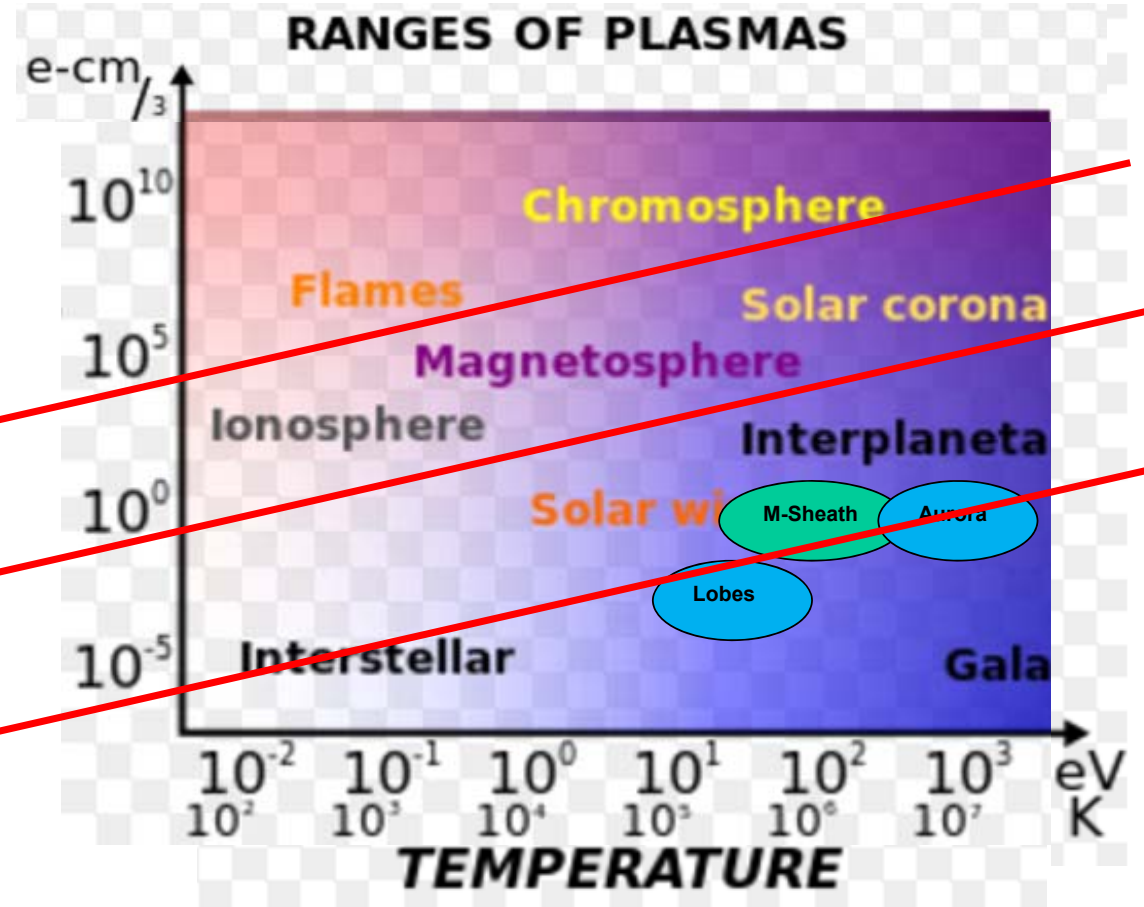
$$\Phi(r) = \frac{Q}{4\pi\epsilon r} e^{-r/\lambda_D}$$

$$\lambda_D = 2 \cdot 10^{-3} \text{ m}$$

$$\lambda_D = 0.7 \text{ m}$$

$$\lambda_D = \sqrt{\frac{\epsilon_0 k T}{n e^2}} \sim 69 \sqrt{\frac{T [\text{K}]}{n [\text{m}^{-3}]}}$$

$$\lambda_D = 218 \text{ m}$$

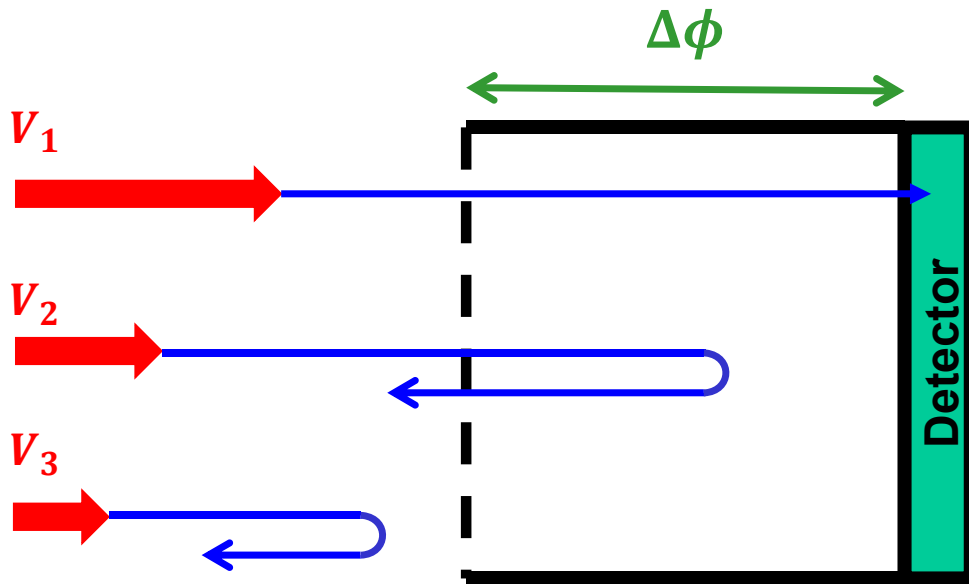


Particle Energy Distributions

A plasma particle is defined by

- a mass
- an electric charge
- a speed or kinetic energy

$$\frac{1}{2}mV(r)^2 + Ze\phi(r) = \frac{1}{2}mV_{\infty}^2$$



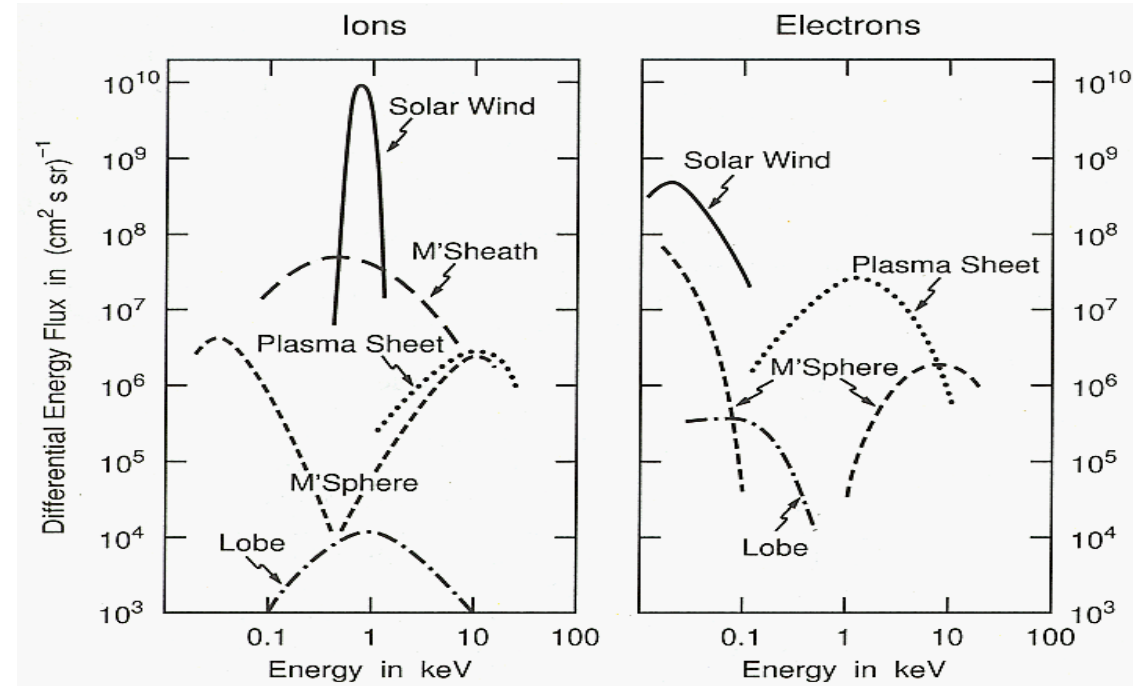
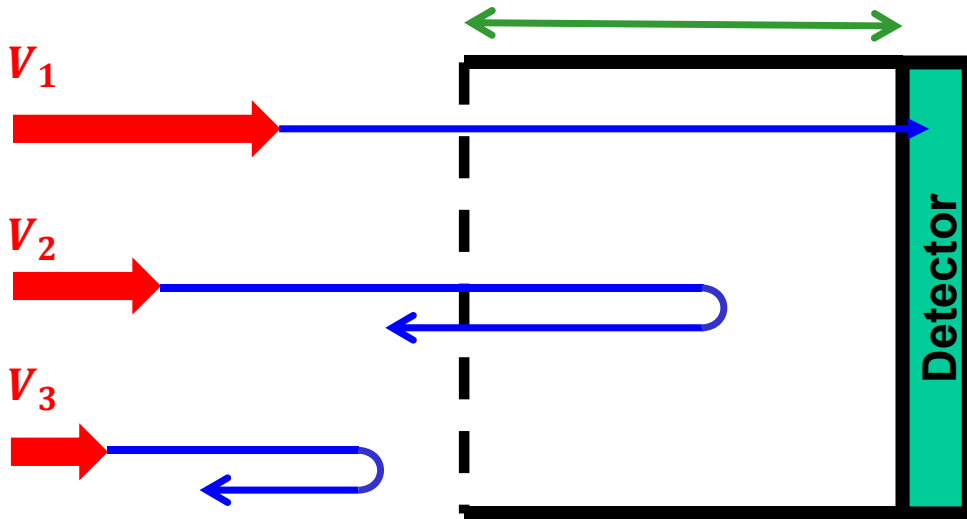
Particle Energy Distributions

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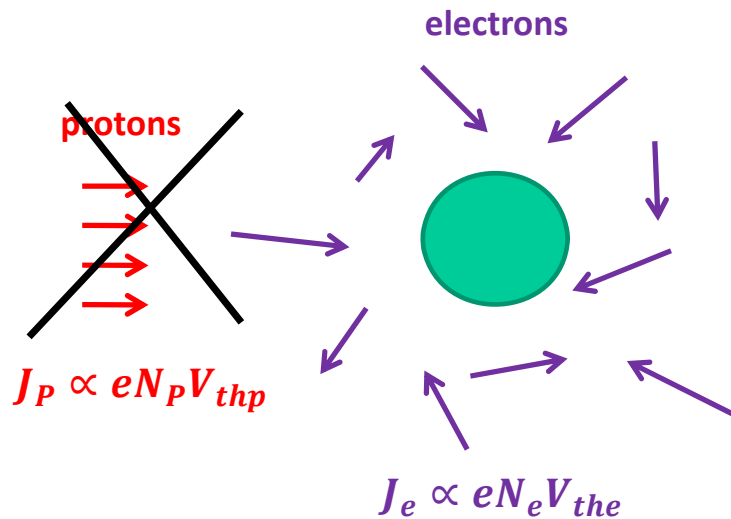
$$\frac{1}{2}mV(r)^2 + Ze\phi(r) = \frac{1}{2}mV_{\infty}^2$$

$\Delta\phi$



The broad range of energies and fluxes
require different instruments
← Faraday cup principle

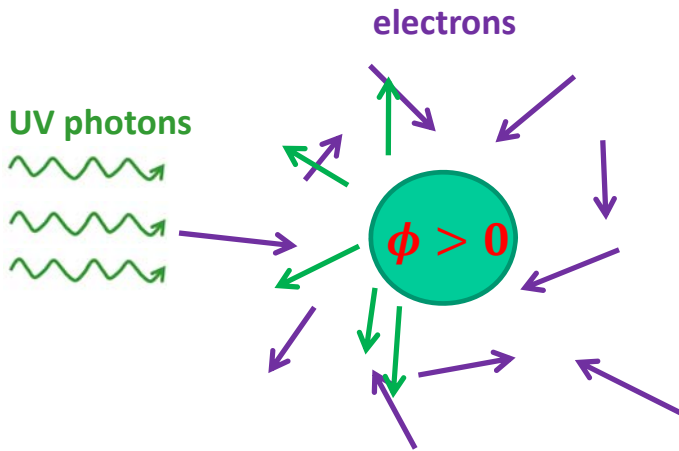
Electric charging in space : an example in the Solar Wind



For a plasma in quasi-equilibrium with $T_e \sim T_i$

- $V_{the} = \sqrt{\frac{2kT_e}{m_e}} \gg V_{thi} = \sqrt{\frac{2kT_i}{m_i}}$
- In the Solar Wind $T_e \sim T_p \sim 10^5 K$
 $\rightarrow V_{the} \sim 1740 \text{ km/s} \gg V_{thp} \sim 40 \text{ km/s}$

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$$V_e^* = \frac{4\pi}{N_e^*} \int_{v_\phi}^{\infty} v' f'(v') v'^2 dv'$$

$$J_e = \frac{2e}{\sqrt{\pi}} N_e V_{the} \left(1 + \frac{e\phi}{kT_e} \right)$$

Finally the potential is obtained by assuming

$$J_e(\phi) = J_{ph}(\phi)$$

$$\frac{e\phi}{k_B T_{ph}} - \ln \left[\frac{1 + \frac{e\phi}{k_B T_{ph}} \frac{K J_{ph0}}{N_e \sqrt{T_e}}}{1 + \frac{e\phi}{k_B T_e}} \right] = 0$$

What are the average electron velocity V_e^* and density N_e^* at the surface of the Sphere?

