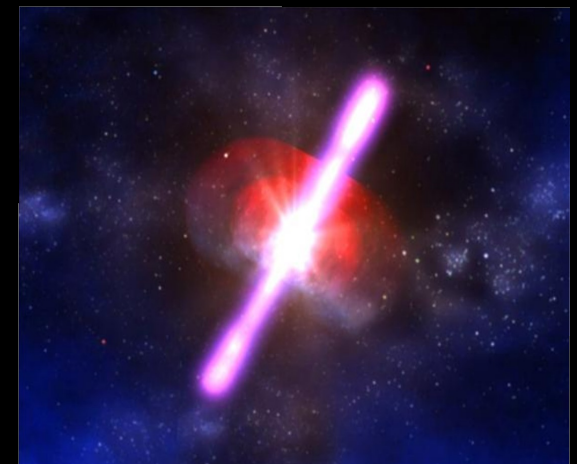
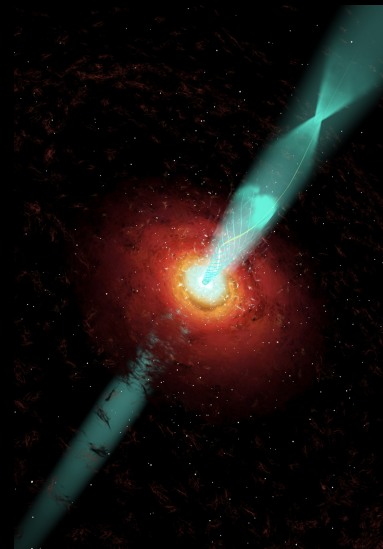
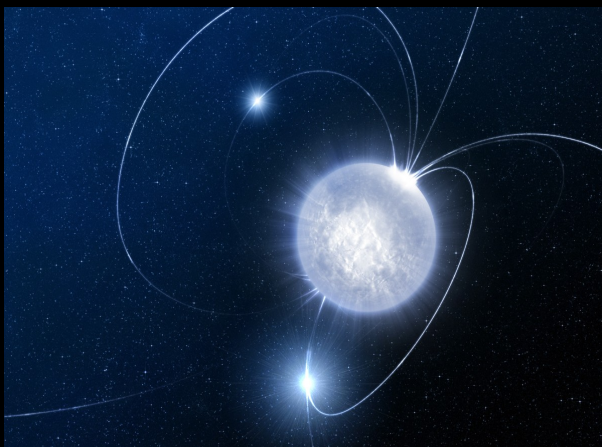


Particle acceleration in reconnection sites

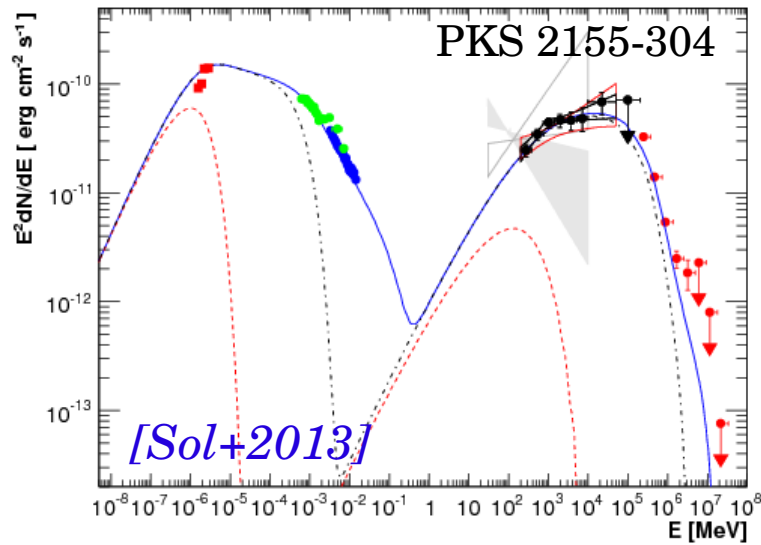
Benoît Cerutti

Institut de Planétologie et d'Astrophysique de Grenoble

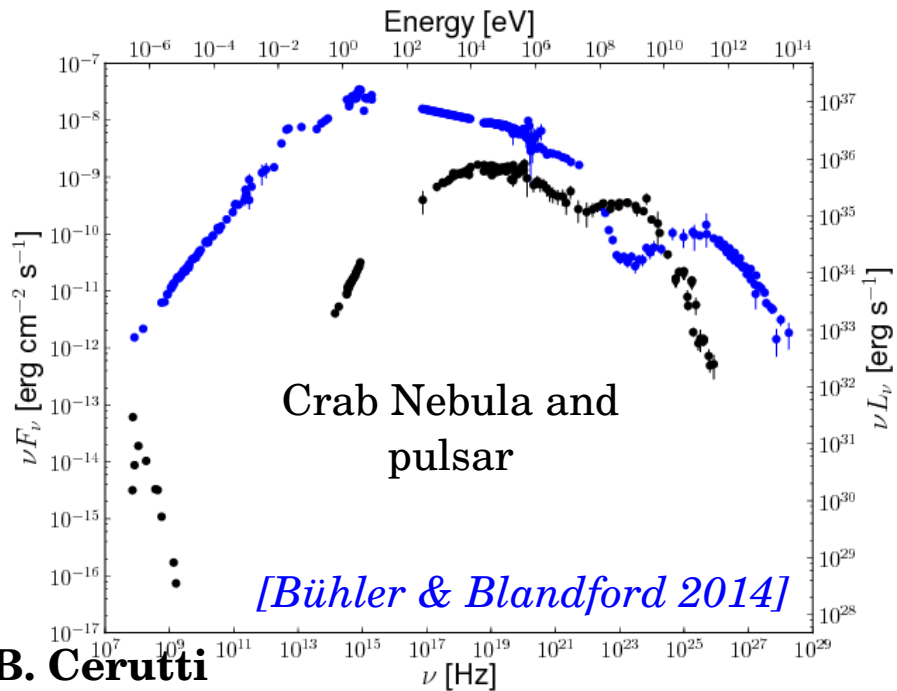


Broad non-thermal distributions

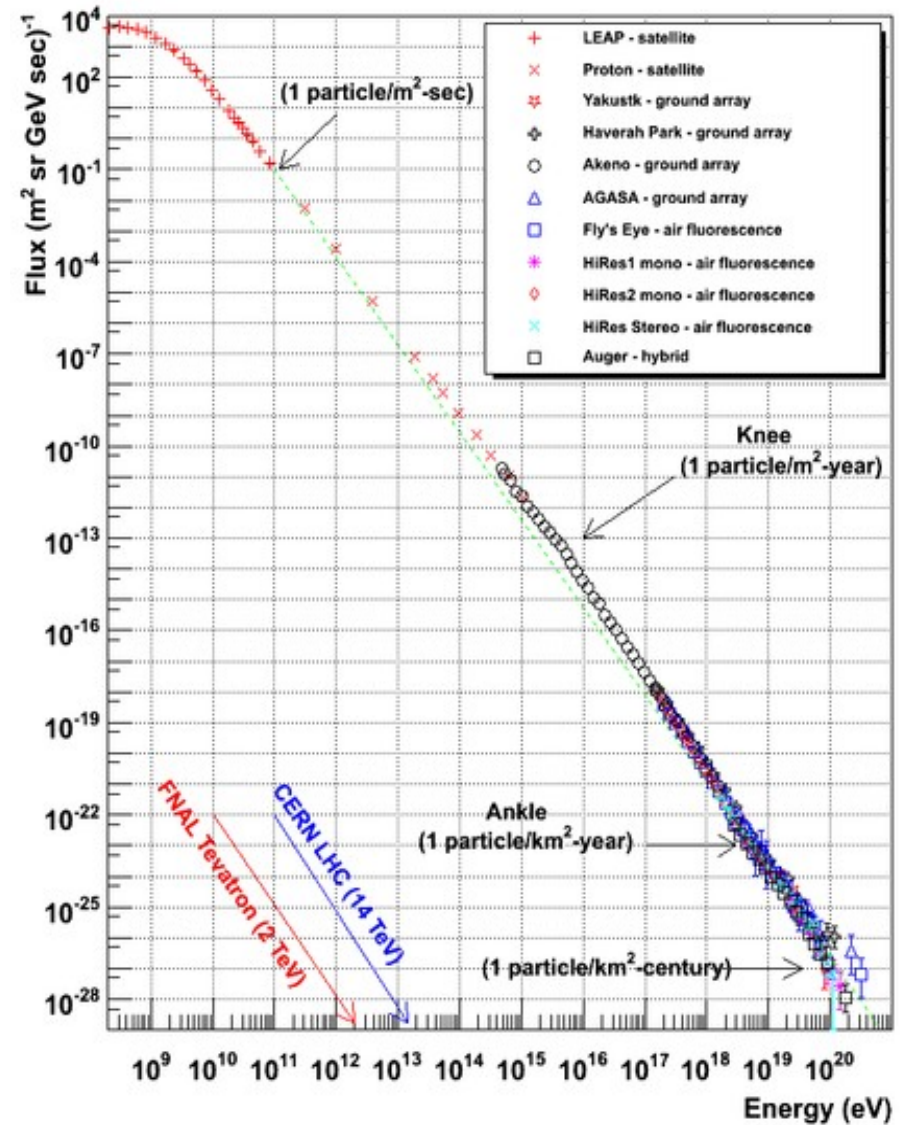
Blazars



Pulsars & Pulsar Wind Nebulae



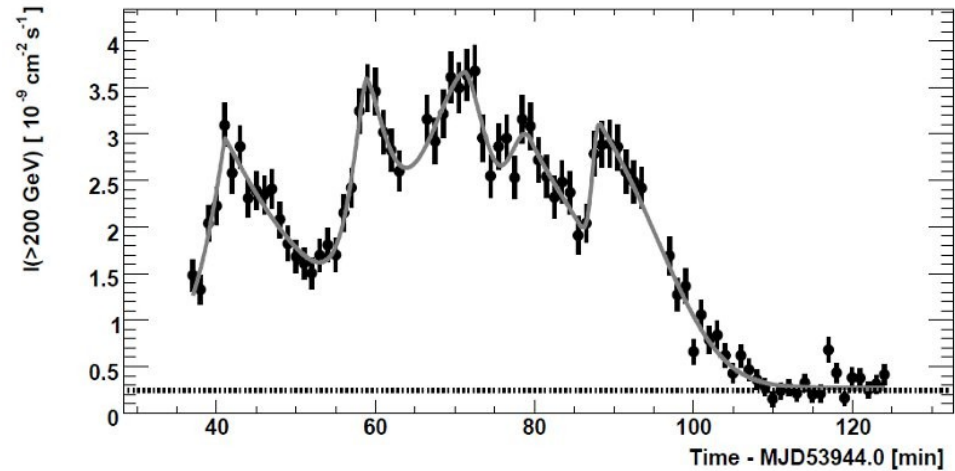
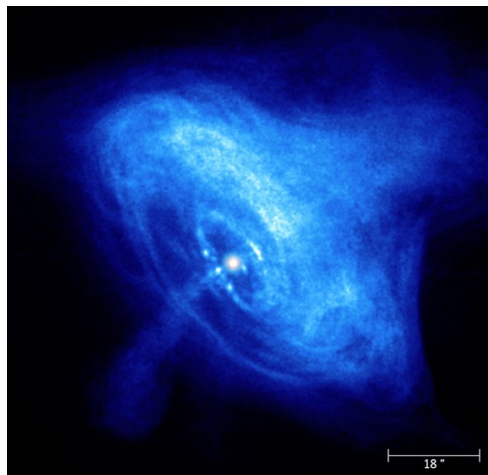
Cosmic Ray Spectra of Various Experiments



[<http://www.physics.utah.edu/~whanlon/spectrum.html>]

Diffuse shock acceleration cannot explain :

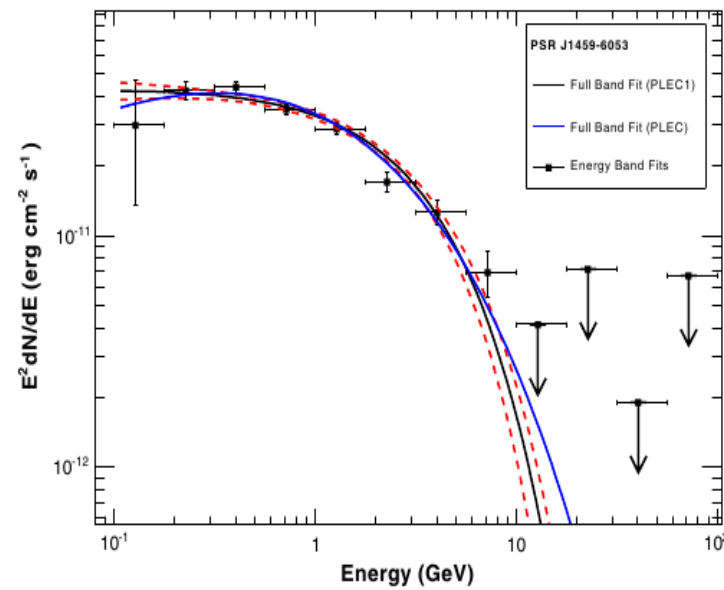
Ultra-rapid time variability
Too slow !