$J_e = e N_e^* V_e^*$ can be obtained applying
 Louville Theorem

$$\frac{1}{2} m_e v'^2 + Ze\phi = \frac{1}{2} m_e v'^2$$

$$v'_{\min} = \sqrt{2e\phi/m_e} \sim V_{the}/\sqrt{2}$$

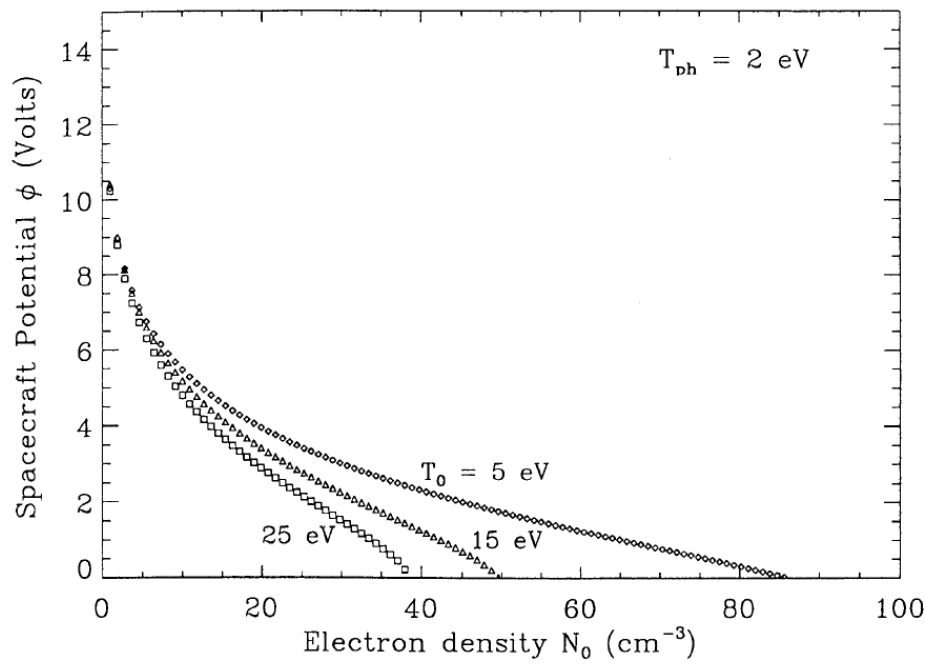
And the photoelectron current ?

$$J_{ph} = J_{ph0} \left(1 + \frac{e\phi}{k_B T_{ph}} \right) e^{-e\phi/k_B T_{ph}}$$

- Depends on the photo-electric properties of the material
- $T_{ph} \sim 2 \text{ to } 3 \text{ eV}$ and $J_{ph} \sim \text{few } \mu\text{A/m}^2$

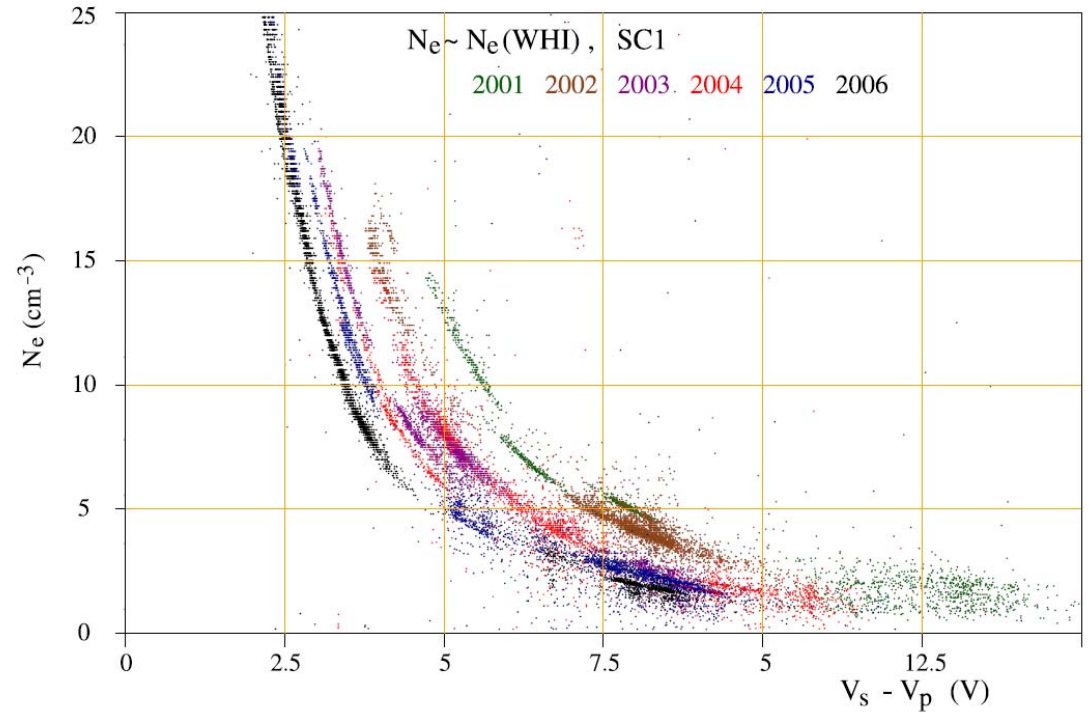
To be fully exhaustive one shall include the secondary electrons currents (important closer to the Sun)

$$\frac{e\phi}{k_B T_{ph}} - \ln \left[\frac{1 + \frac{e\phi}{k_B T_{ph}} \frac{KJ_{ph0}}{N_e \sqrt{T_e}}}{1 + \frac{e\phi}{k_B T_e}} \right] = 0$$



Salem et al., 2001

Cluster data



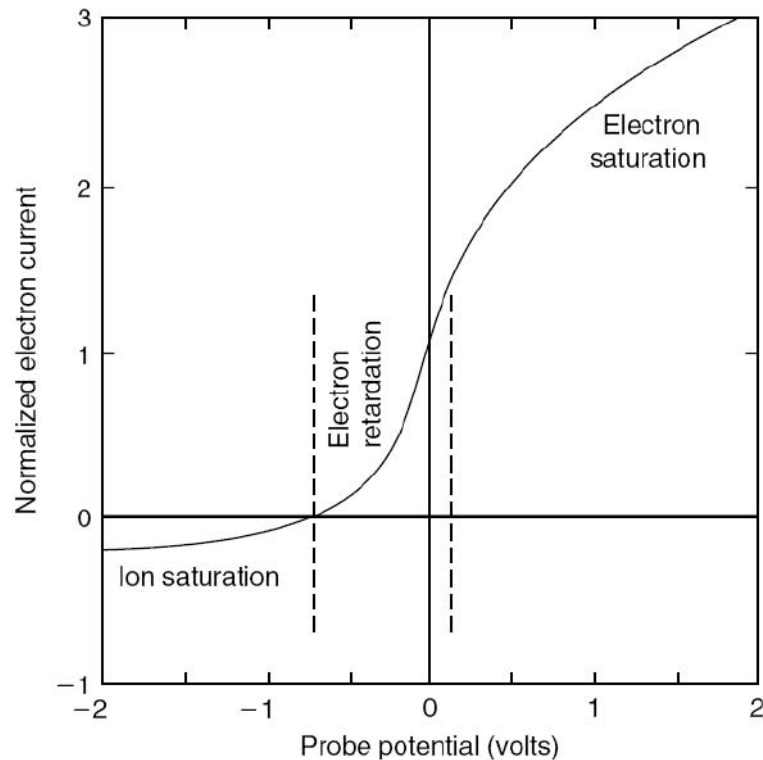
Pedersen et al., 2008

Electron diagnostic in cold plasmas – The Langmuir Probe



Sweeping V and measuring I

Typical Langmuir probe I-V characteristic



- For a plasma in quasi-equilibrium with $T_e \sim T_i$

$$V_{te} = \sqrt{\frac{2kT_e}{m_e}} \gg V_{ti} = \sqrt{\frac{2kT_i}{m_e}}$$

- in the retardation region

$$I_e \propto N_e e V_{the} e^{-eV/k_B T_e}$$

- In the saturation region

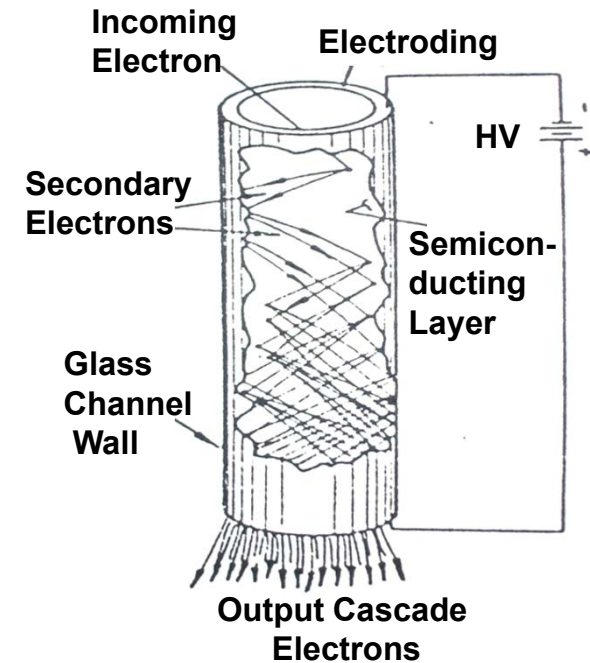
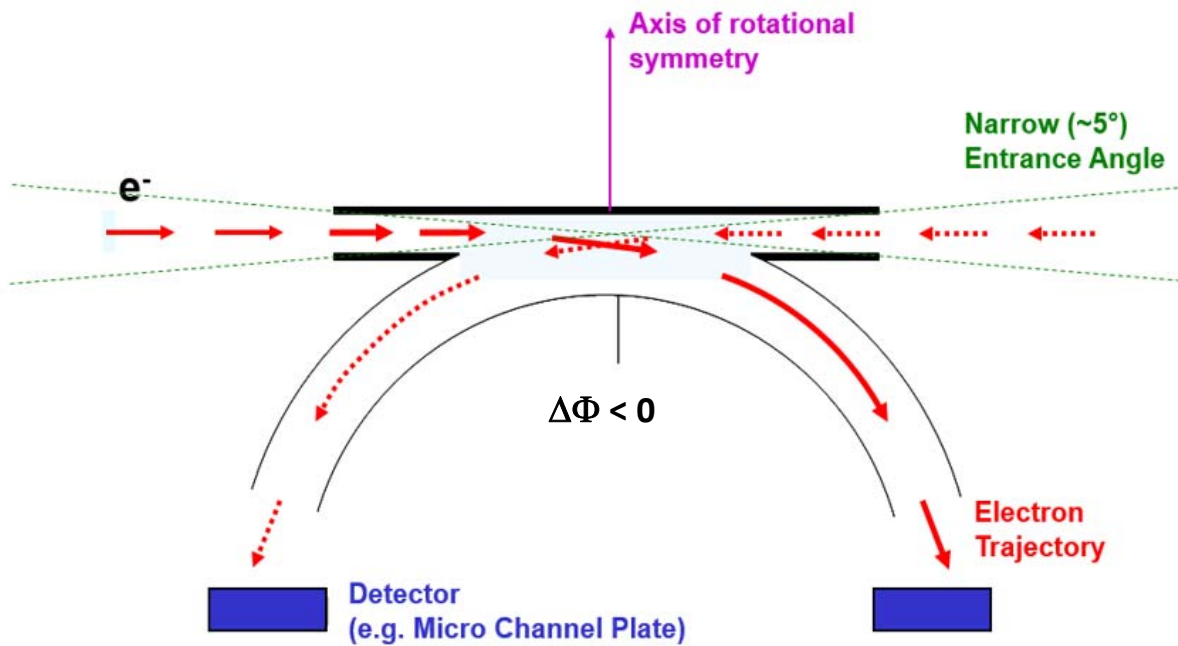
$$I_e \propto N_e e V_{the} \left(1 + \frac{eV}{k_B T_e} \right)^{1/2}$$

Segmented langmuir probe in order to measure the ion bulk speed



Measuring the full 3D VDF

'Top Hat' detectors – Principles of Operation



Hollow cylinders

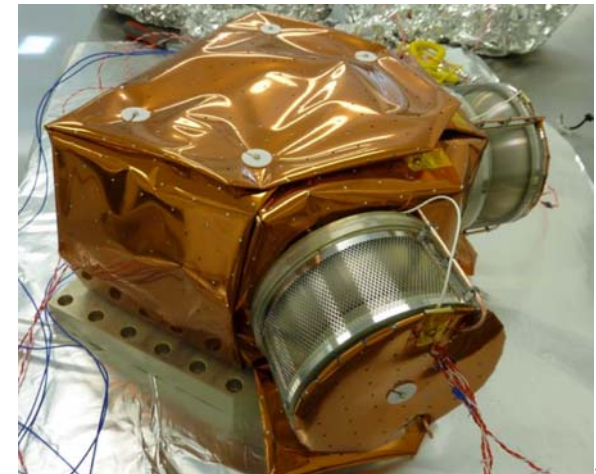
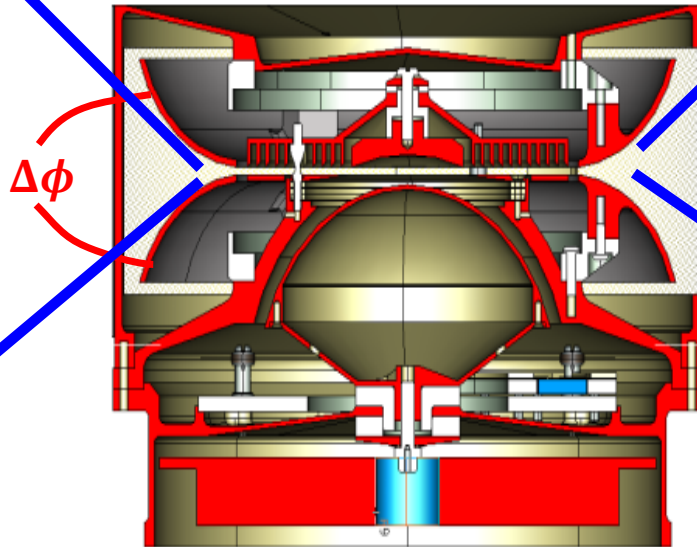
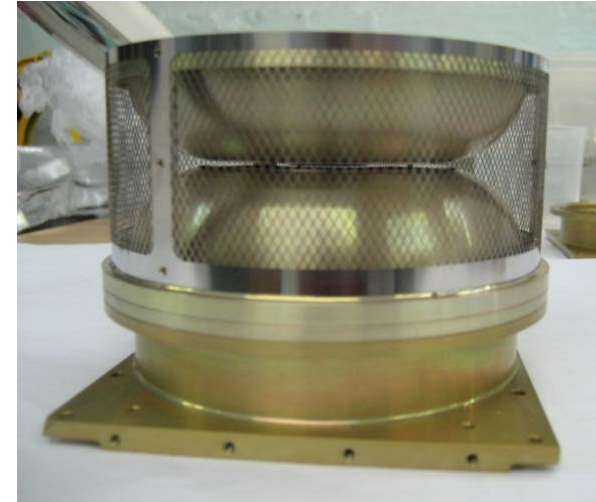
A full 3D Velocity Distribution function is obtained by

- Scanning the energy ($\Delta\Phi$)
- Using the Top hat axis of symmetry and the spacecraft spin



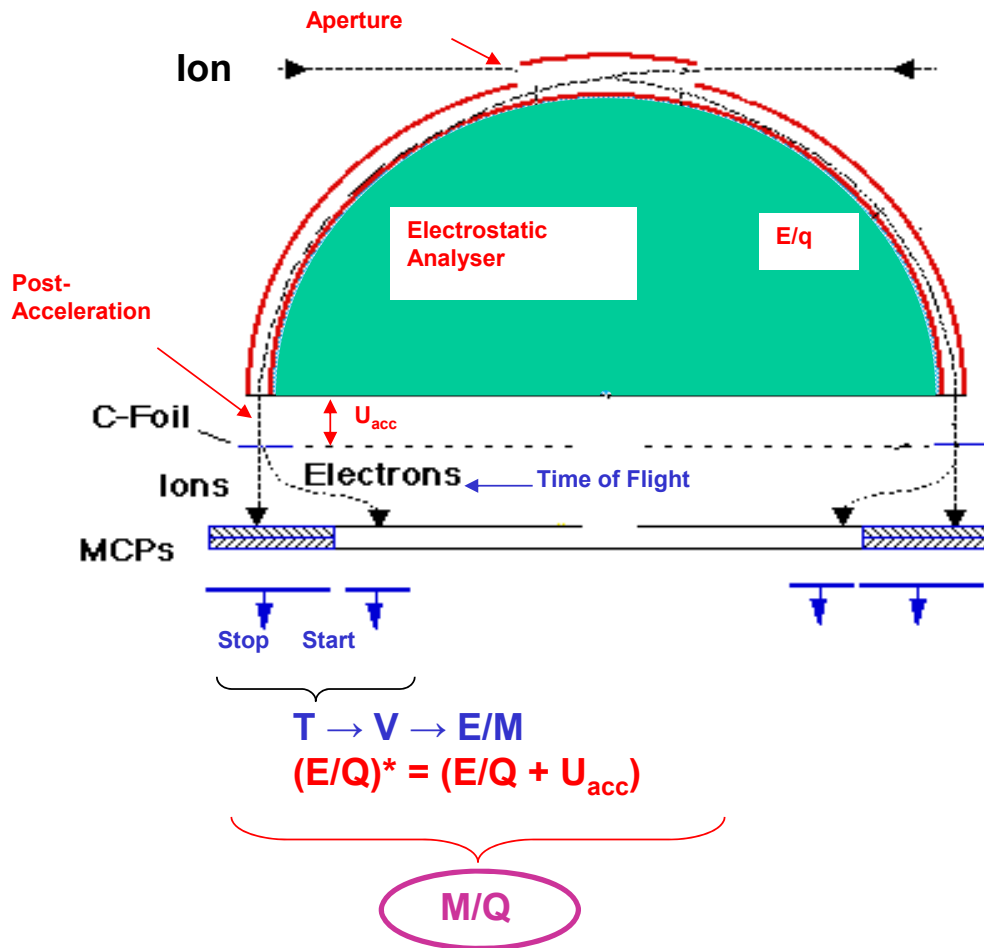
The Solar Orbiter EAS instrument

*Christopher « Chris » Owen,
Mullard Space Sciences Lab., UK*



Particle Composition :

Add a time-of-flight detector beneath the top hat analyser:

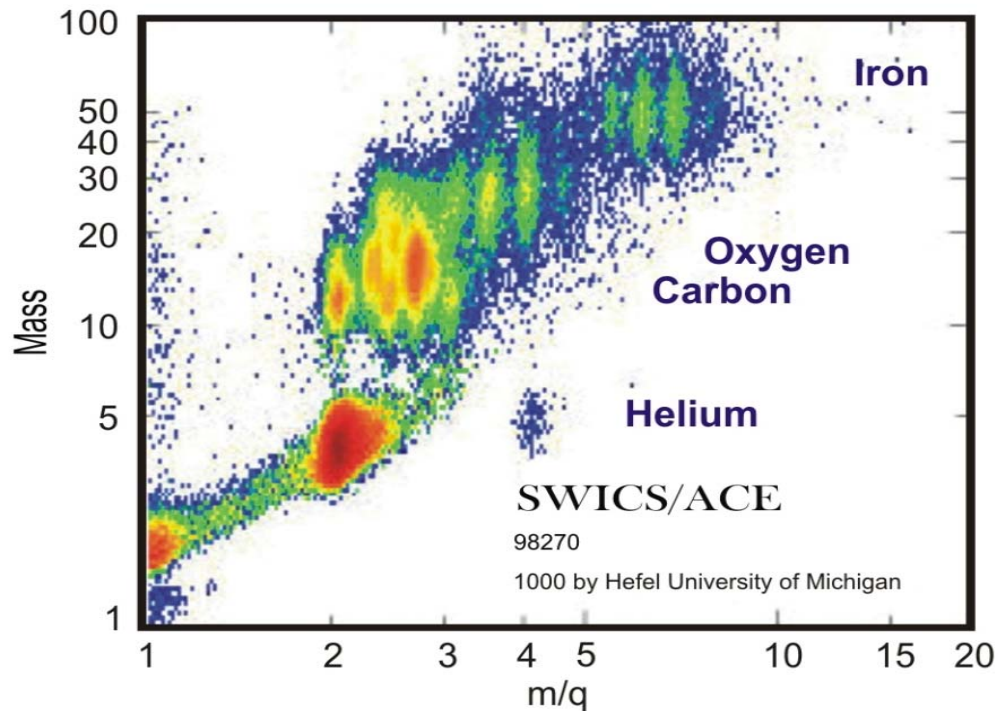


- Top hat section makes E/q selection as before : $E/Z e \propto \Delta\Phi$
- Ions are then accelerated by an electric field into a thin carbon foil
- On passing through the foil the ion knocks out an electron
- The difference in travel time to the detector between the ion and electron can be used to determine the ions velocity, and hence E/M for the ion;
- Combining the two measurements gives M/Q

Particle Composition :

Add a time-of-flight detector beneath the top hat analyser:

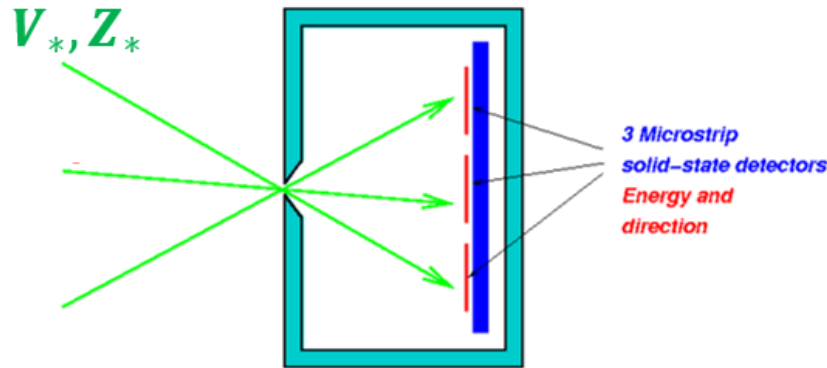
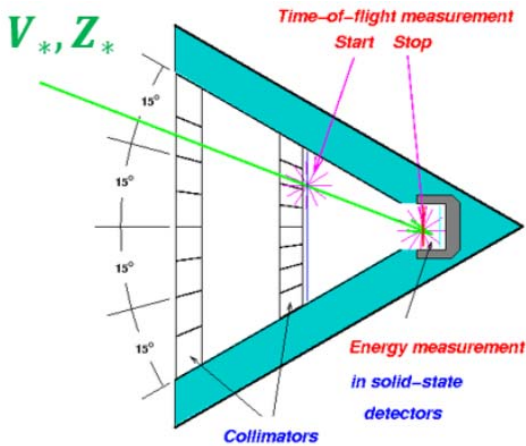
Composition Plot – colours represent relative abundance of ions



- Top hat section makes E/q selection as before : $E/Ze \propto \Delta\Phi$
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High-Energy Particle Telescopes

Robert "Bob" Wimmer-Schweingruber
University of Kiel, Germany



Energy loss Bethe-Block equation

Javier Rodríguez-Pacheco
University of Alcala, Spain

$$\frac{dE}{dx} = - \frac{Z_*^2 e^4 n_e}{4\pi\epsilon_0^2 V_*^2 m_e} \left[\ln \left(\frac{2eV_*^2}{\langle E_B \rangle} \right) - \ln(1 - \beta^2) - \beta^2 \right] \text{ with } \beta = \frac{V_*}{c}$$

$\langle E_B \rangle$ is the « first energy ionization » of the target

Outline

☐ Particles measurements

- Electrons
- Ions & composition
- Energetic particles

☐ Waves measurements

- Magnetic DC & AC sensors
- Measuring DC-LF electric fields
- Measuring AC Electric Field: The Thermal Noise as an example

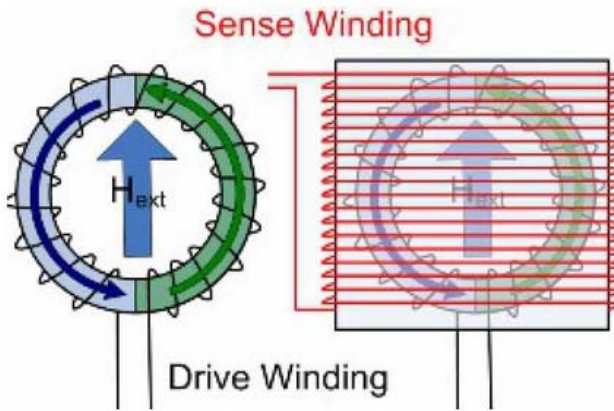
☐ Surprises, tricky data corrections and future challenges

Measuring DC/LF Magnetic Fields in space

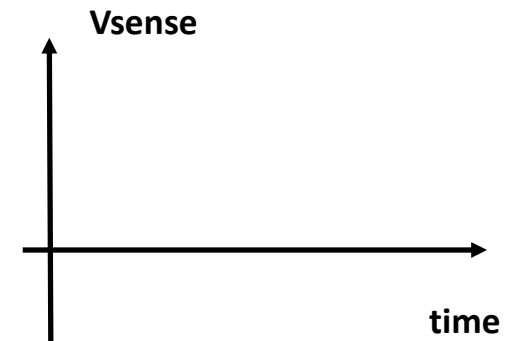
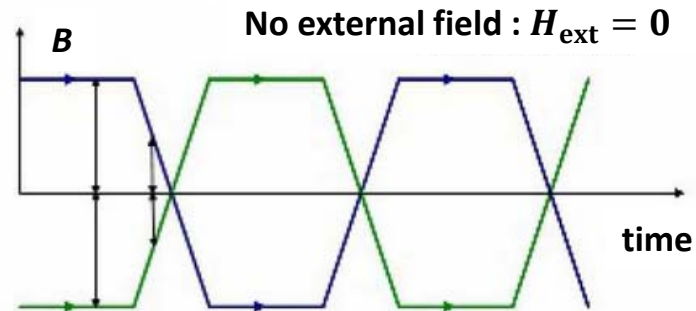
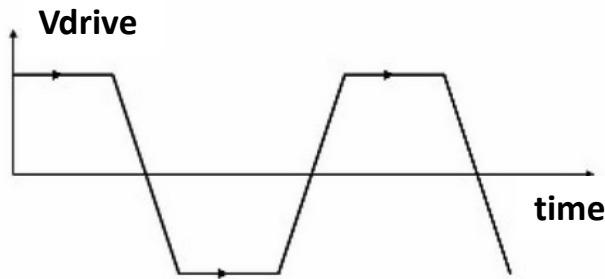
Flux Gate Magnetometer

Nobody knows... best theory is that it is little pixies inside the cores.

Timothy "Tim" Horbury
Imperial College, UK



ring cores of a highly magnetically permeable alloy around which are wrapped two coil windings

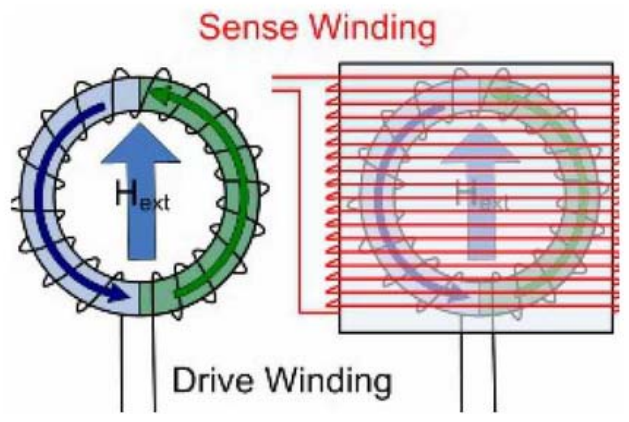
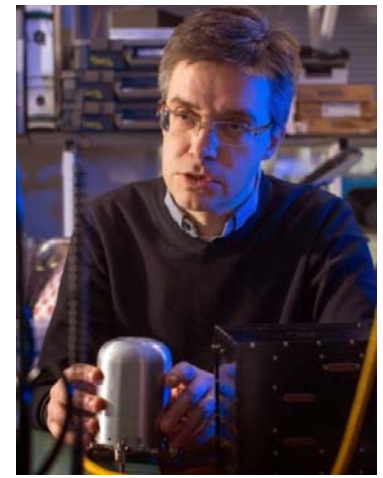