Aharonian et al. 2007 (H.E.S.S. Collaboration)

Particle acceleration at ultra-relativistic
... and magnetized shock.

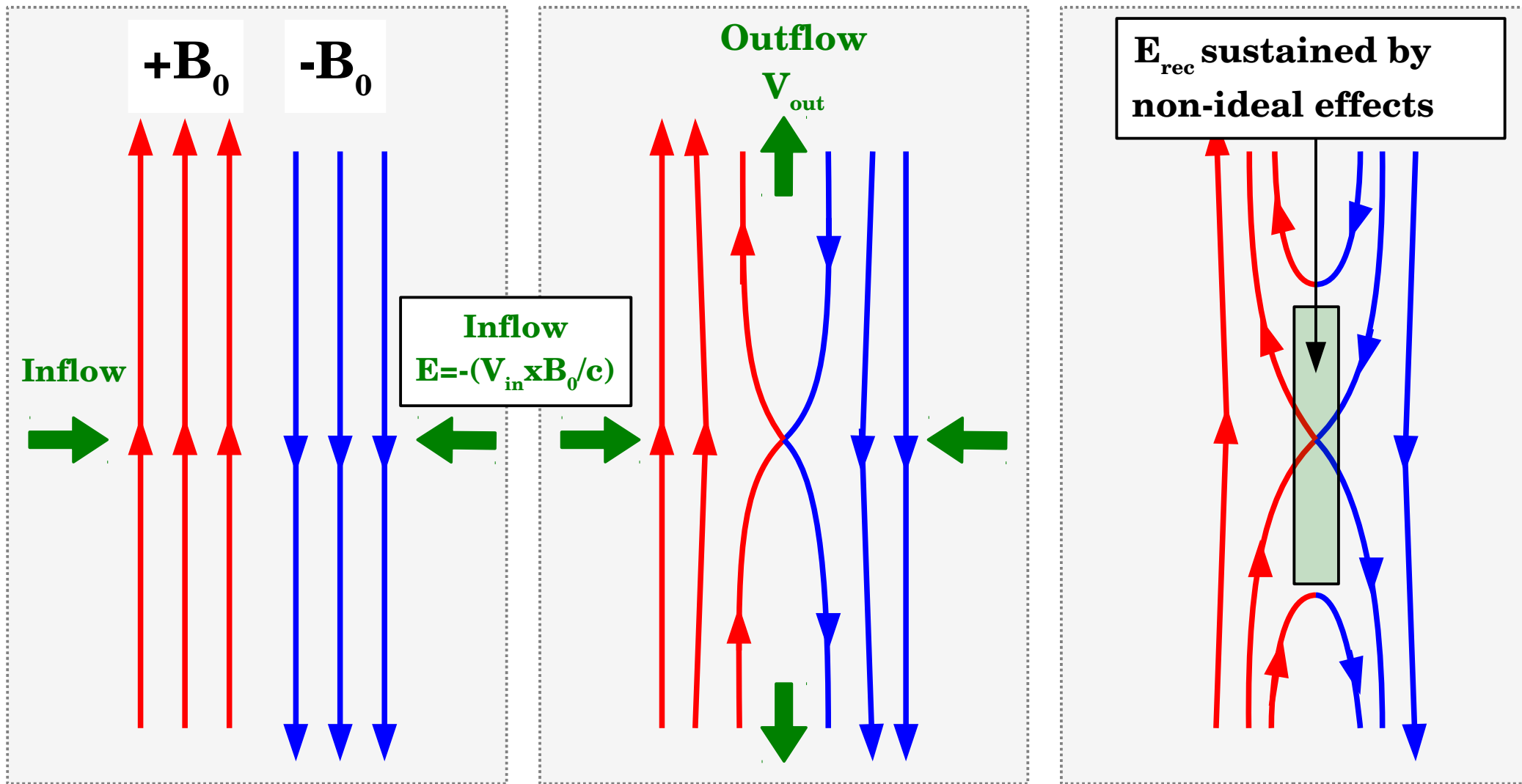
Hard particle spectra
 $p < 2$



I. Relativistic reconnection

Collisionless magnetic reconnection

[See reviews by Zweibel & Yamada, 2009, Yamada 2010, Kagan et al. 2015]



Reconnection rate: $\beta_{rec} = V_{in} / V_{out}$

Magnetic energy \Rightarrow Plasma kinetic energy (heating+non-thermal particles)

Relativistic collisionless reconnection

Upstream plasma magnetization parameter :

$$\sigma = \frac{B_0^2}{4\pi(nmc^2 + P)}$$

If $\sigma \gg 1$, magnetic energy density **exceeds** the rest mass energy density (and pressure) of the plasma. [*Blackman & Field 1994 ; Lyutikov & Uzdensky 2003 ; Lyubarsky 2005*]

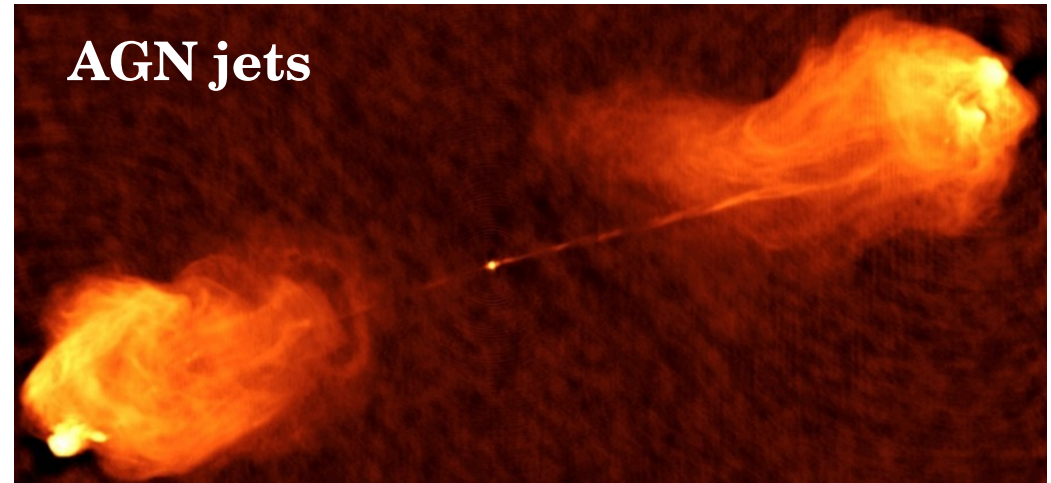
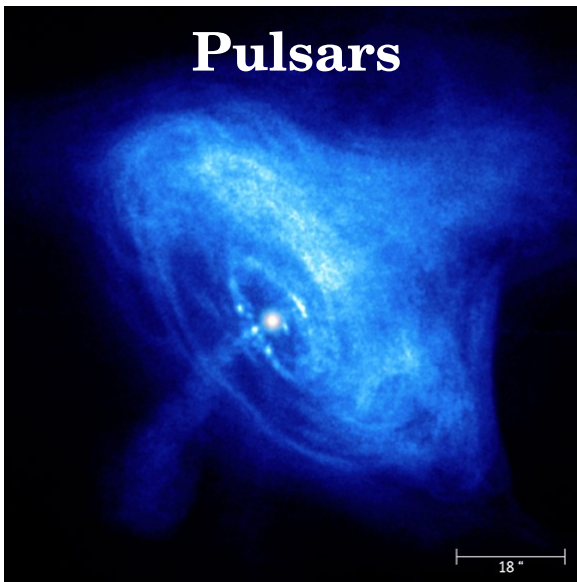
The Alfvén speed approaches the speed of light : $V_A = c \sqrt{\frac{\sigma}{1+\sigma}} \approx c$

Dissipation of magnetic energy => relativistic particles !

Focus on :

- Relativistic **electron/positron pair** plasmas : no mass ratio effects, i.e., no Hall effect. Also relevant to **ultra-relativistic electron/ion plasmas**
- **Collisionless plasmas** : Mean free path \gg system size
- **Optically thin**

Relevant astrophysical environments



Questions :

- **How fast is reconnection ?**
- **How is the magnetic energy transferred to particles ?**
- **How are the particles accelerated and what is the spectrum ?**
- **What are the radiative signatures ? Photon spectra, lightcurves...**
- **Where does reconnection operate in astrophysical systems ?**

Reconnection dynamics

The Particle-In-Cell (PIC) approach

Computational domain

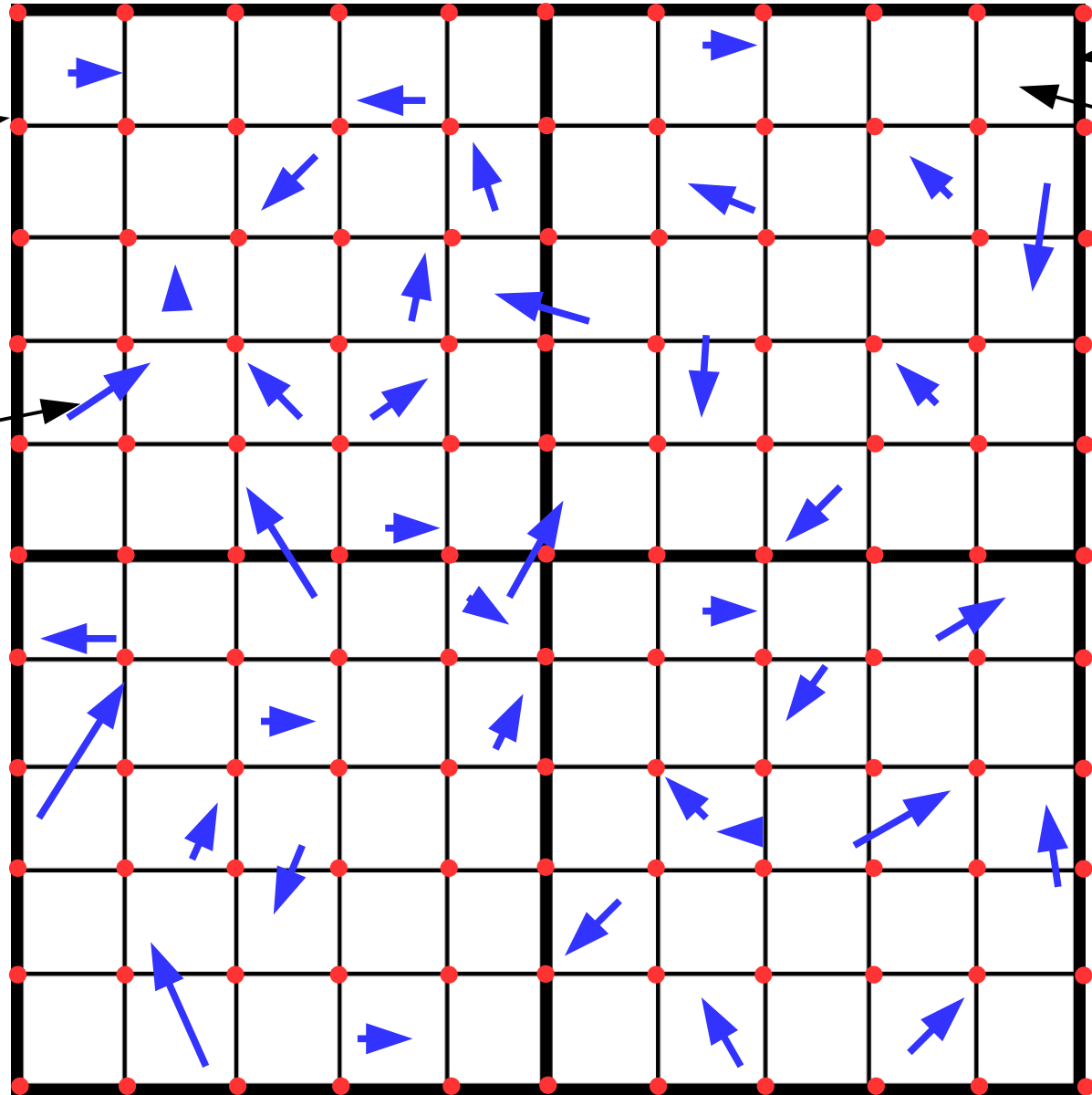
y

Grid

Cell

(E,B) fields known
on the grid
(Eulerian approach)