Measuring DC/LF Electric Field in space

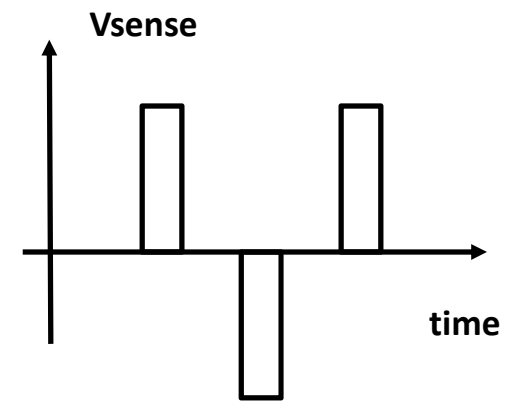
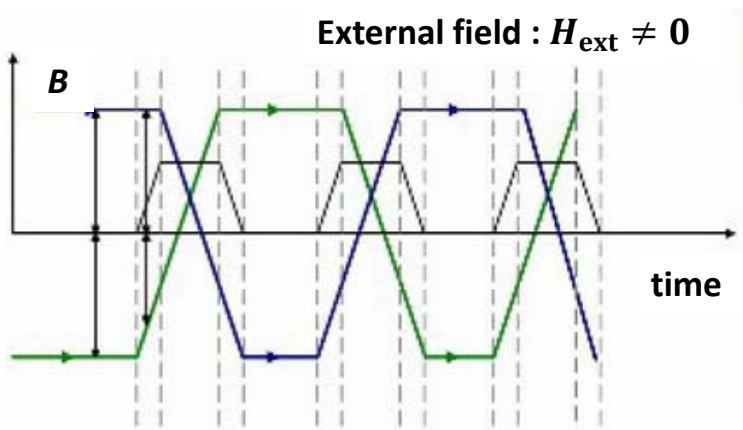
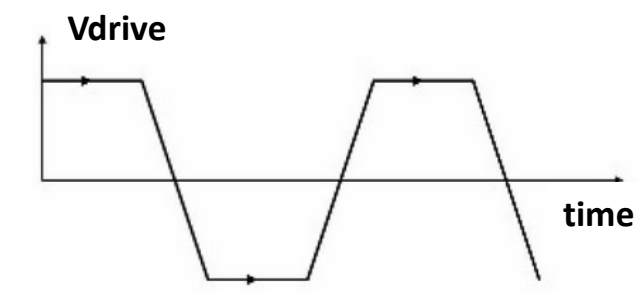
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ring cores of a highly magnetically permeable alloy around which are wrapped two coil windings



Measuring AC Magnetic Fields in space

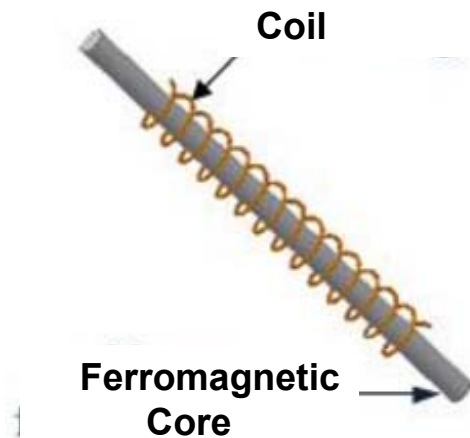
The Search Coil Magnetometer (SCM) is an inductive sensor based on Faraday's law of induction.

$$\mathcal{E} = -N \frac{d\phi}{dt} = -NS \frac{dB}{dt}$$

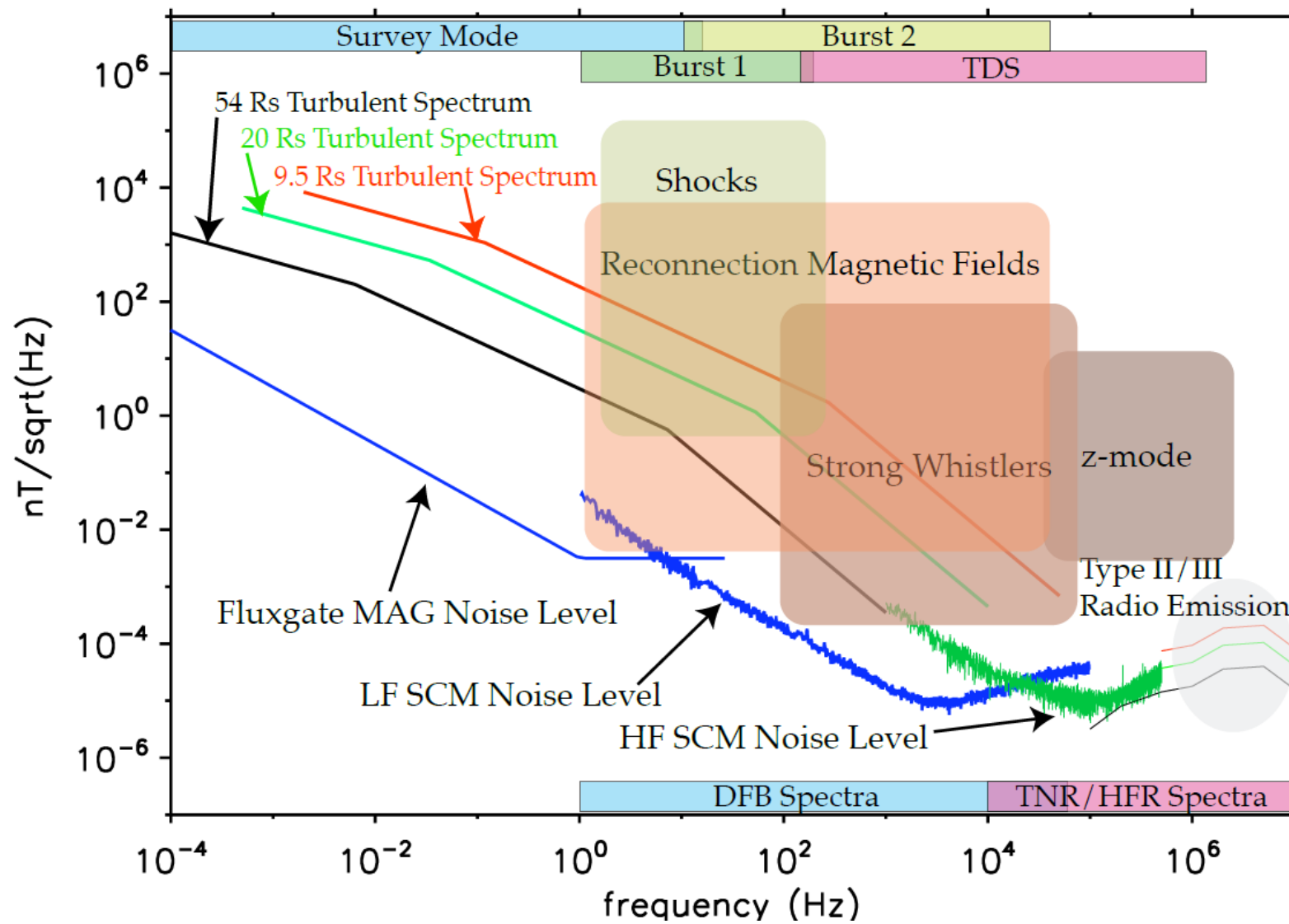
Vladimir "Volodya" Krasnoselskikh
LPC2E Orléans, France



SCMs On Solar Orbiter
and Solar Probe Plus



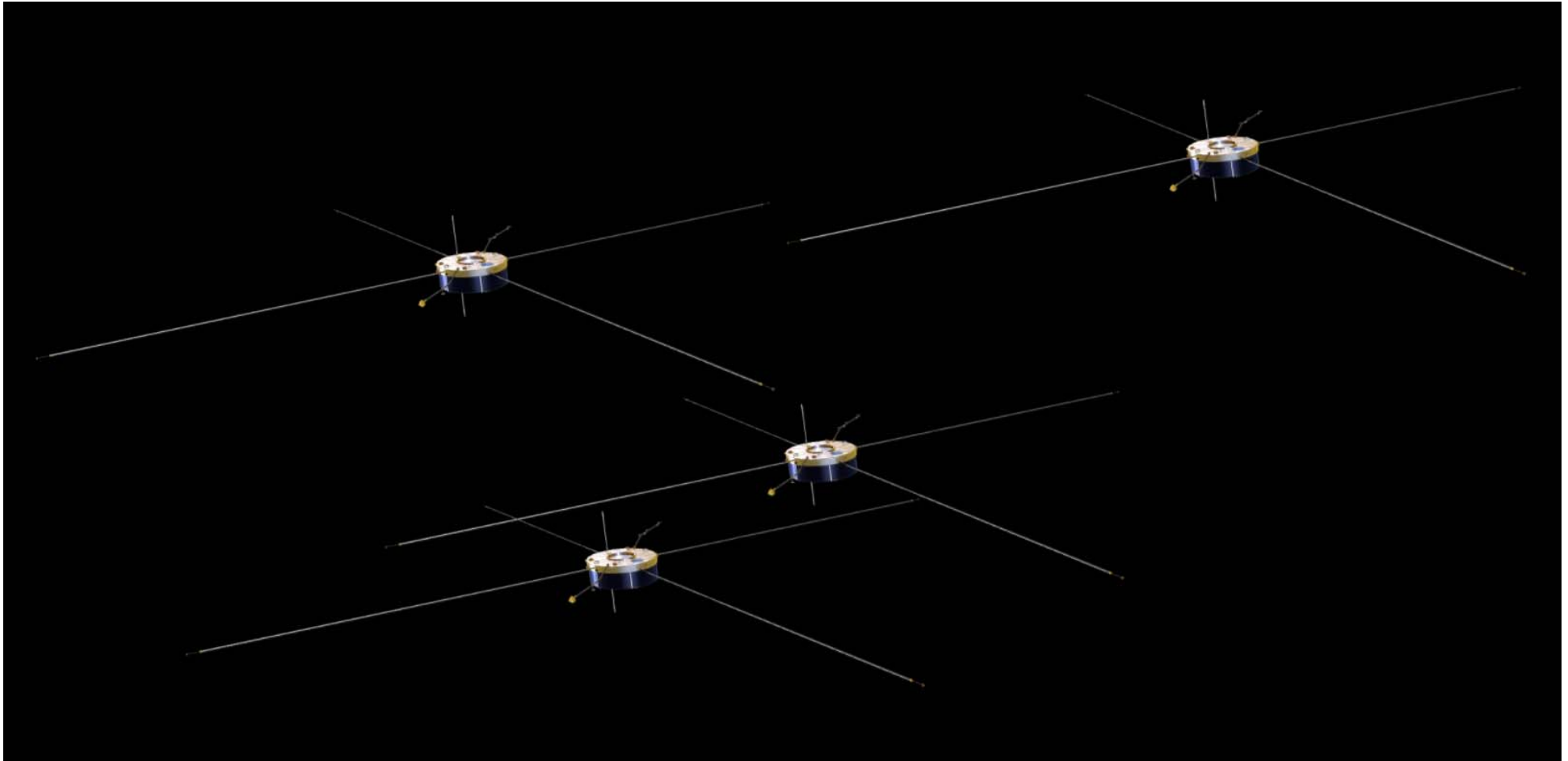
Maksimovic : Space plasmas measurement techniques

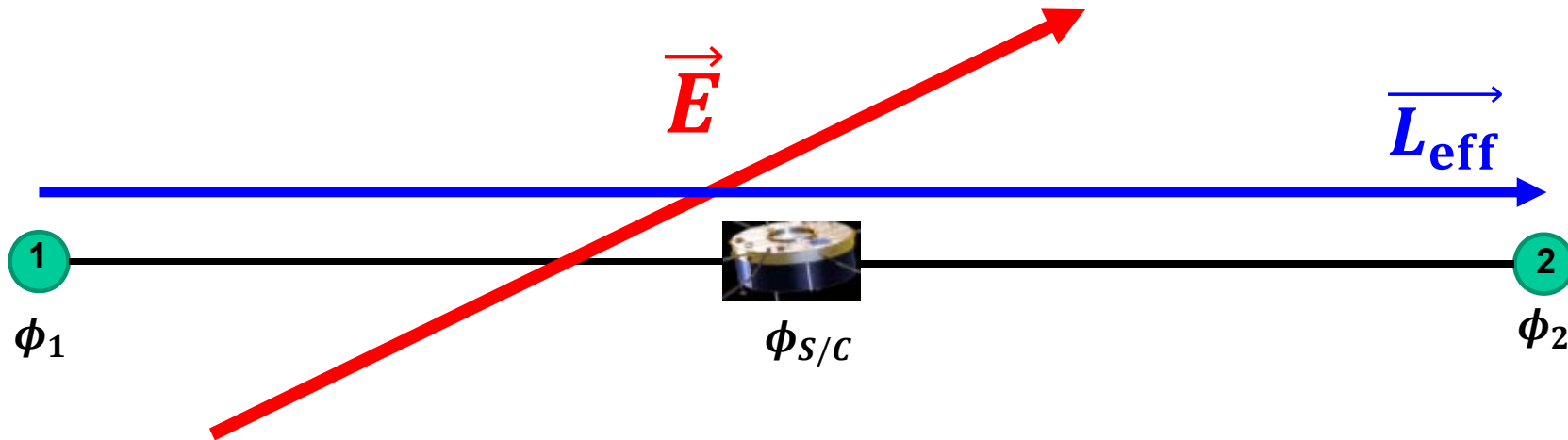


Maksimovic : Space plasmas measurement techniques



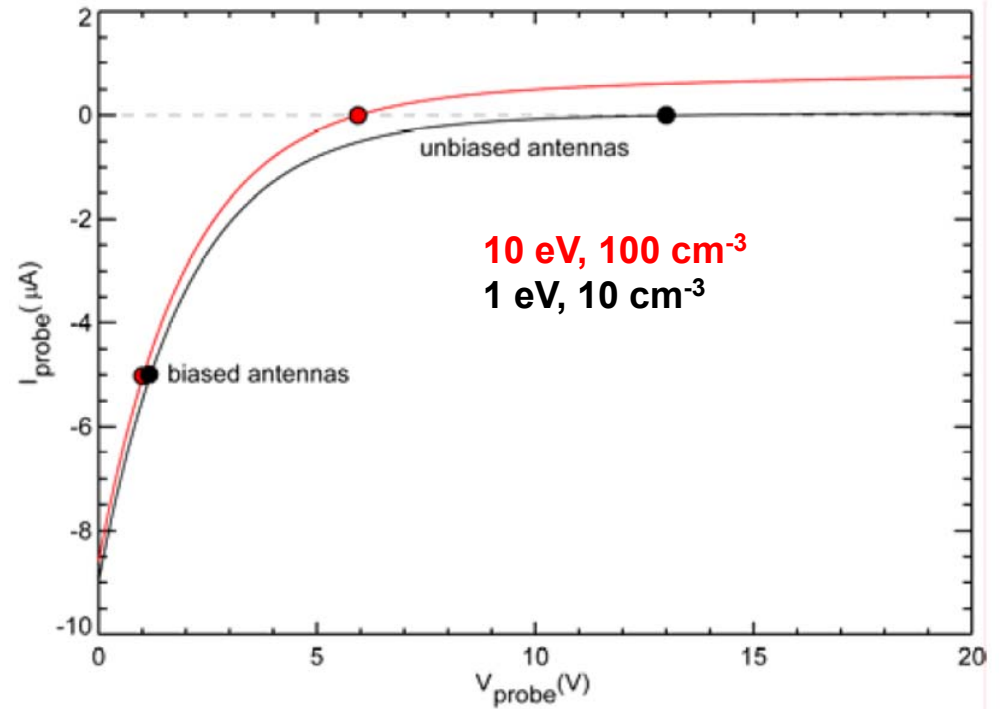
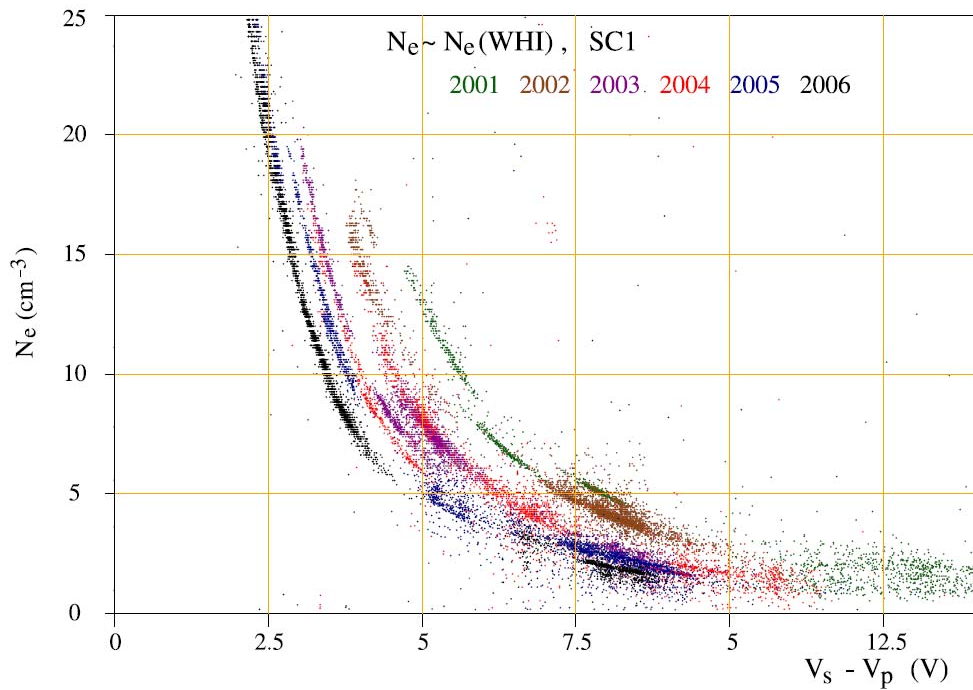
Measuring DC/LF Electric Field in space





- If the two probes are equally illuminated then $\phi_1 = \phi_2$ and not necessarily equal to $\phi_{S/C}$
- Actually we do not measure ϕ_1 and ϕ_2 but rather $V_1 = \phi_{S/C} - \phi_1$ and $V_2 = \phi_{S/C} - \phi_2$ where we use the S/C as the potential ground
- If the ground is the same for the two probes then $V_1 - V_2 = \phi_2 - \phi_1 = 0$
- If an external \vec{E} is applied then $\phi_1 \neq \phi_2$
- Actually $\phi_2 - \phi_1 = \delta_{\phi_E} = \vec{E} \cdot \vec{L}_{\text{eff}}$
- So if the experimental setup is appropriate then :
 - $V_1 - V_2 = \delta_{\phi_E} = \vec{E} \cdot \vec{L}_{\text{eff}} !$

But life is not so simple because typical electric fields are tiny (a few mV/m) !



Adding a Biasing current on the probe will solve the problem of the density fluctuations in the medium

$$J_e(\phi) = J_{ph}(\phi) + J_{BIAS}$$

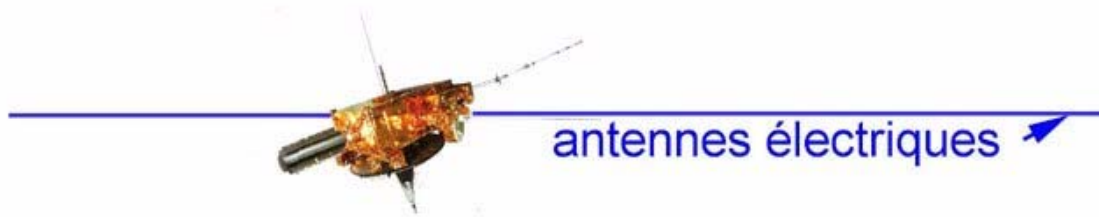
Maksimovic : Space plasmas me

- So if the experimental setup is appropriate and if one injects a BIAS current in the probes then

- $V_1 - V_2 = \delta\phi_E = \vec{E} \cdot \vec{L}_{eff} !$

- $\frac{V_1 + V_2}{2} = \phi_{S/C} - \cancel{\phi_2} + \frac{\delta\phi_E}{2} \sim \phi_{S/C}$

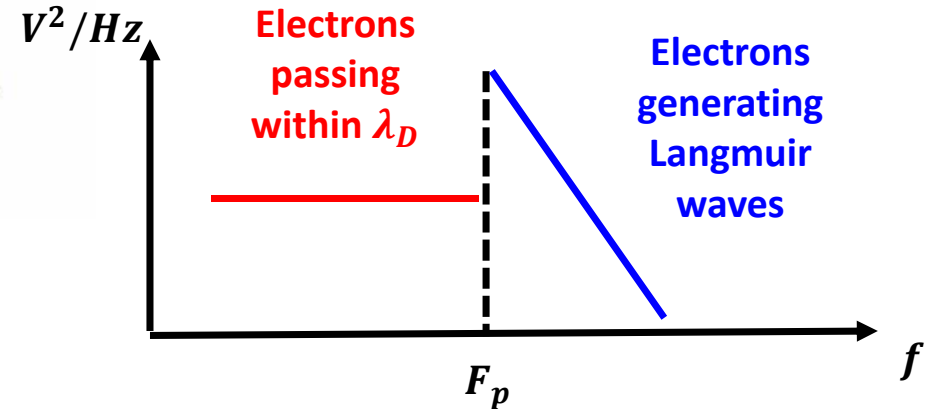
Measuring AC Electric Field: The Thermal Noise as an example



In drifting plasma, spectral density at the antenna ports:

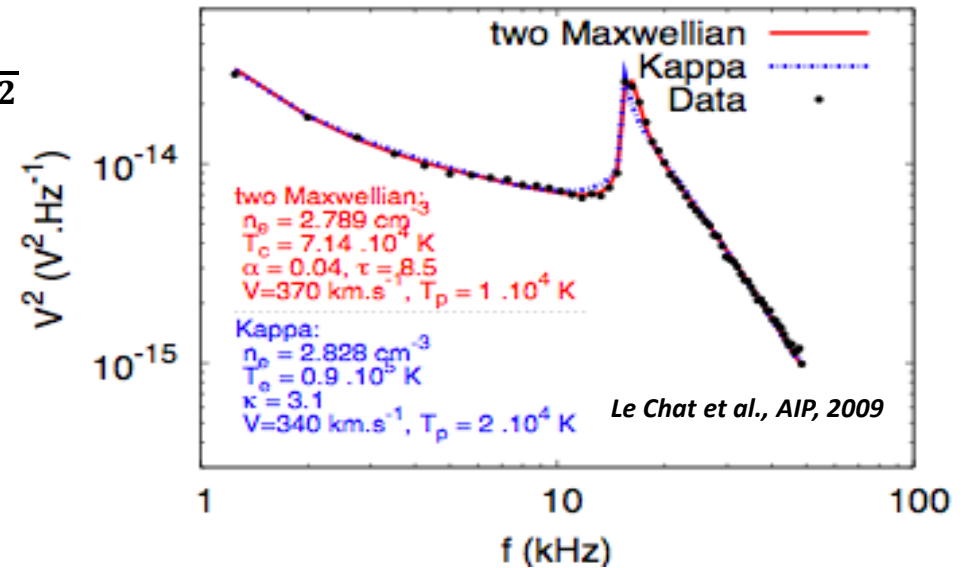
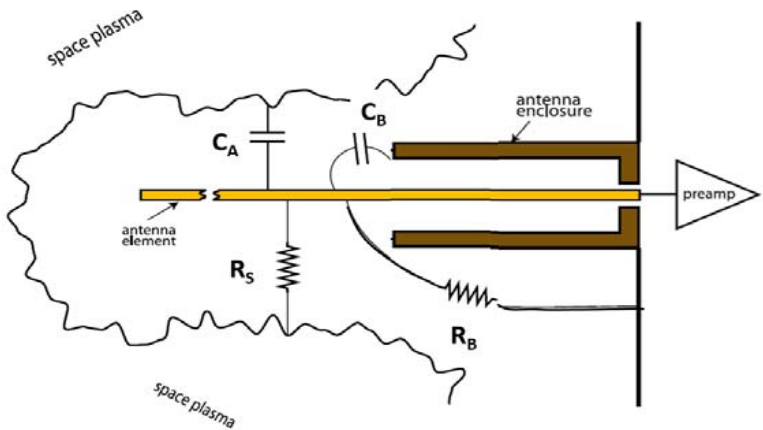
$$V_{\omega}^2 = \frac{2}{(2\pi)^3} \int d^3k \left| \frac{\mathbf{k} \cdot \mathbf{J}}{k} \right|^2 E^2(\mathbf{k}, \omega - \mathbf{k} \cdot \mathbf{V})$$

antenna plasma



Resistive $G = \frac{1}{(1 + R_S/R_B)^2}$

$$V_R^2(\omega) = \frac{V_{\omega}^2}{(1 + C_B/C_A)^2}$$



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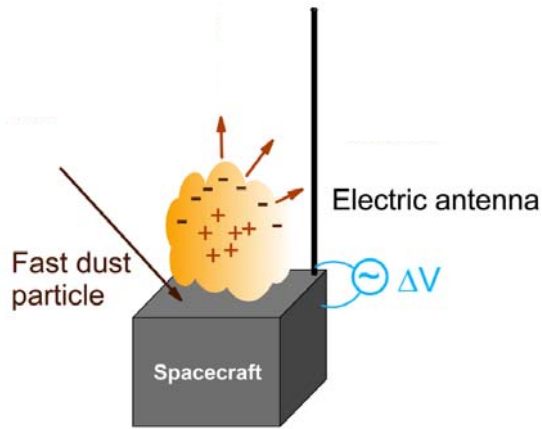
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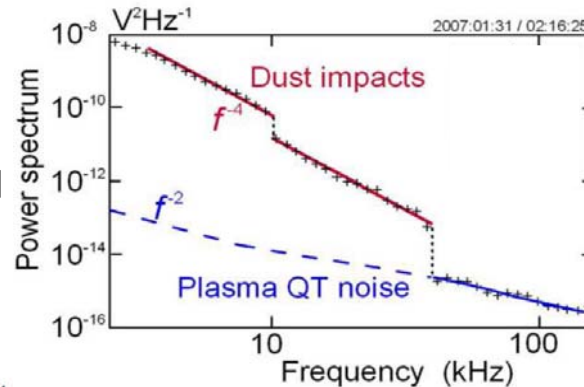
☐ **Surprises, tricky data corrections and future challenges**

Surprises : when an instrument allow unexpected observations

Nanodusts with STEREO



Spectral domain

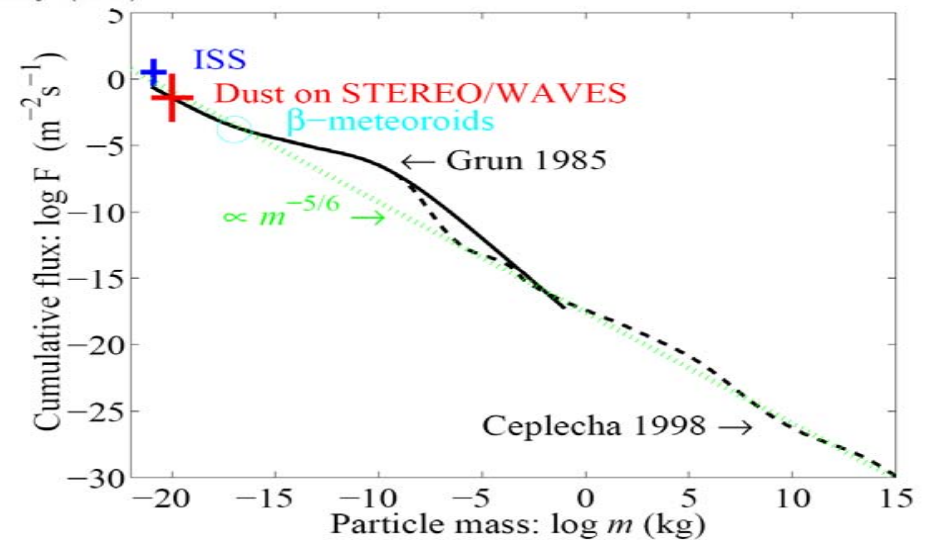
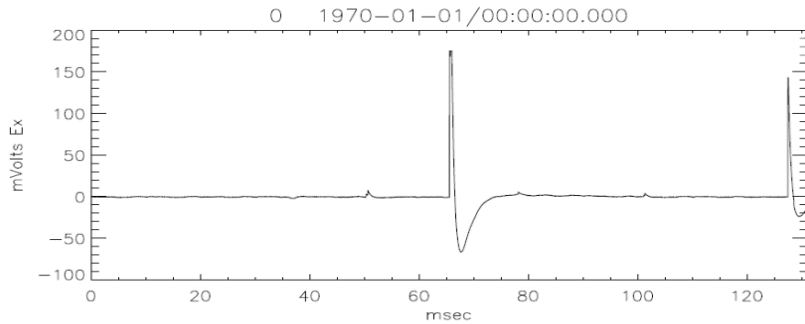