Particles evolve in
continuous space
(Lagrangian approach)



x

Computation procedure per timestep in PIC

Step 1

Solve Newton's equation

$$\frac{d\mathbf{p}}{dt} = q \left(\mathbf{E} + \frac{\mathbf{v} \times \mathbf{B}}{c} \right)$$

$$\frac{\partial \mathbf{B}}{\partial t} = -c \nabla \times \mathbf{E}$$

$$\frac{\partial \mathbf{E}}{\partial t} = c \nabla \times \mathbf{B} - 4\pi \mathbf{J}$$

$\Delta \mathbf{t}$

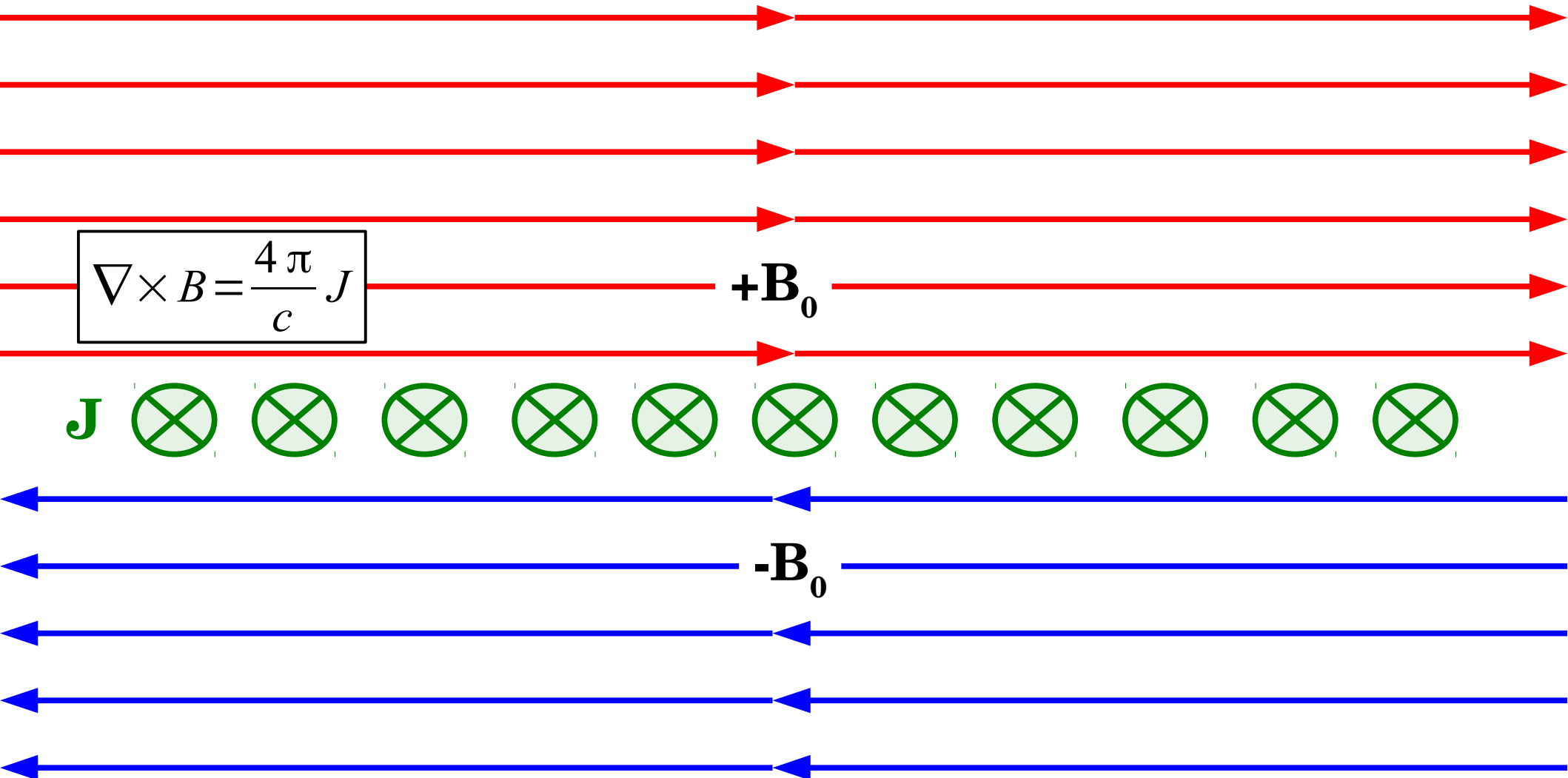
Solve Maxwell's equations (E,B)

Step 3

Deposit Charge and current densities (ρ, \mathbf{J})

Step 2

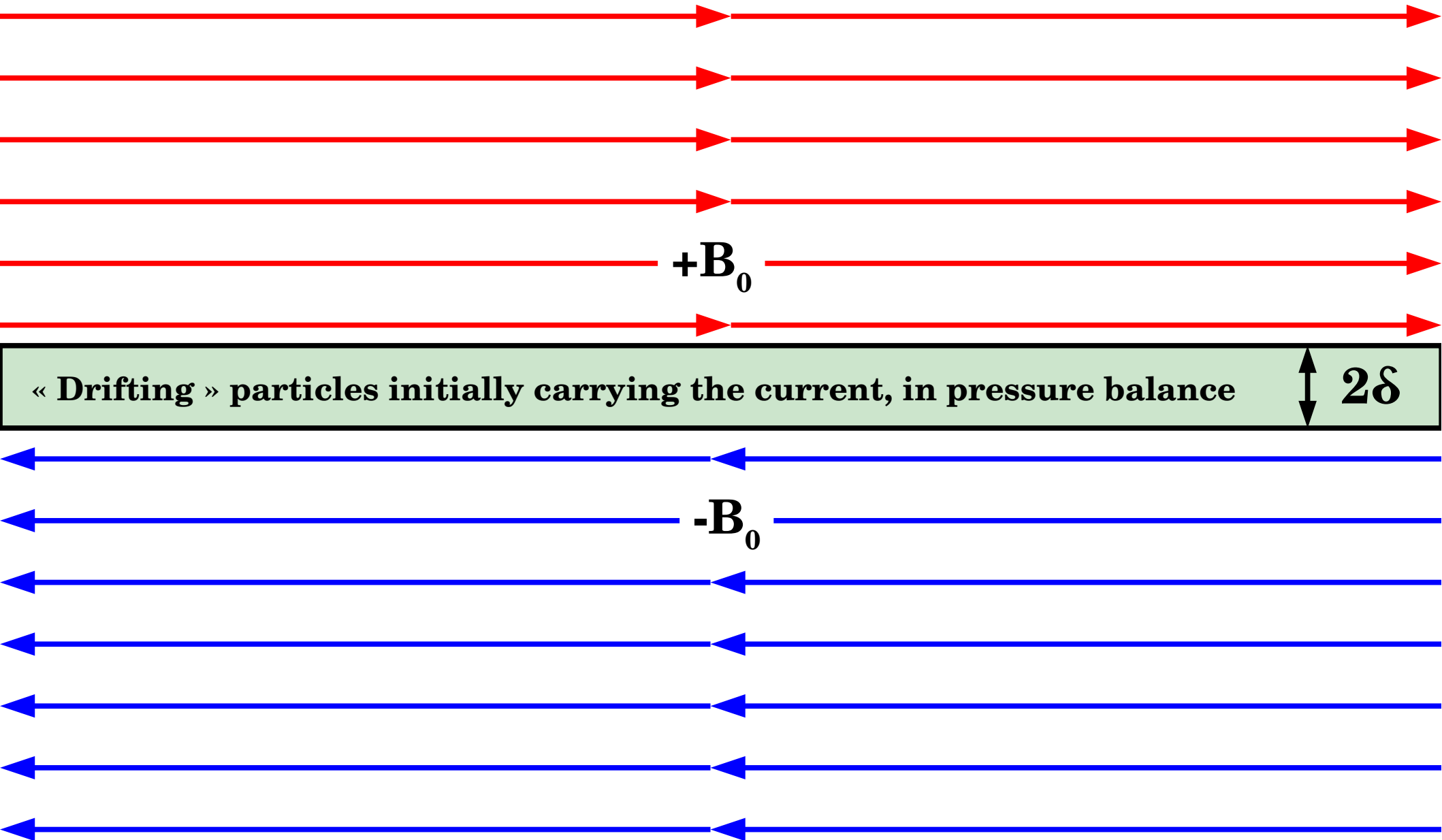
Numerical setup : Harris sheet



Other setup investigated :

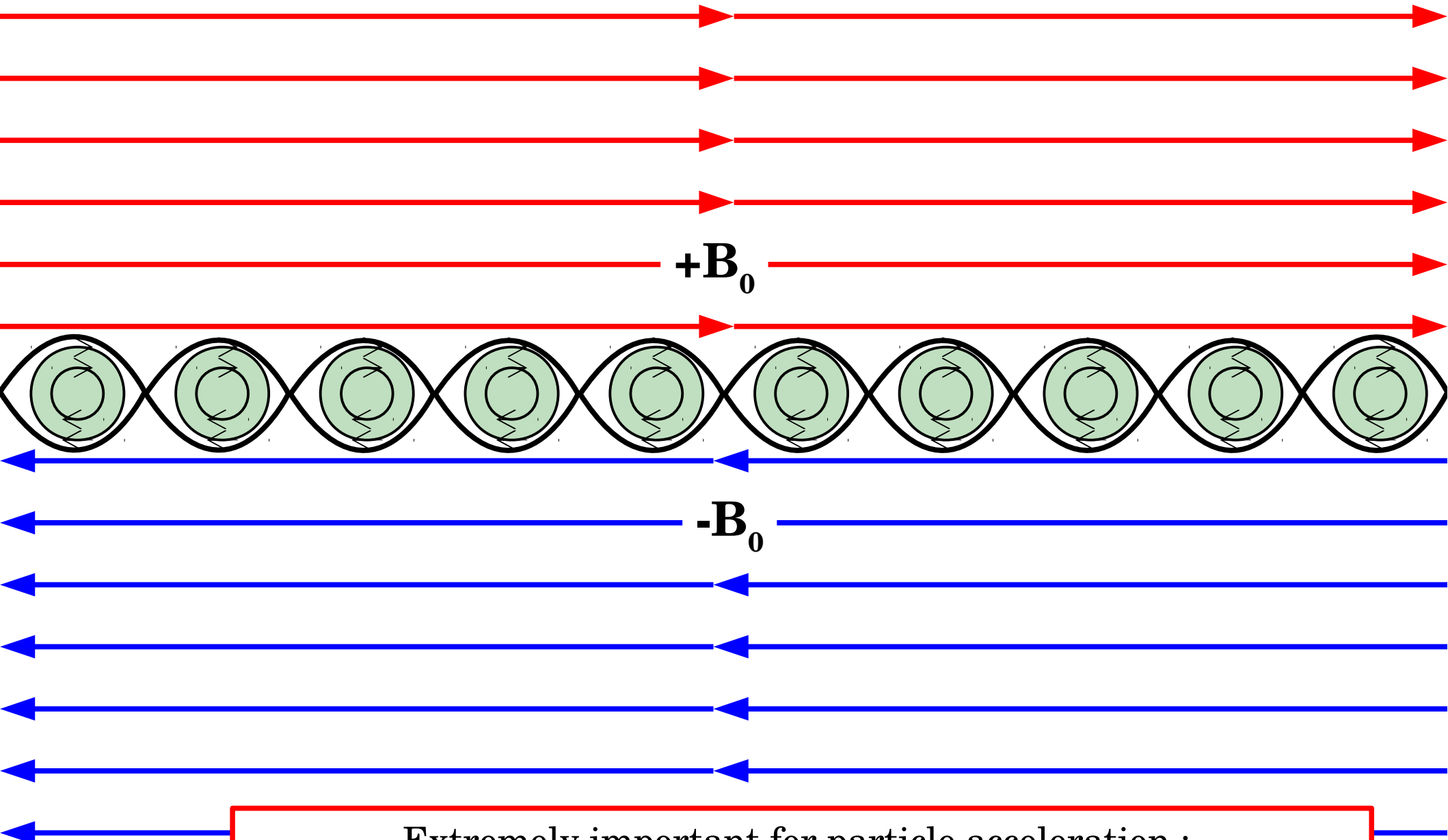
- « **Force-Free** » current sheet (e.g., *Guo et al. 2015*)
- « **ABC** » equilibrium setup **see Journal club** (*Nalewajko+2016 ; Lyutikov+2016*)