→ A nanoparticle @ 300 km/s ~ a grain of mass 10^4 greater @ 20 km/s

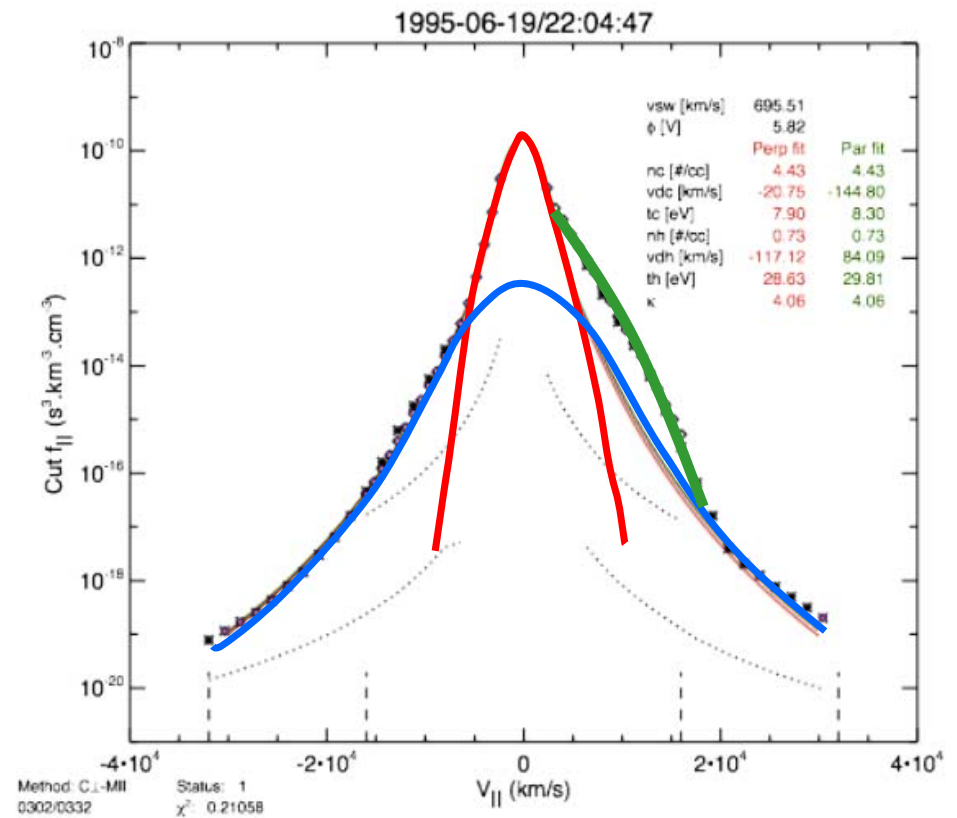
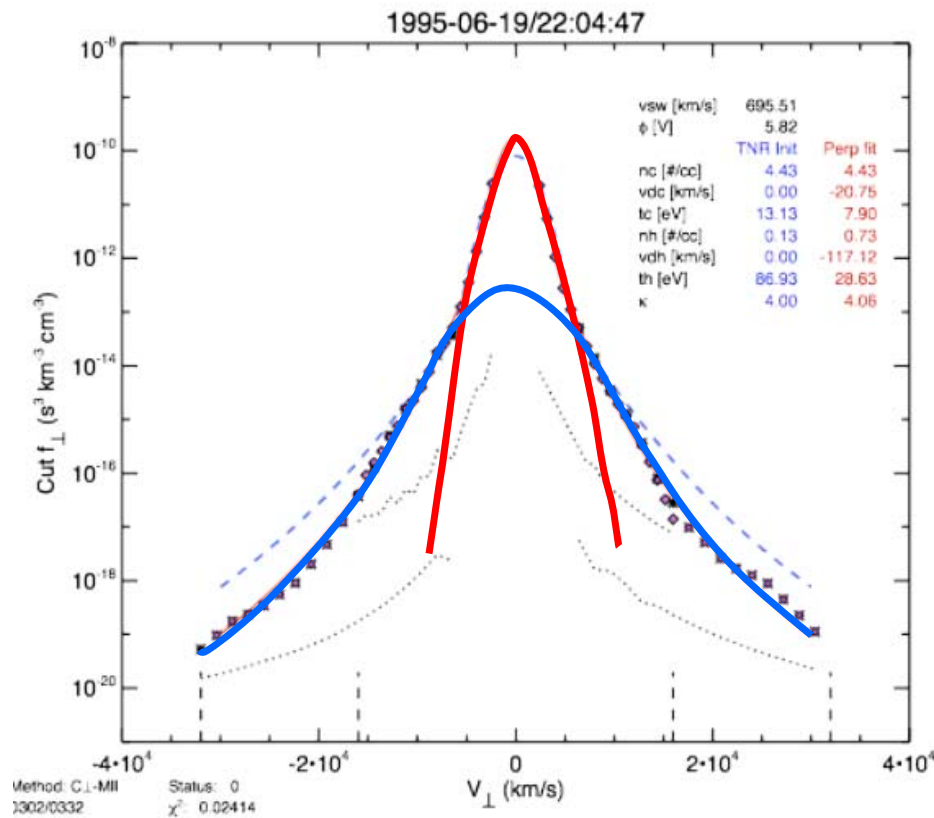
Picked-up by the VXB field

Released charge : $Q \simeq 0.7m^{1.02}v^{3.48}$

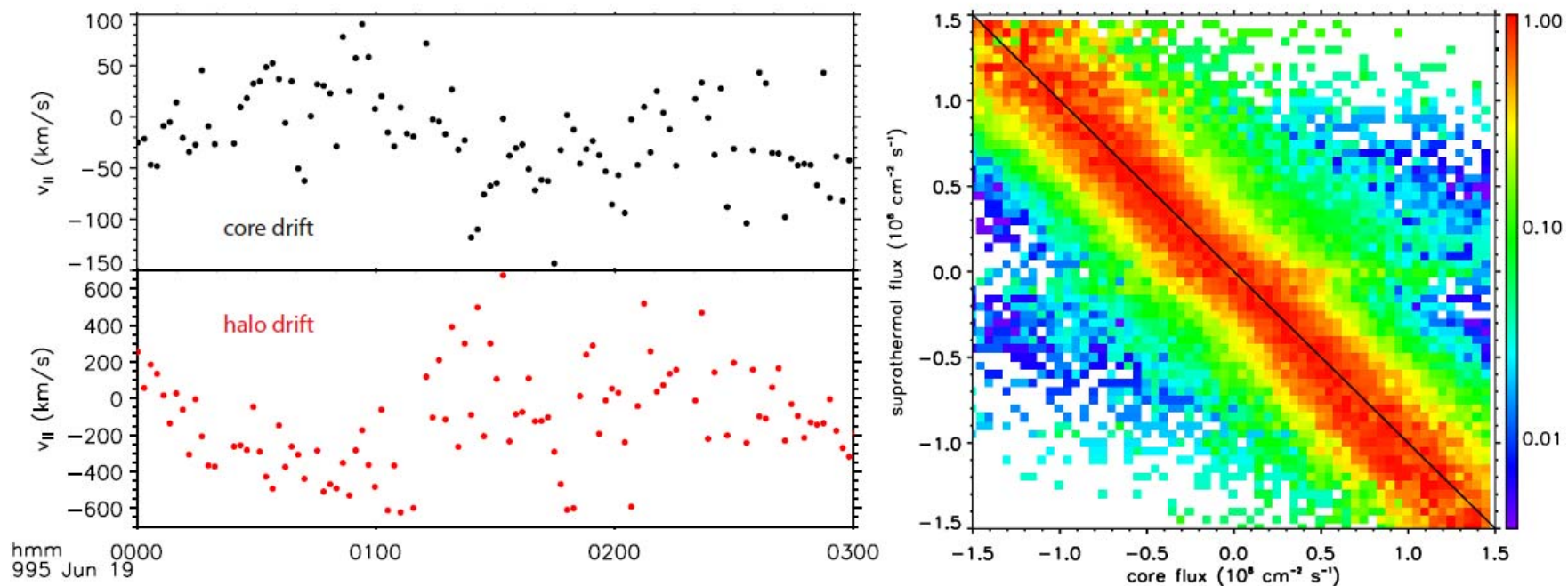
Induced voltage pulse on S/C of capacitance C : $\delta V \sim -Q/C$



Tricky data corrections techniques

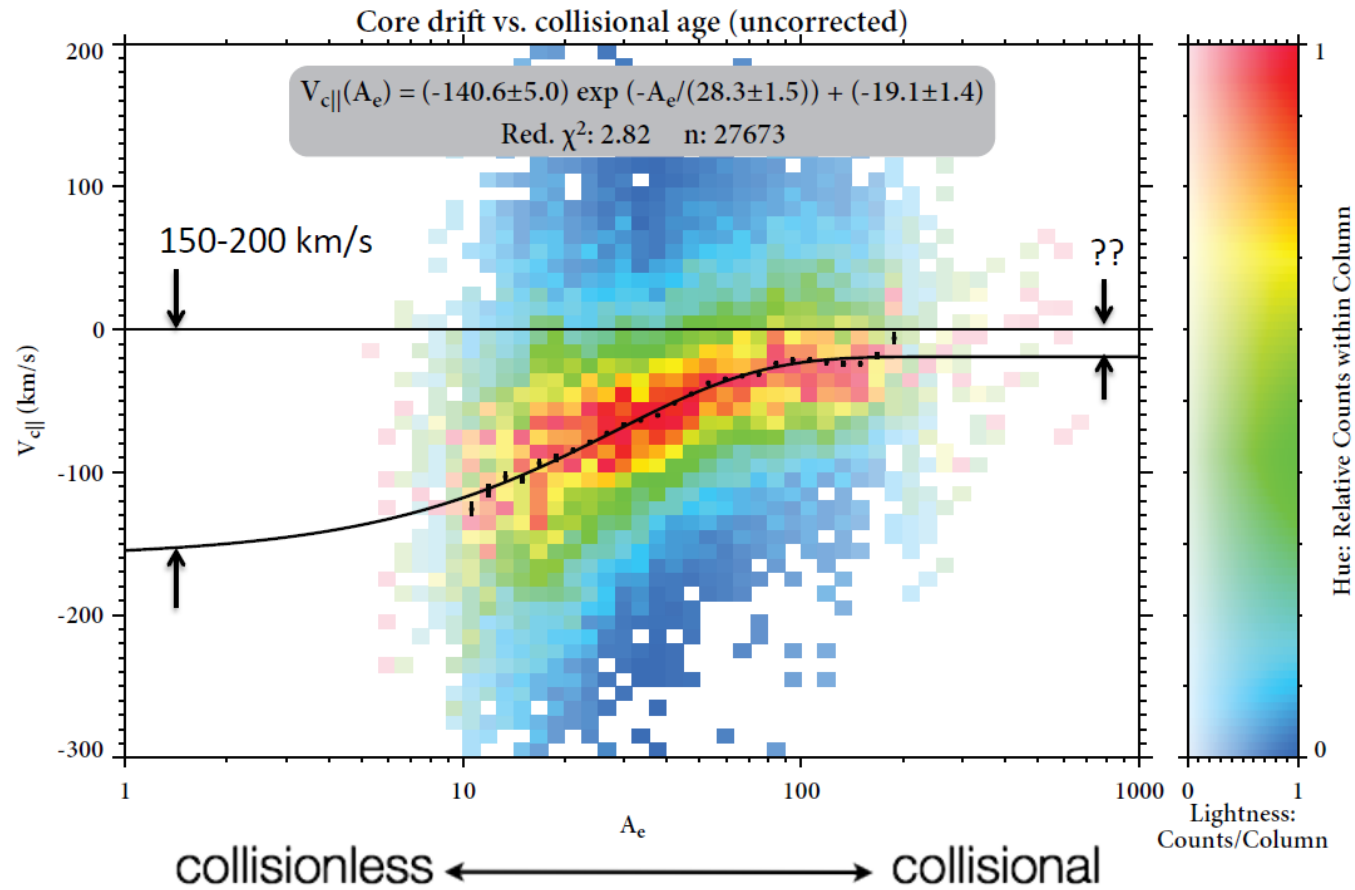


Current balance in the proton frame

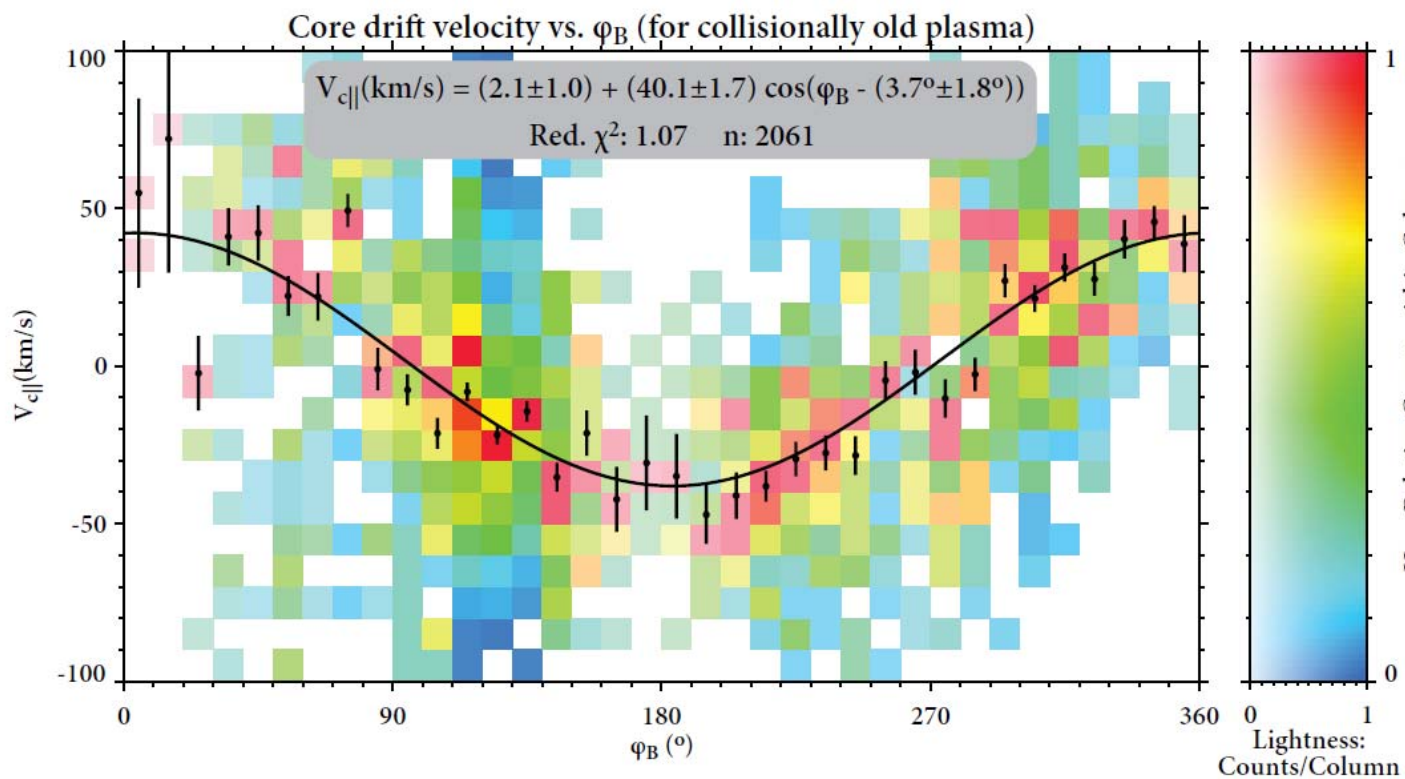


In the charge-center ($\sim \text{cm}$) frame, we expect zero net current: $n_c v_c + n_h v_h + n_s v_s = 0$, which seems to be so...

Core electron-proton (||) drift

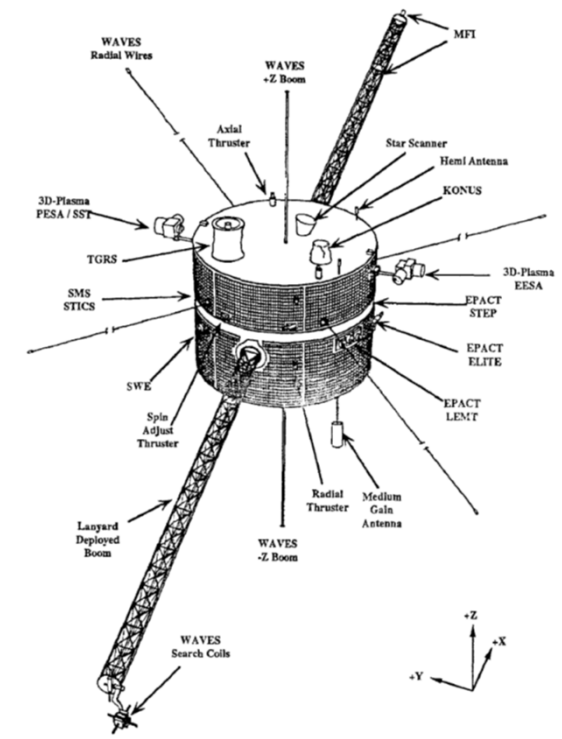


Sunward drift in collisional regime



Dipolar S/C potential correction

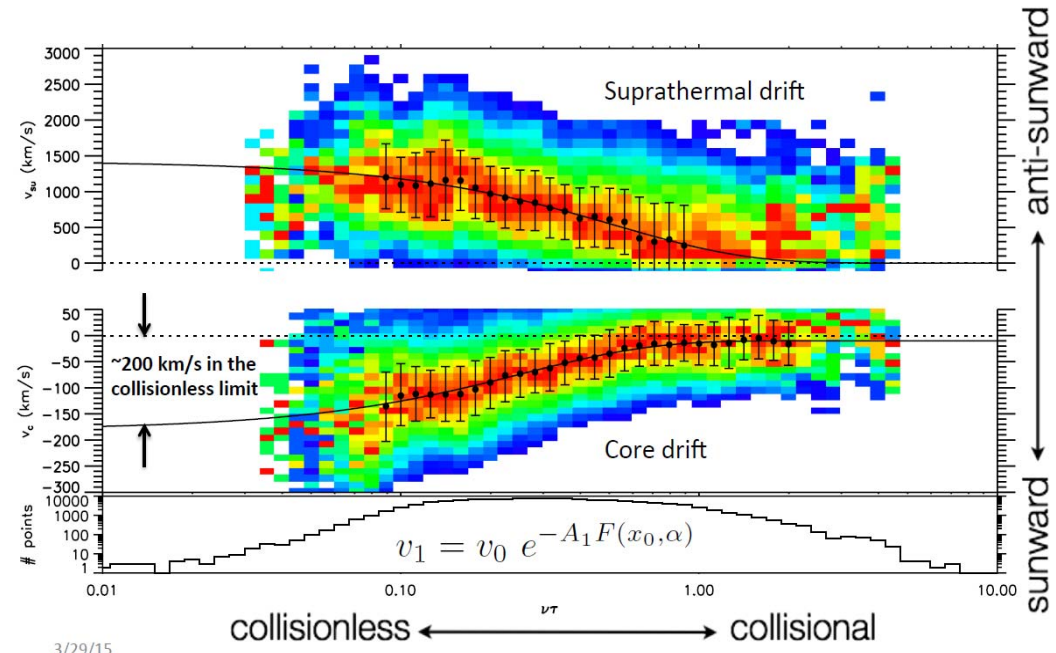
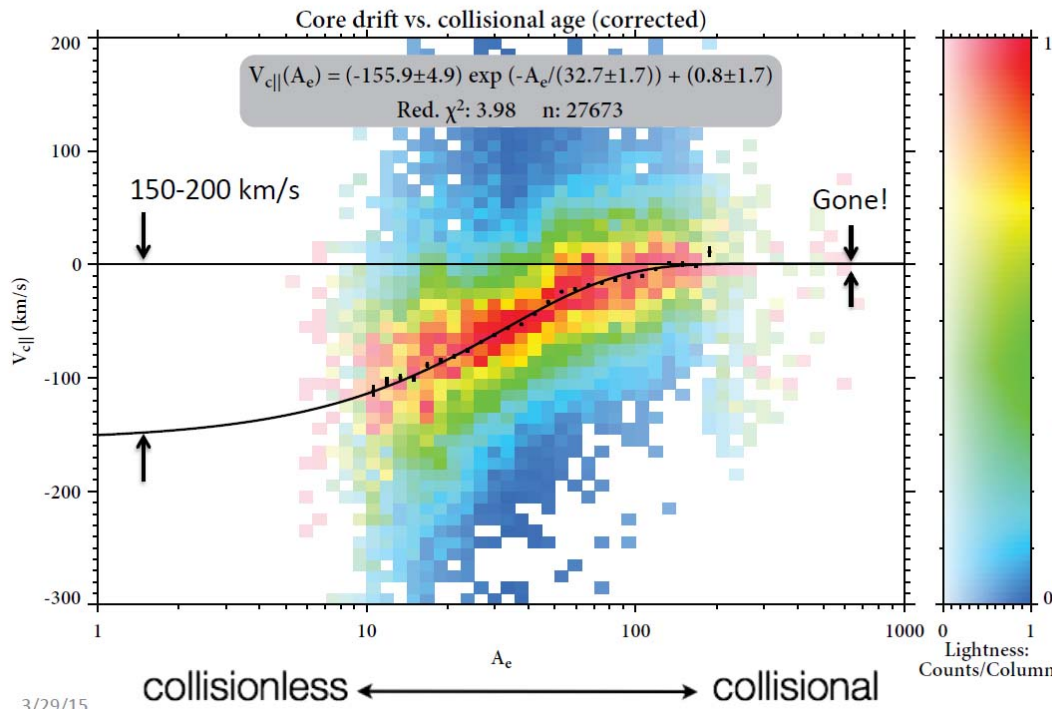
Maksimovic : Space plasmas measurement techniques



$$\phi(\varphi, R) = K \left(\frac{1}{R} - \frac{A \cos \varphi}{R^2} \right)$$

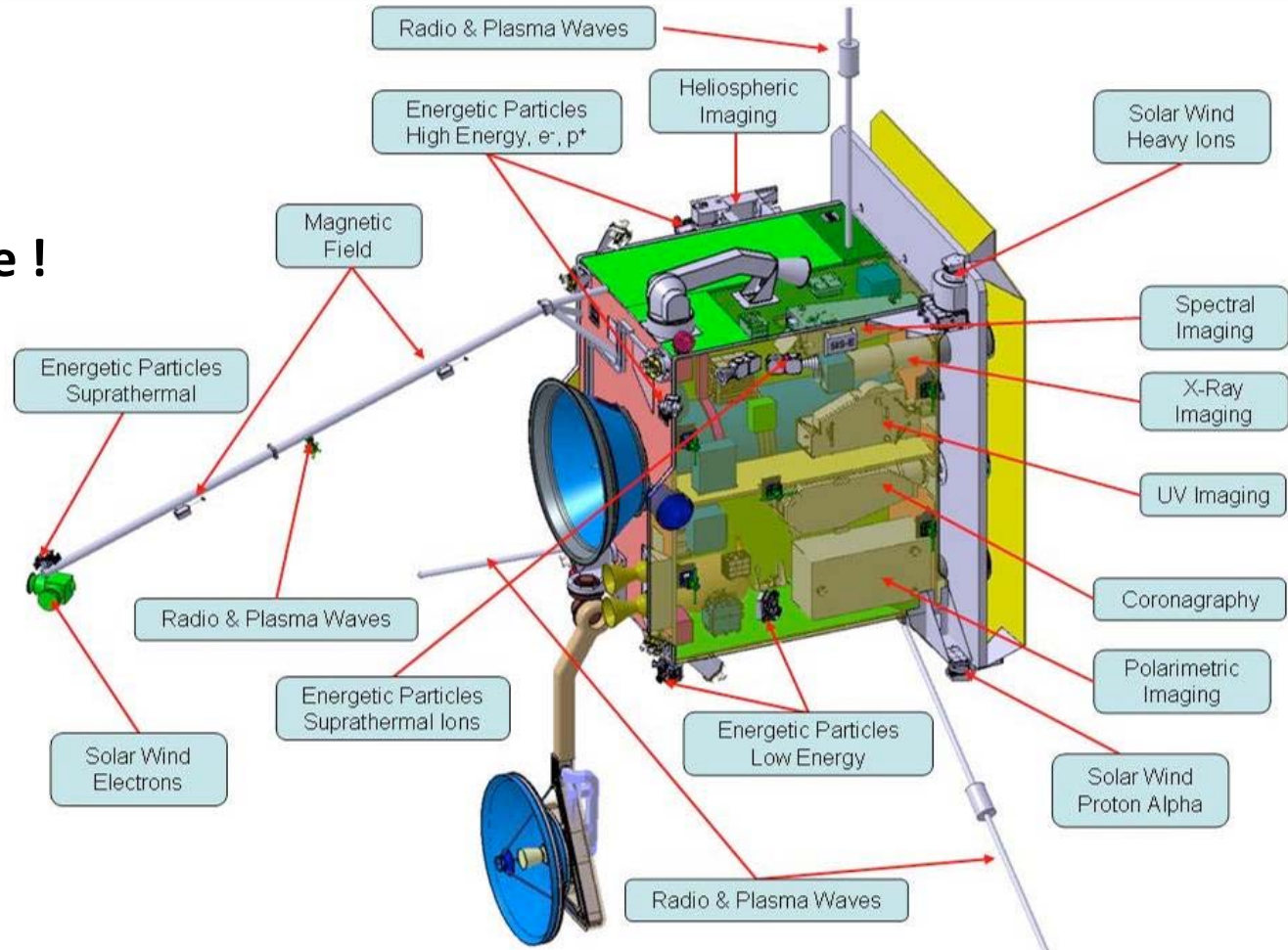
$$\phi(\varphi) = \phi_0 - \phi_1 \cos \varphi$$

Core electron-proton (||) drift

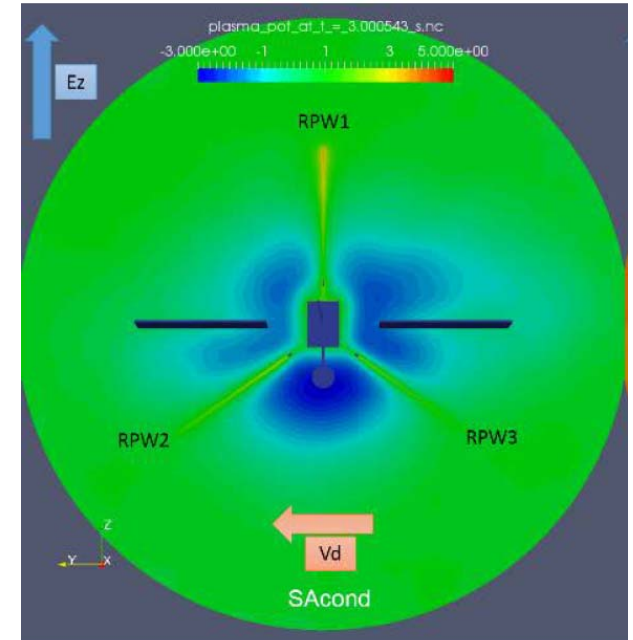
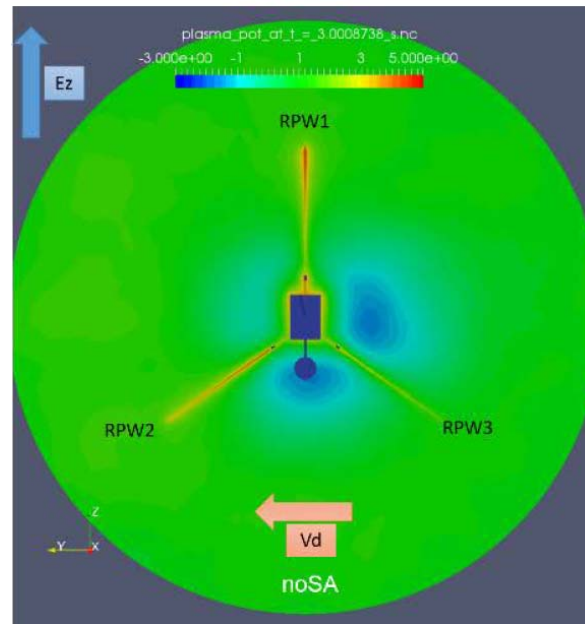
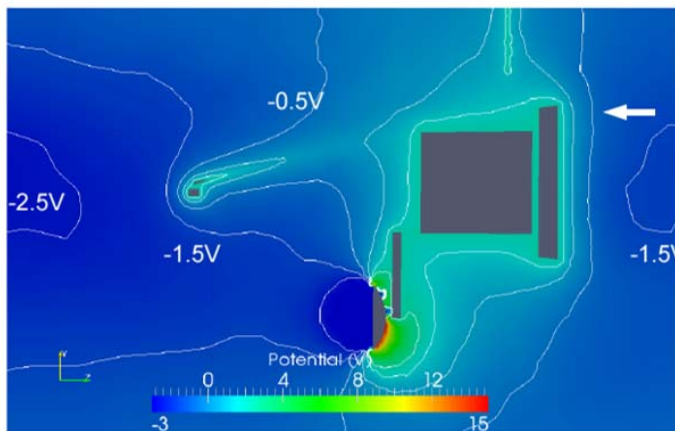
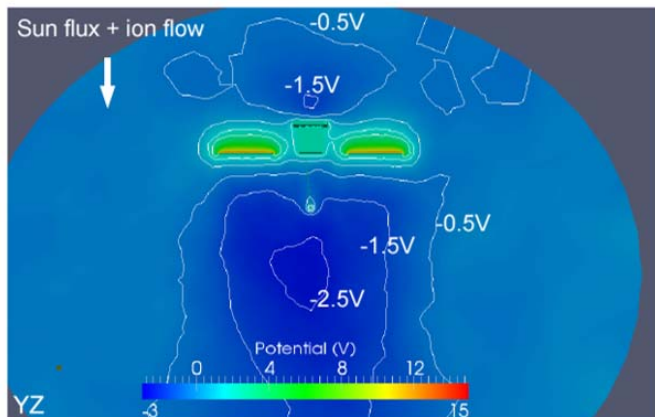