Numerical setup : Harris sheet



A long thin current sheet is tearing unstable

[Zelenyi & Krasnoselskikh 1979 ; Zenitani & Hoshino (2007) ; Pétri & Kirk 2007]



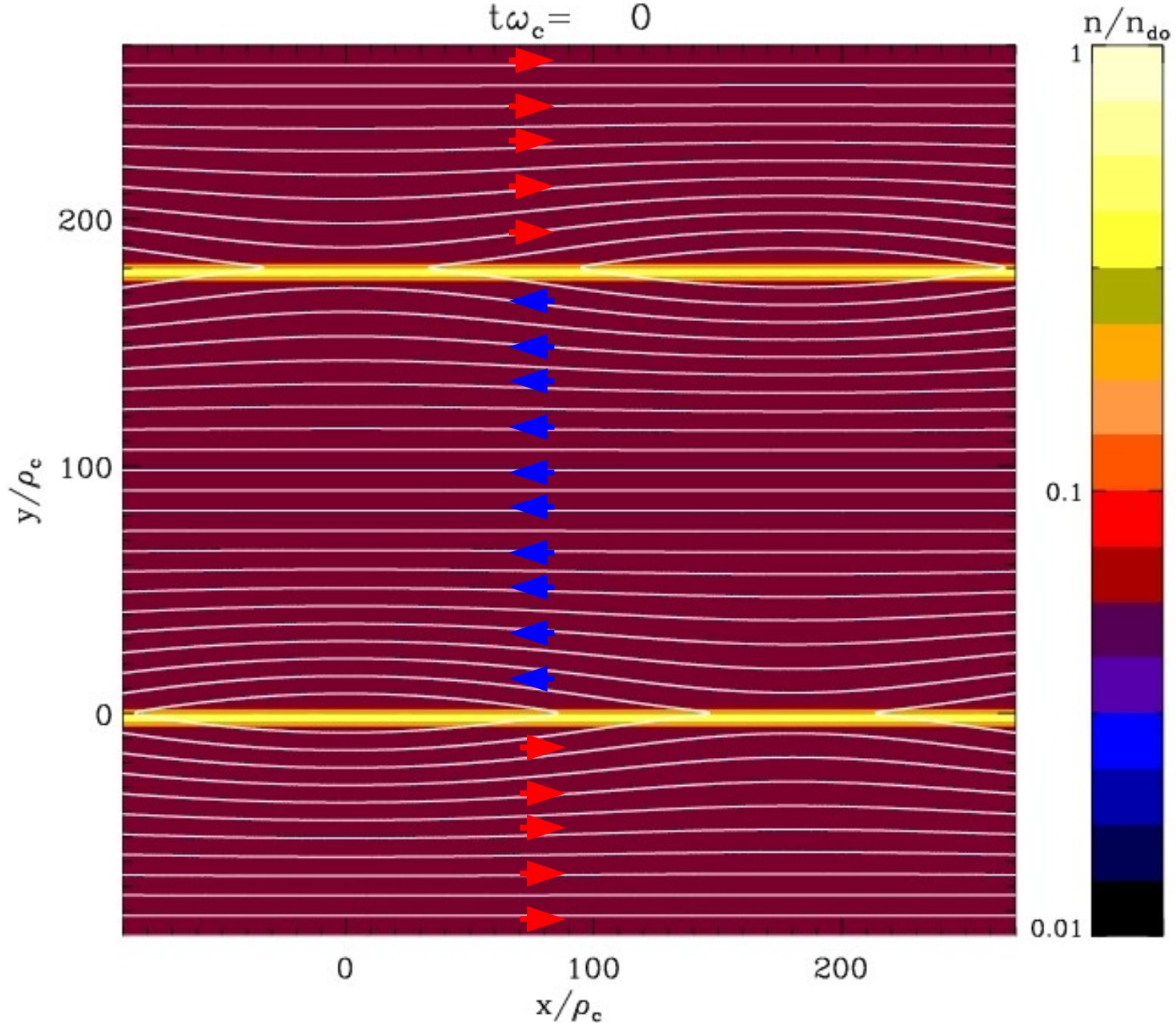
Extremely important for particle acceleration :
Mediate fast reconnection and non-thermal particle acceleration

2D PIC simulation

$t\omega_c = 0$

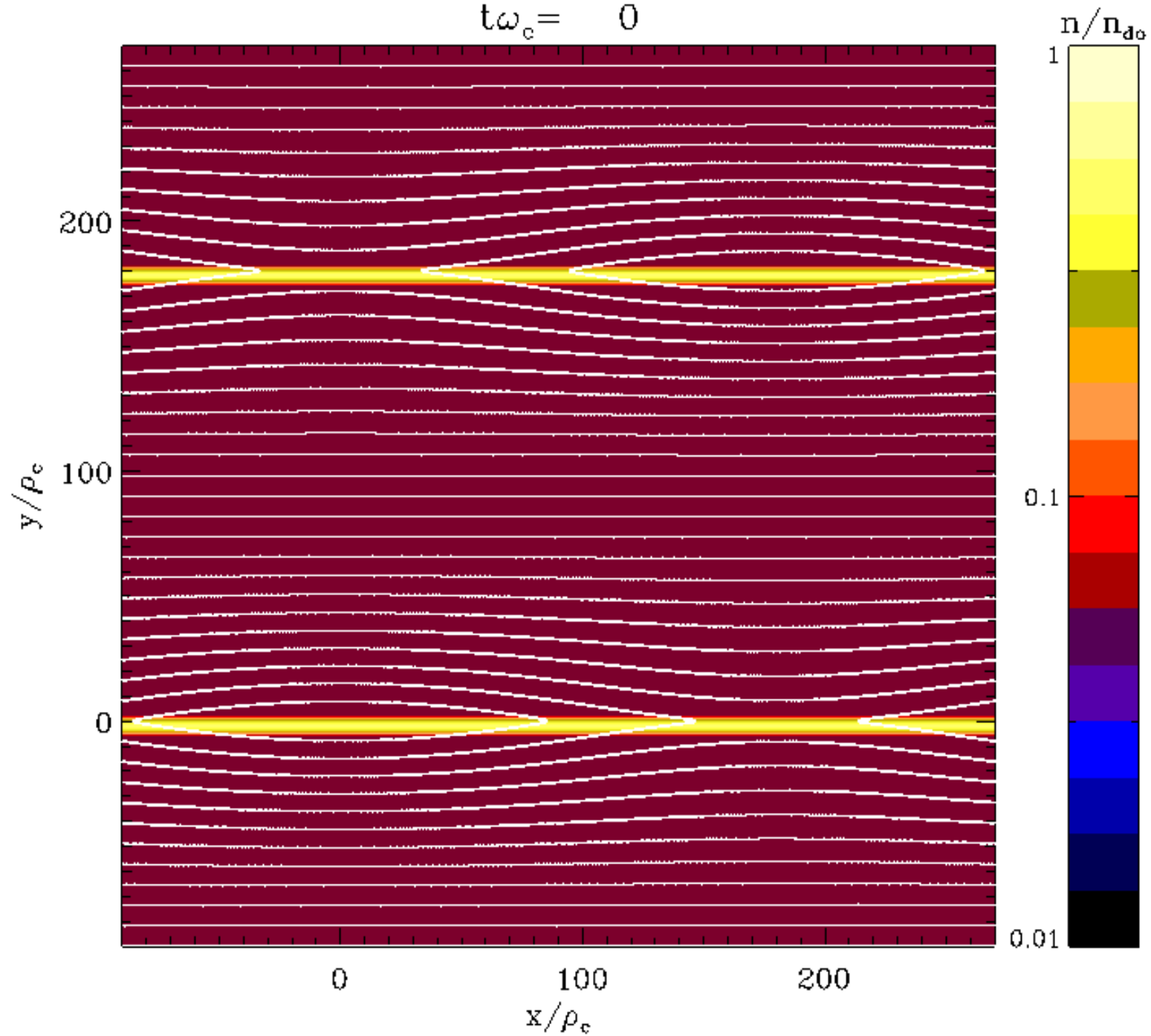
$2\delta \sim$ skin
depth

Periodic

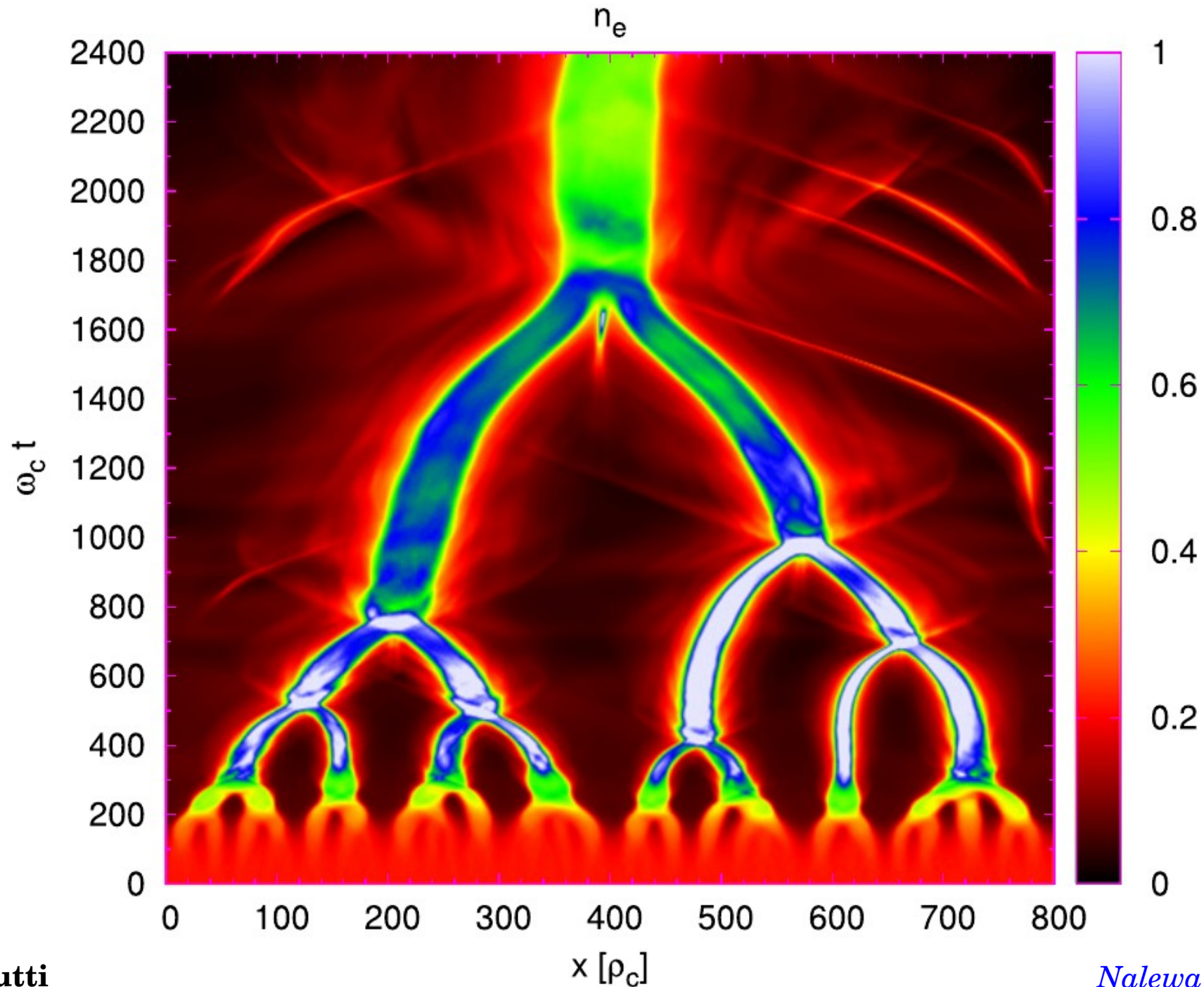


2D PIC simulation

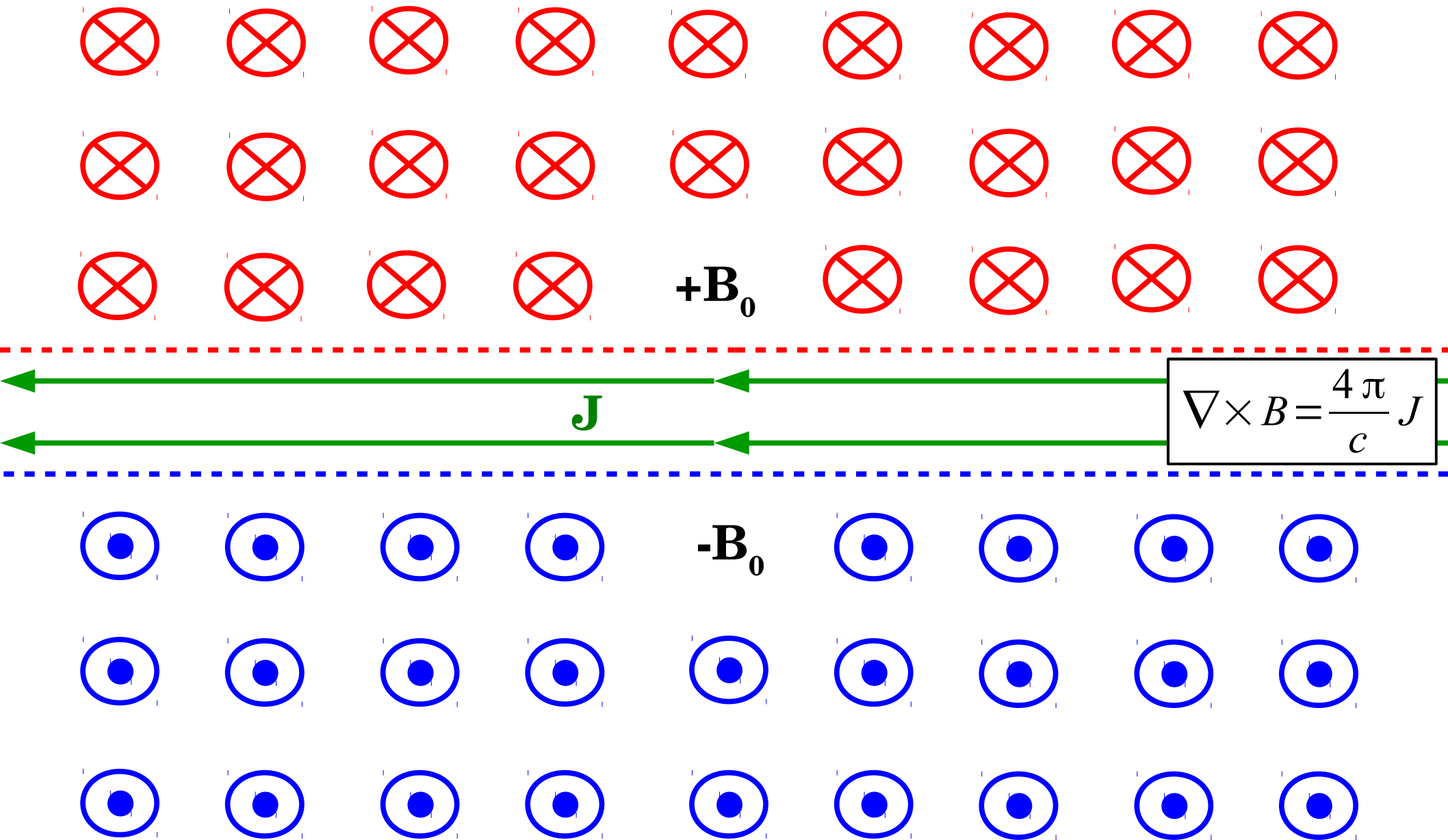
$$t\omega_e = 0$$



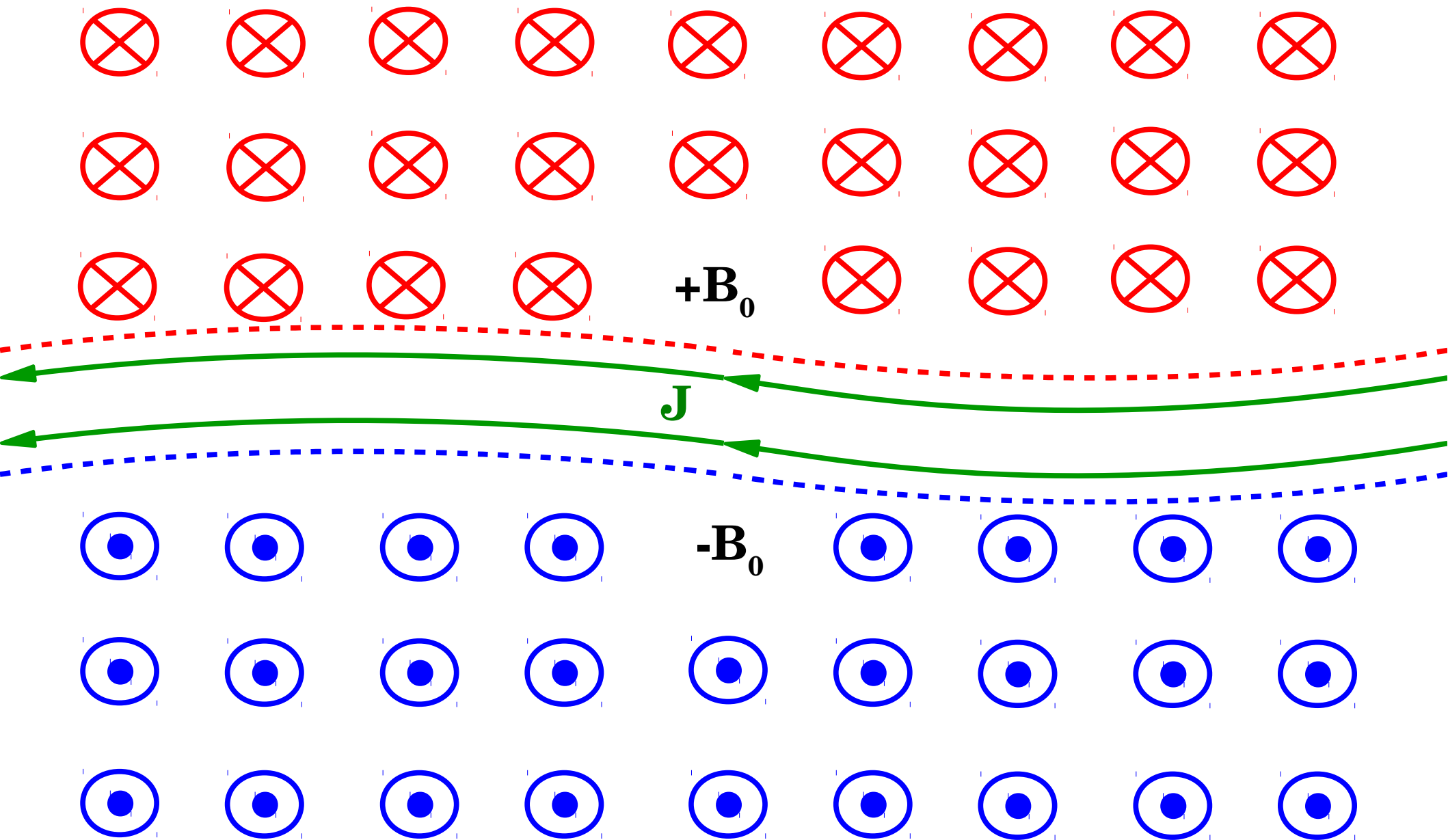
Space-time diagram: Merger tree



The sheet is also kink unstable



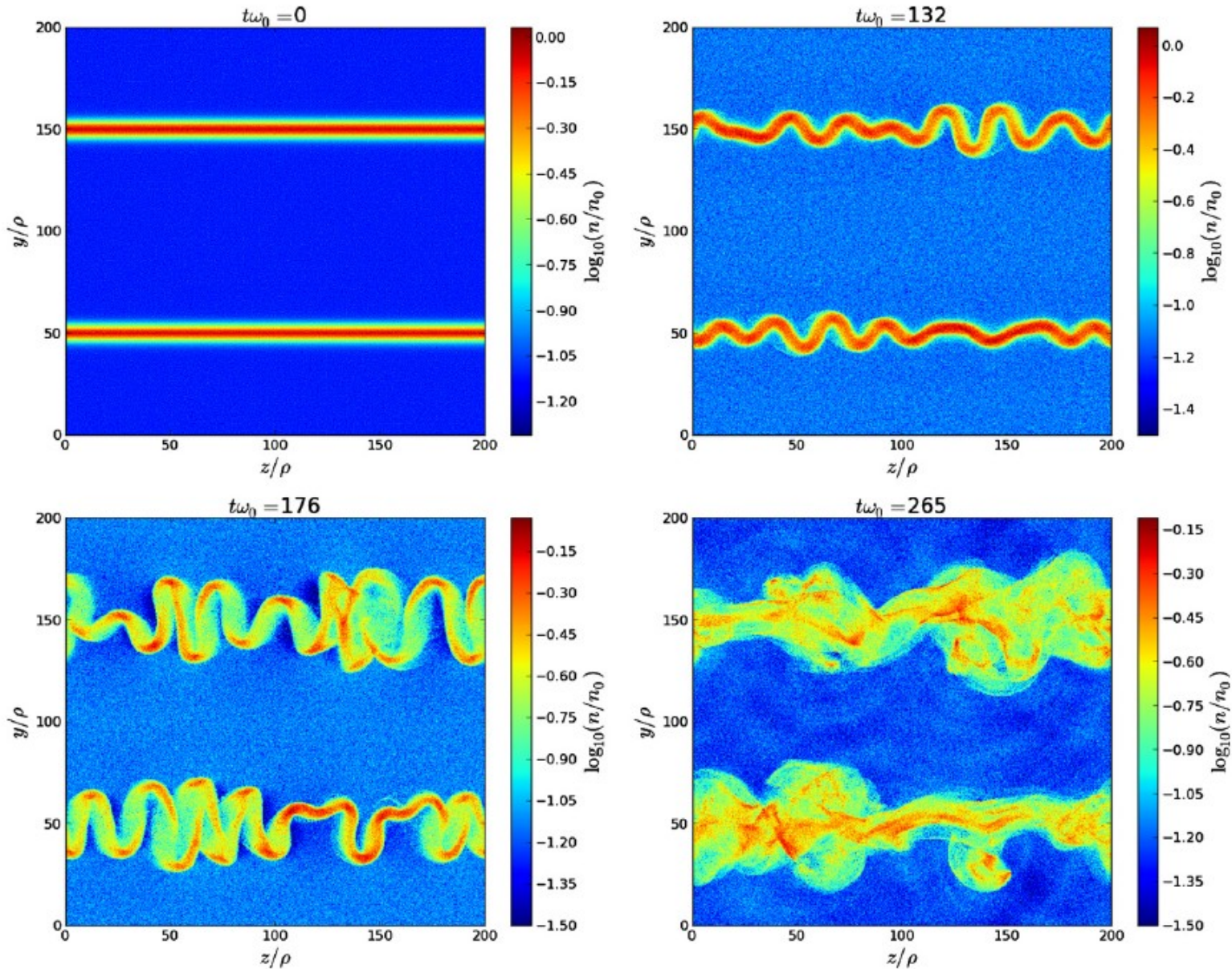
The sheet is also kink unstable



Relativistic drift kink mode

Zenitani & Hoshino 2008 ; Cerutti et al. 2014

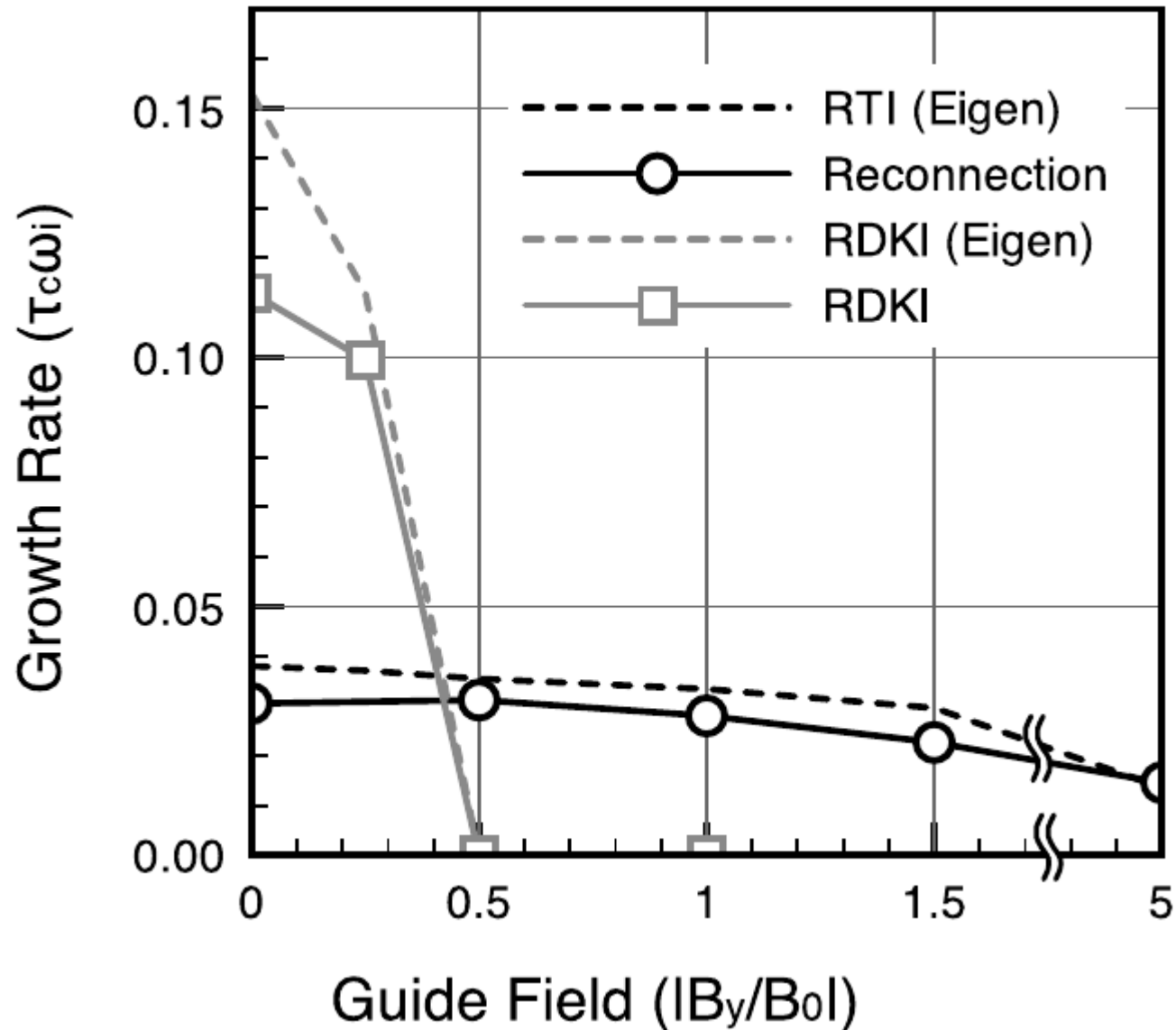