Some Challenges with Solar Orbiter

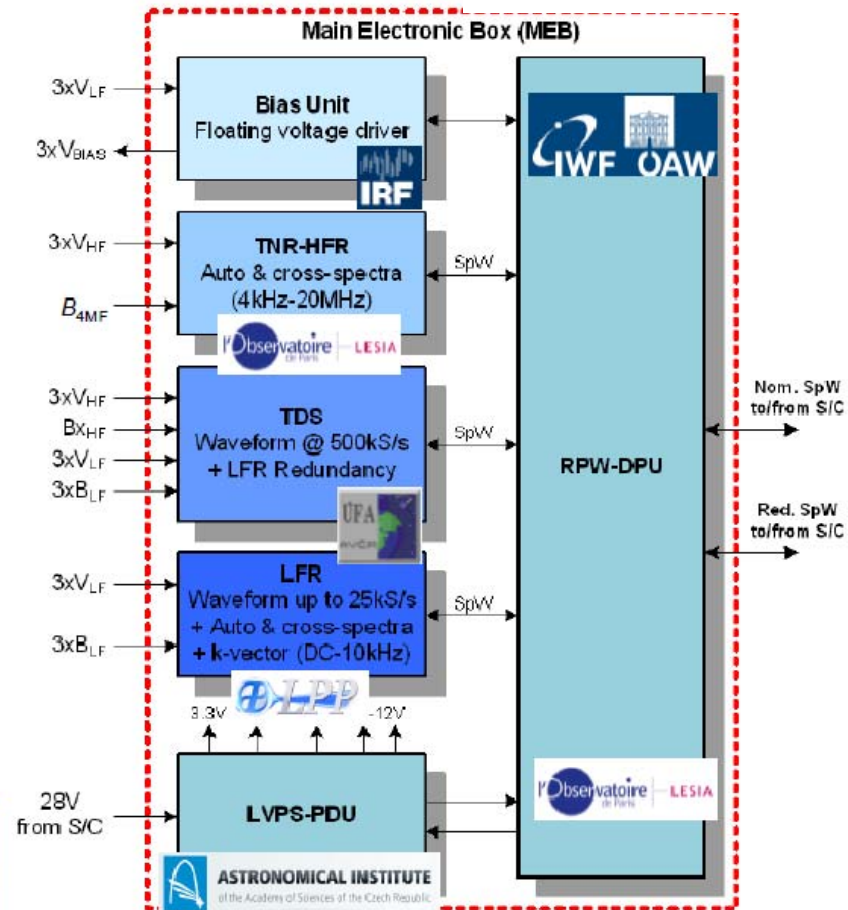
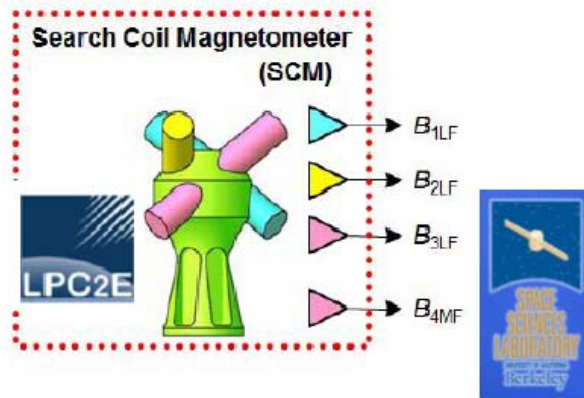
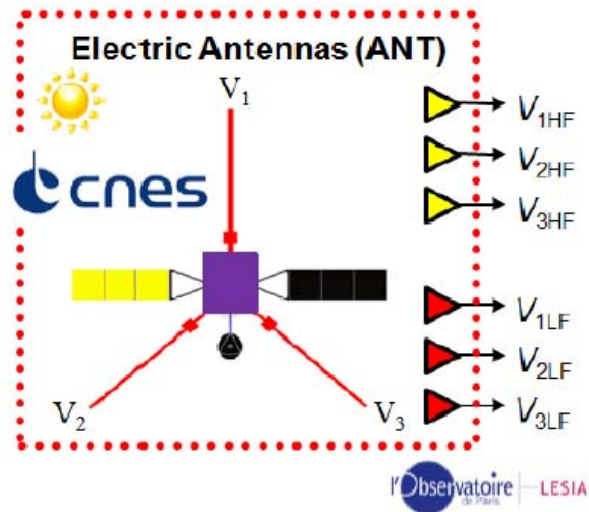
Integrating all 10 instruments on the S/C platform is a real challenge !



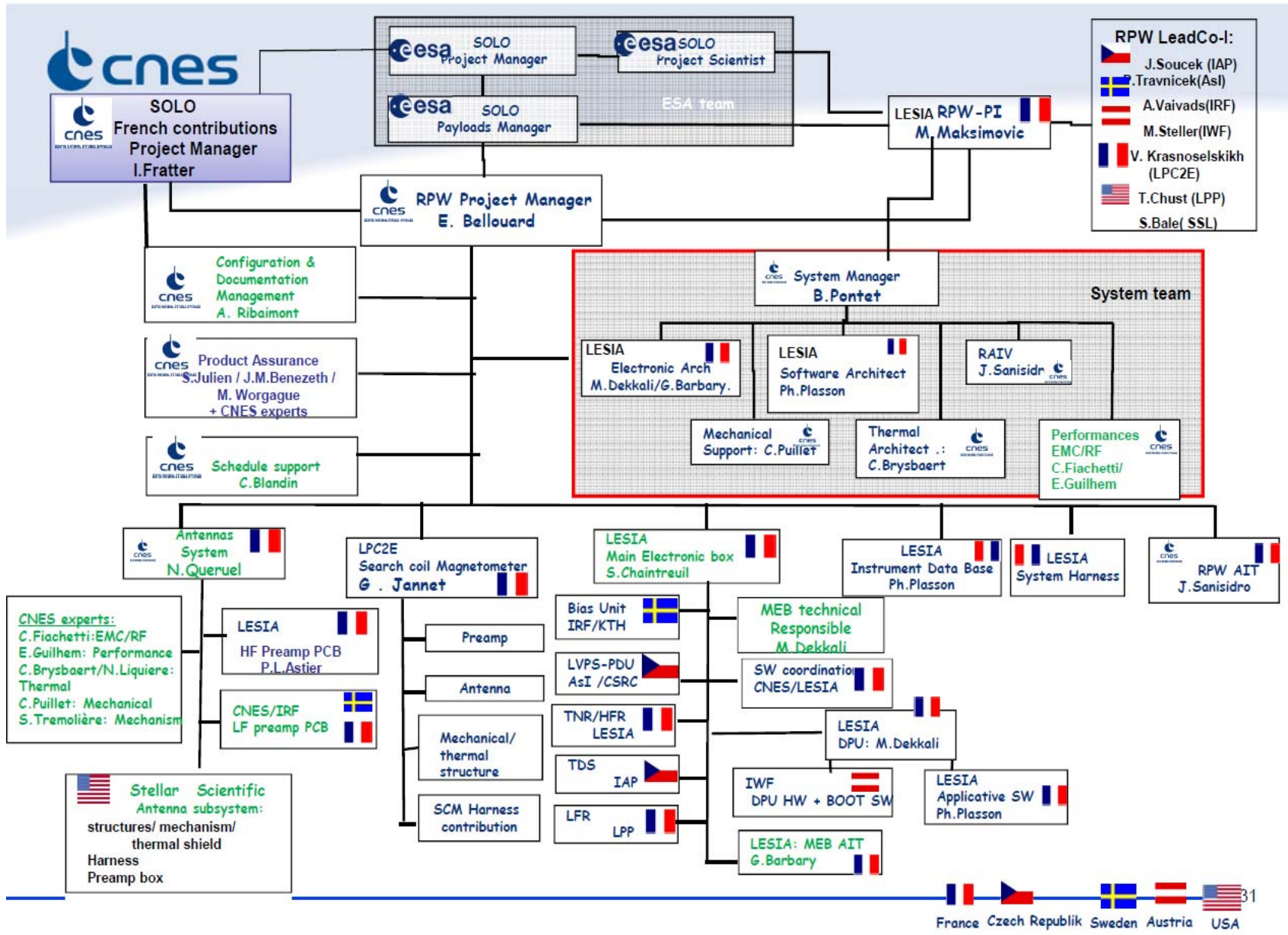
Understanding the electrostatic environment will be a challenge

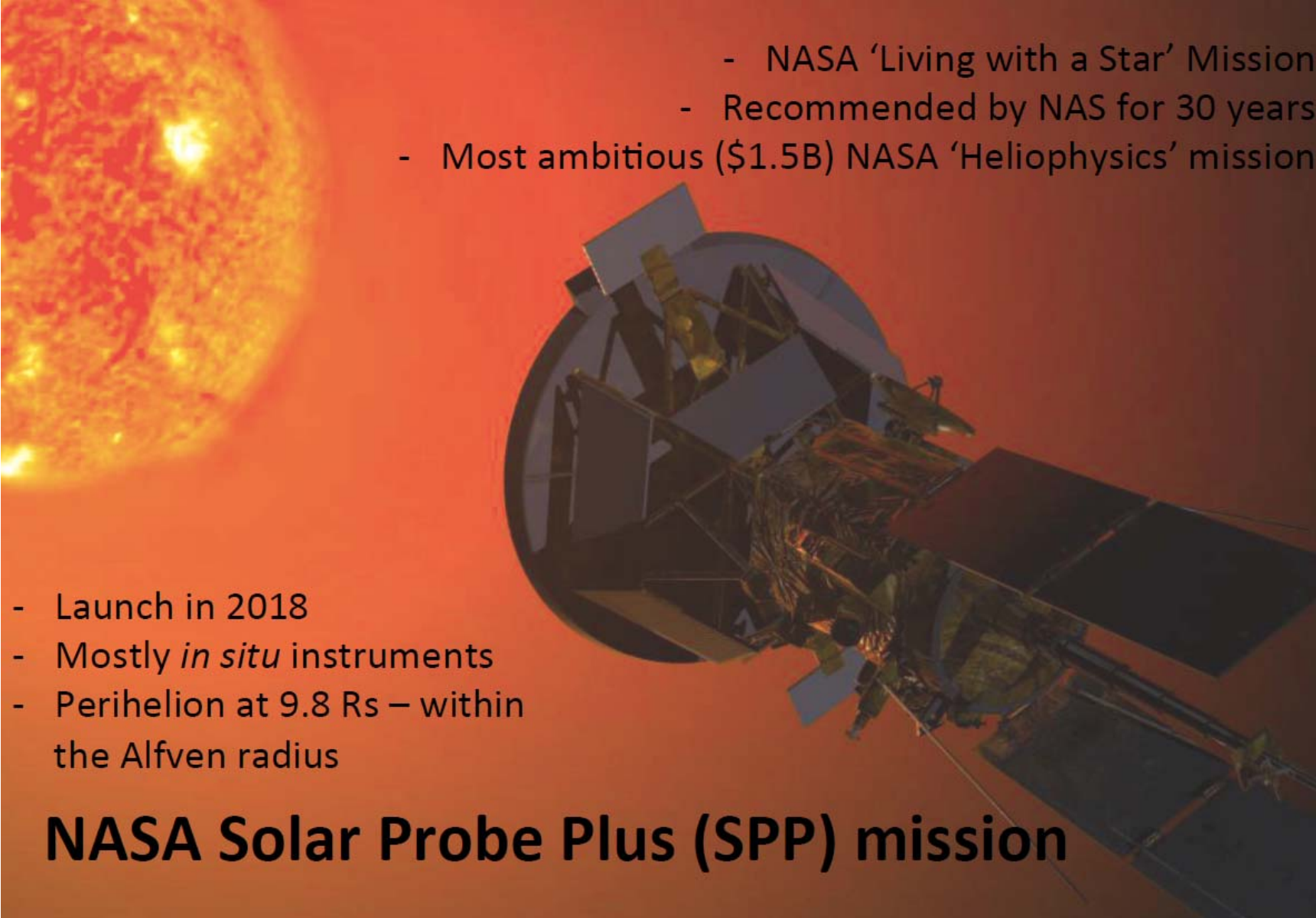


The Radio & Plasmas Waves instrument



DC Efield expertise & calibration, link to SPP



- 
- NASA 'Living with a Star' Mission
 - Recommended by NAS for 30 years
 - Most ambitious (\$1.5B) NASA 'Heliophysics' mission

- Launch in 2018
- Mostly *in situ* instruments
- Perihelion at 9.8 Rs – within the Alfvén radius

NASA Solar Probe Plus (SPP) mission