Broadens or even disrupts the layer



B. Cerutti

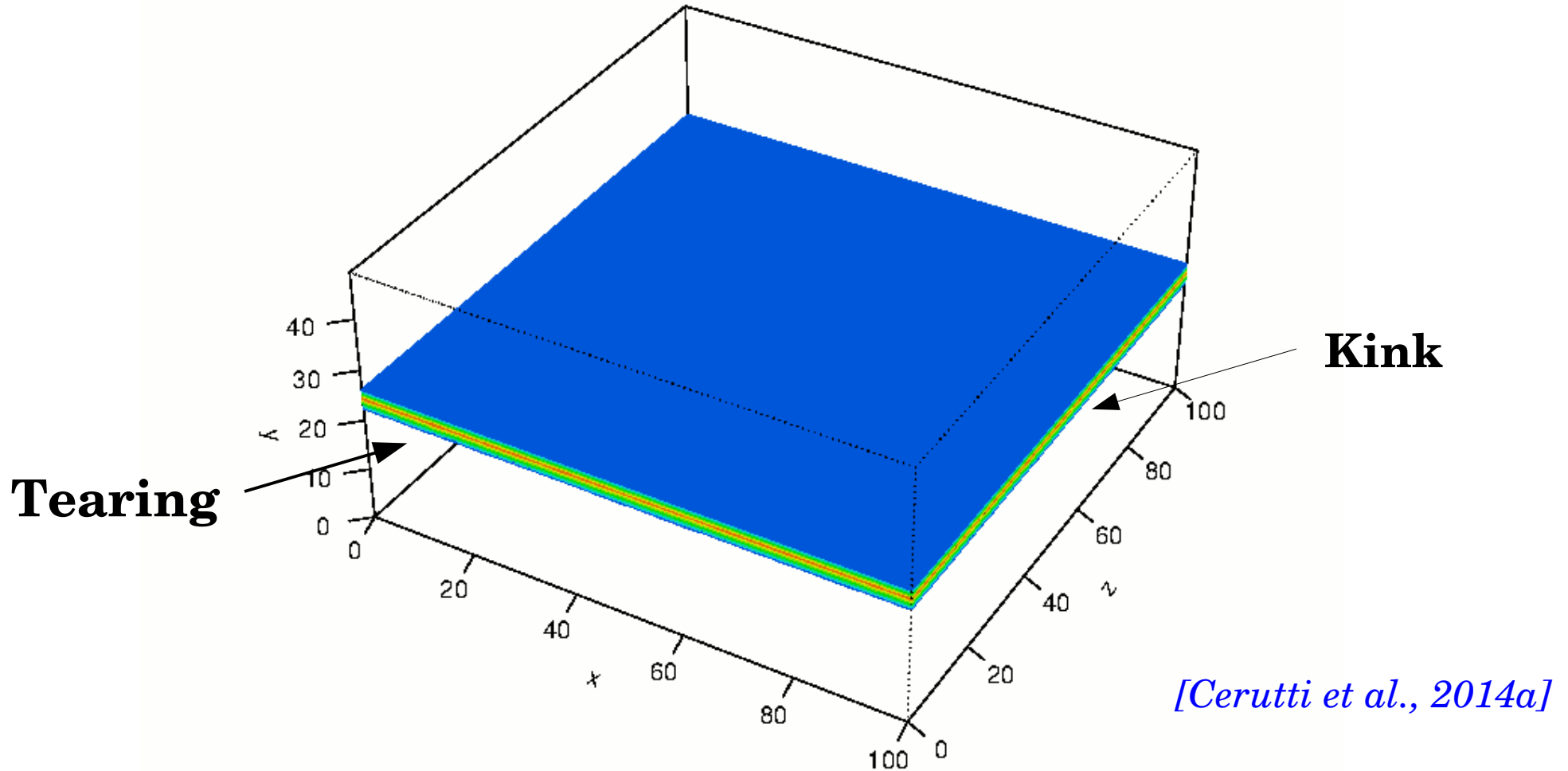
Particle heating, non-thermal **particle acceleration quenched.**

Guide field stabilizes the layer against kink



Zenitani & Hoshino (2008)

3D evolution



The **kink** dominates the early stages of the layer...

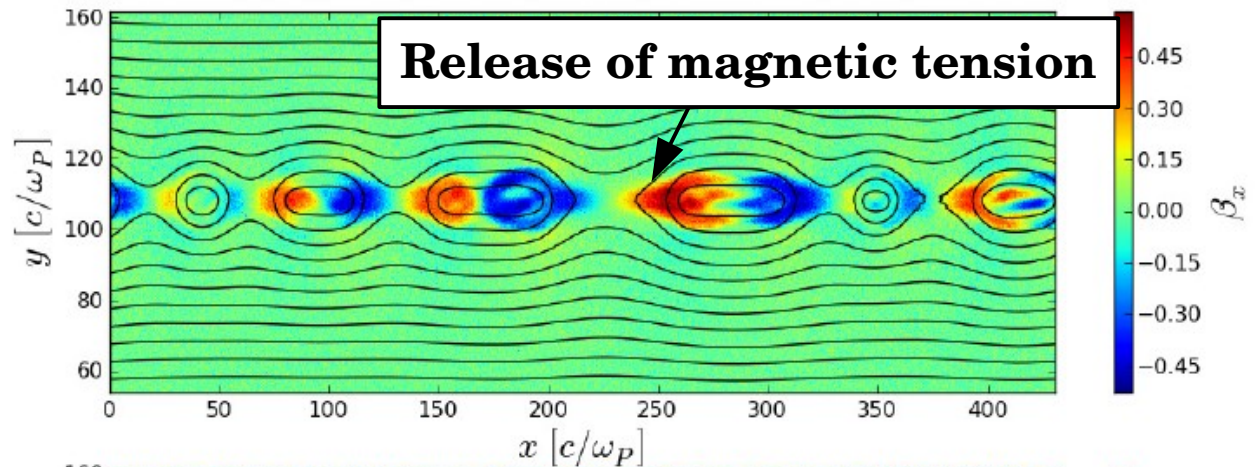
... but the **tearing** controls the later evolution, so that $3D \approx 2D$!

=> Non-thermal particle acceleration still efficient in 3D without guide field

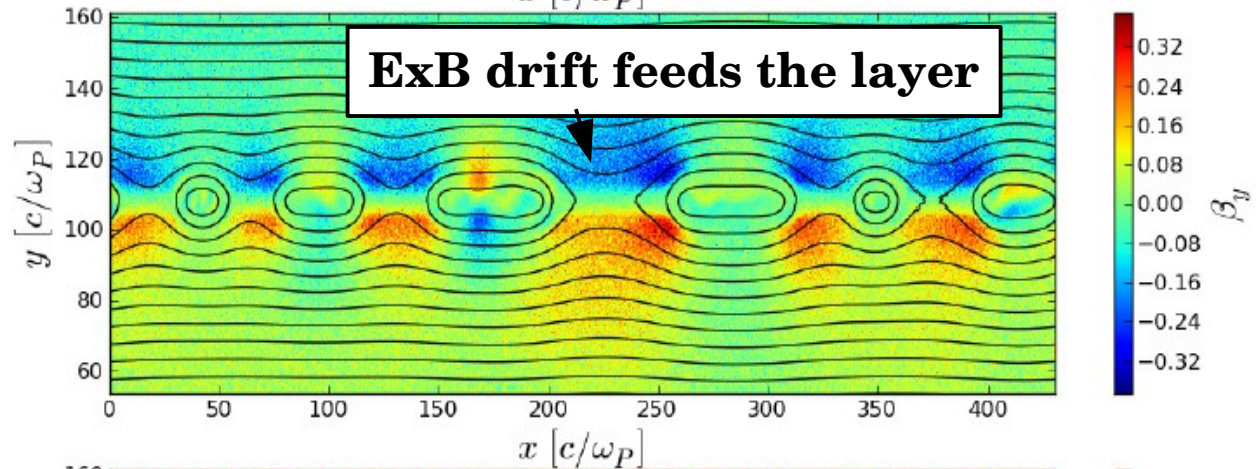
[Sironi & Spitkovsky 2014]

Flow velocity structure

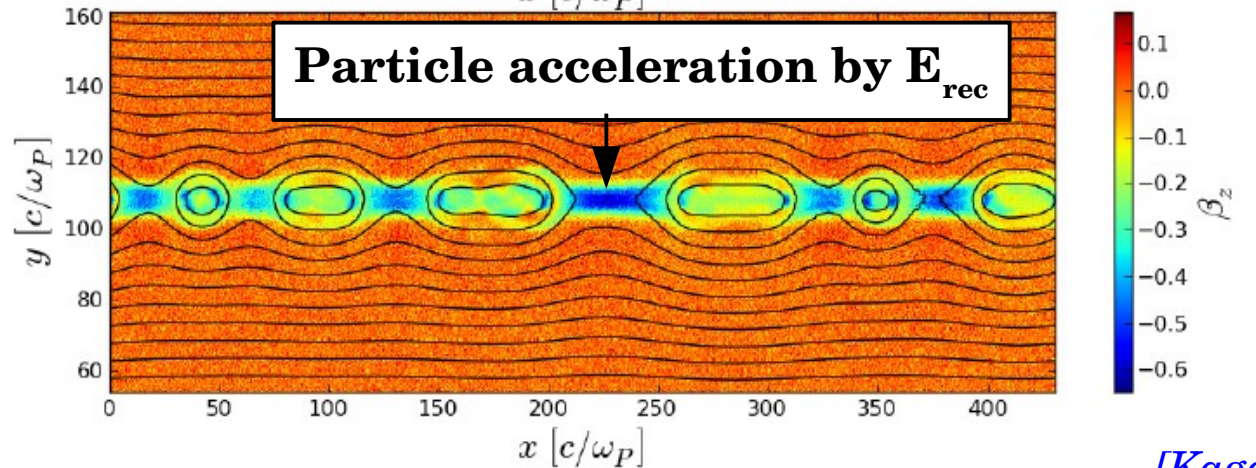
Outflow velocity



Inflow velocity

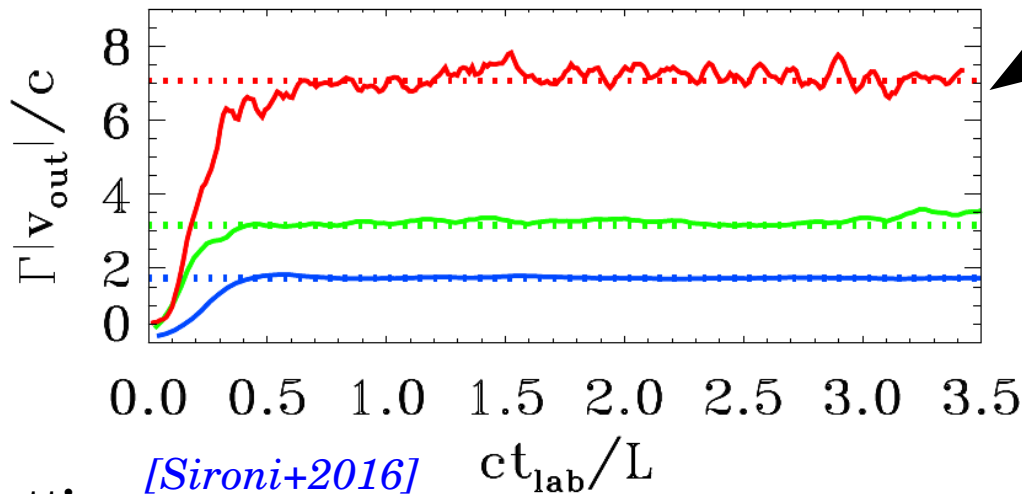
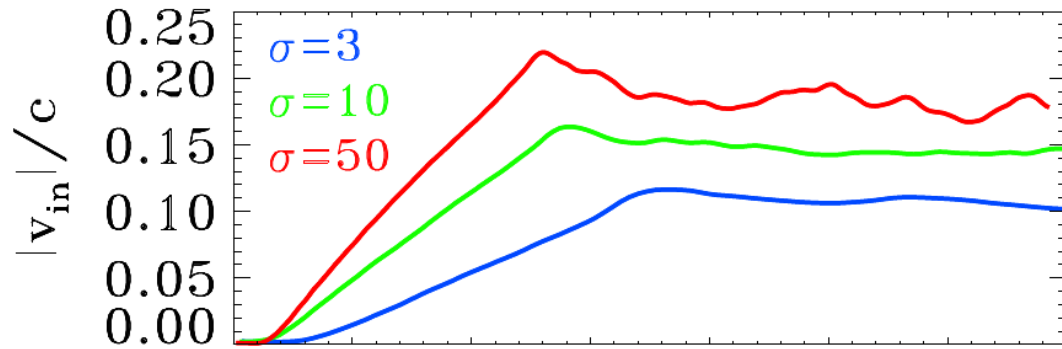
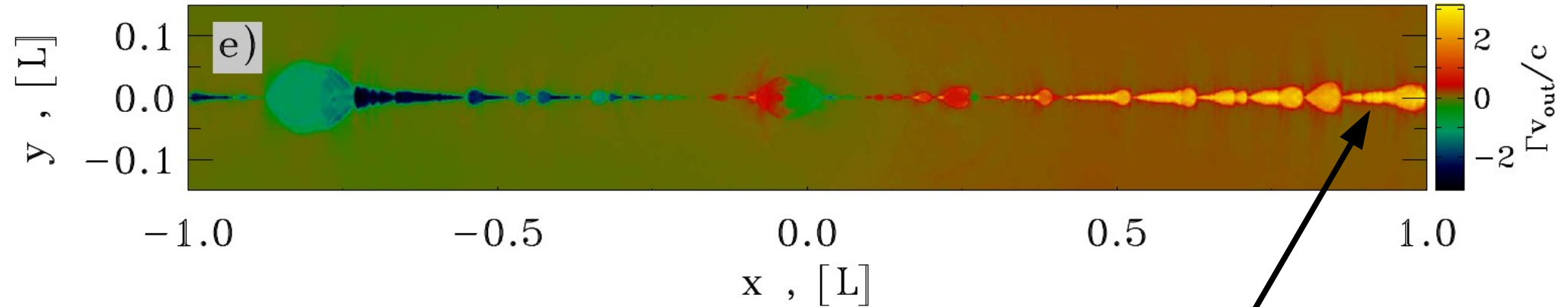


Out-of-plane velocity



Relativistic bulk flows

Large box with open boundaries



Terminal velocity of reconnection outflow is **ultra-relativistic for $\sigma \gg 1$** .

Scales as the Alfvén speed :

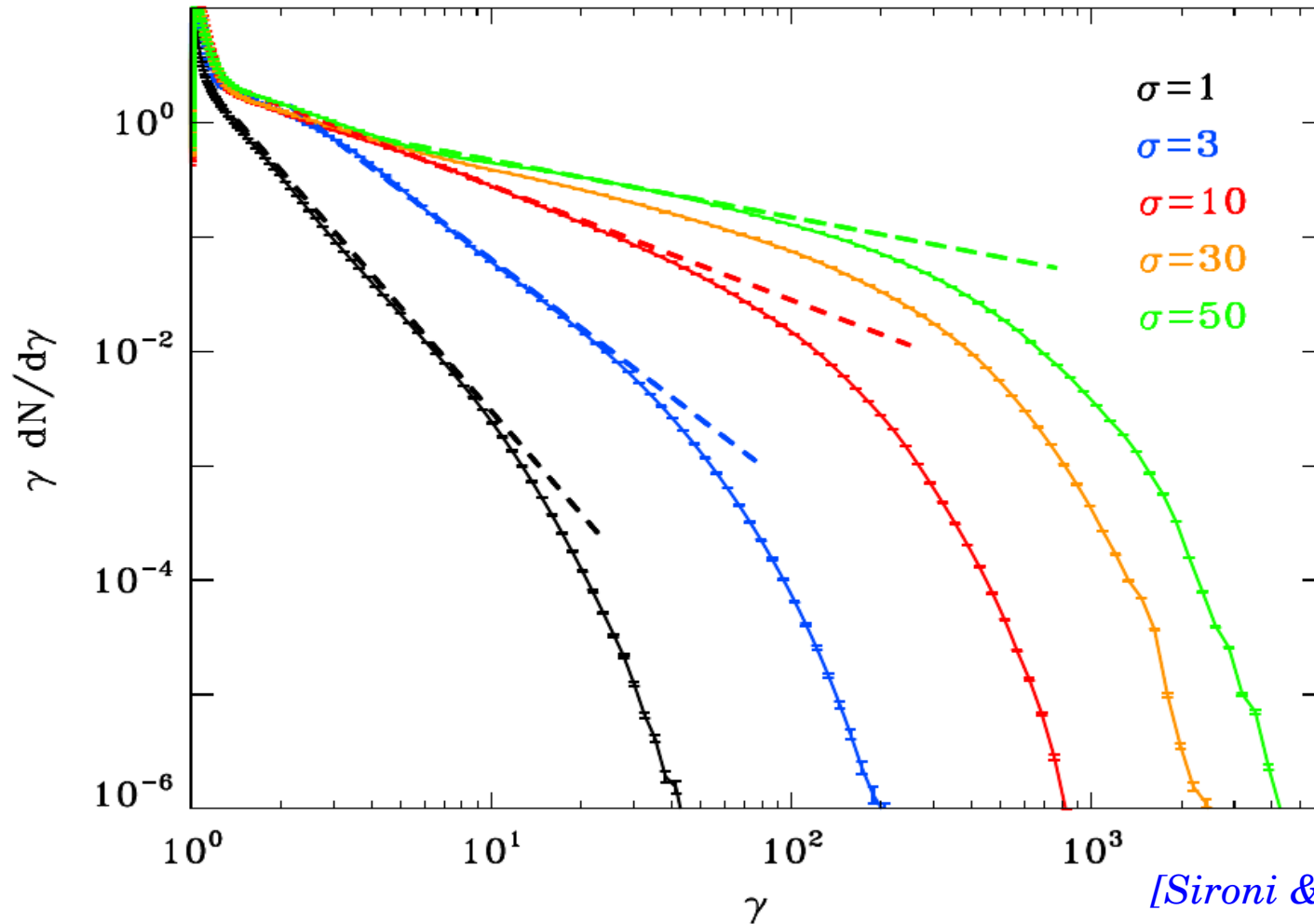
$$\frac{V_A}{c} = \sqrt{\frac{\sigma}{\sigma+1}} \rightarrow \Gamma \approx \sqrt{\sigma}$$

[Lyubarsky 2005]

Particle acceleration

Particle acceleration in relativistic reconnection

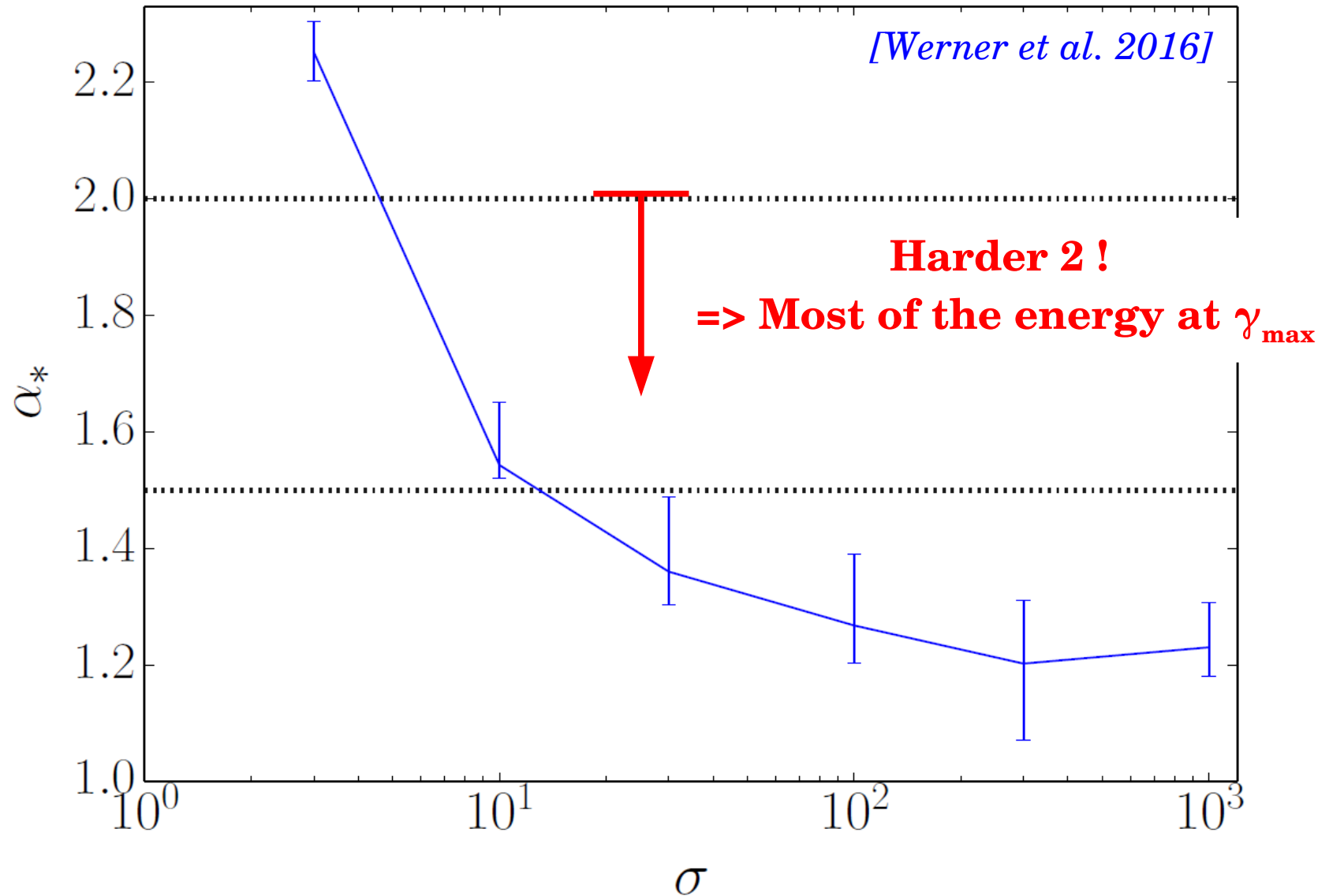
[Zenitani & Hoshino; Jaroschek et al.; Bessho & Bhattacharjee; Pétri & Lyubarsky; Sironi & Spitkovsky; Liu et al., Cerutti et al.; Kagan et al., Guo et al., Werner et al., ...]



[Sironi & Spitkovsky 2014]

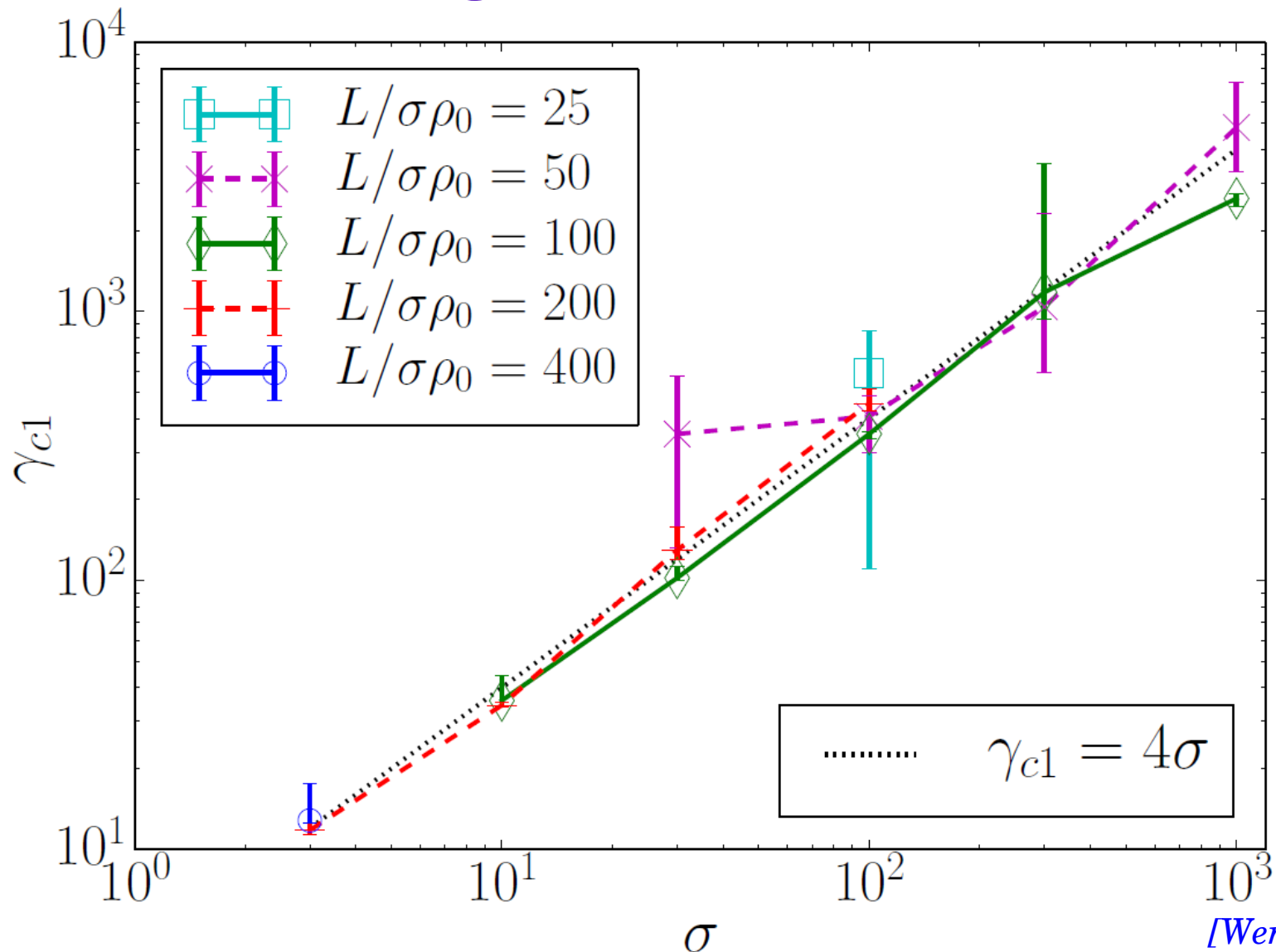
Non-thermal particle acceleration is **efficient for highly magnetized plasmas** ($\sigma \gg 1$),
=> Need for **large** computational box and **late** time evolution of the system.

Spectral hardening with sigma



Spectral index converges towards **$p=1-1.5$** at high magnetization

Large box-size limit

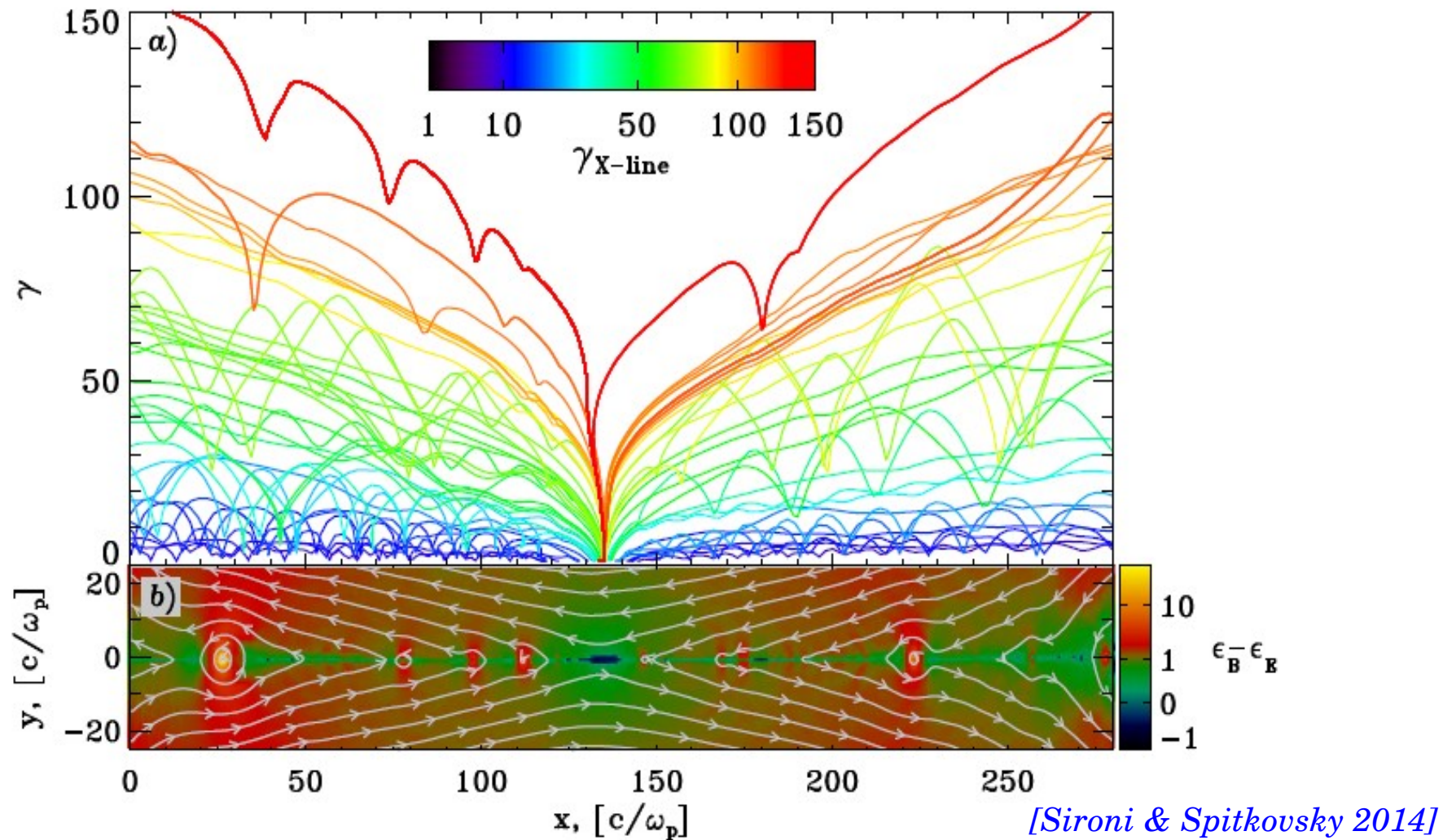


If $L \gg \sigma \rho_0$, the spectrum cuts exponentially with a cut-off given by

$$\gamma_{\max} = 4\sigma$$

Maximum energy limited by the energy budget (σ) !

Acceleration mechanism #1 : X-points



The final particle energy depends how close from the X-point they are injected, leads to a **wide distribution** in energy.

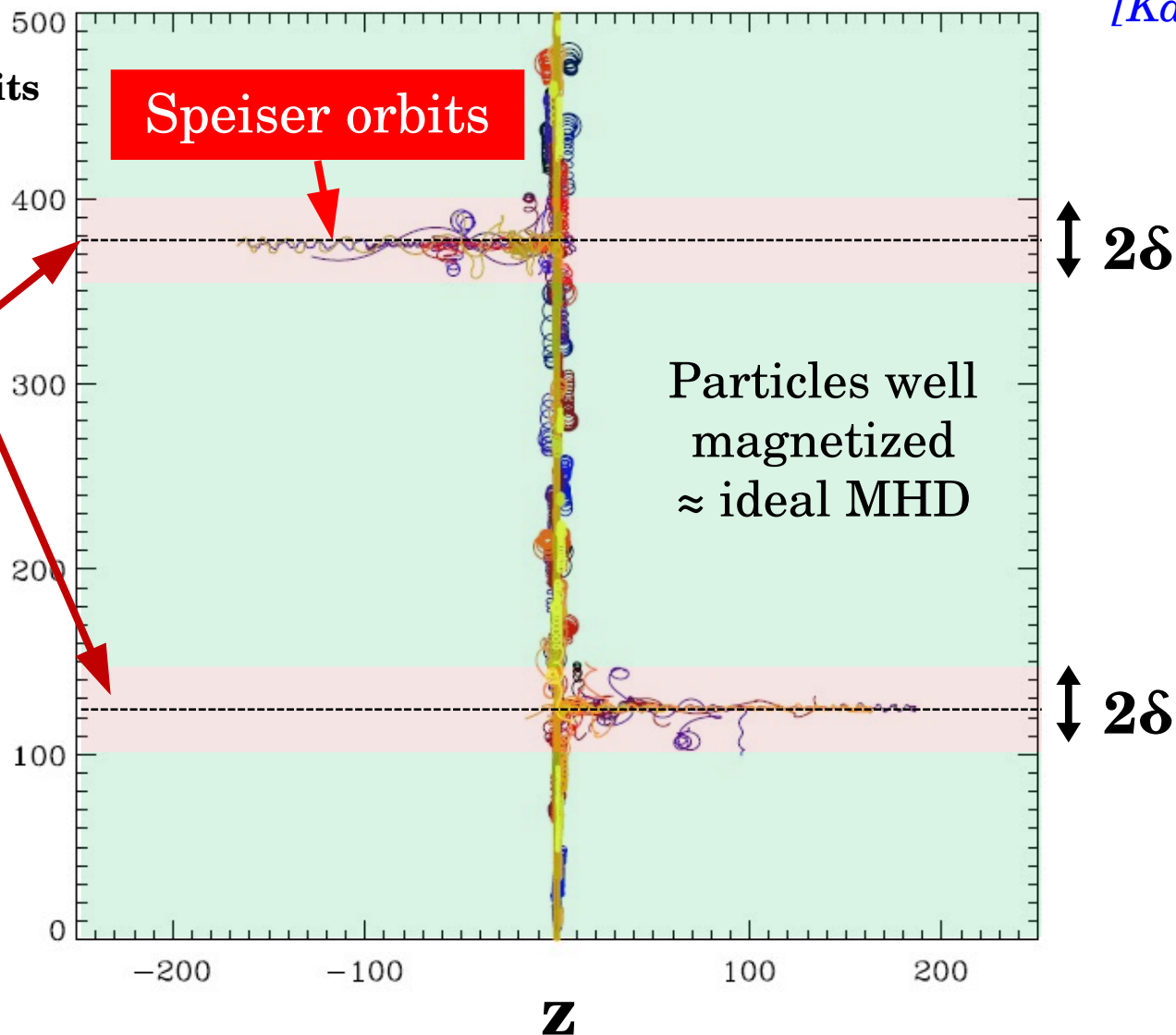
Strong particle acceleration at X-points

[Kagan+2015]

Sample of 150 particle orbits

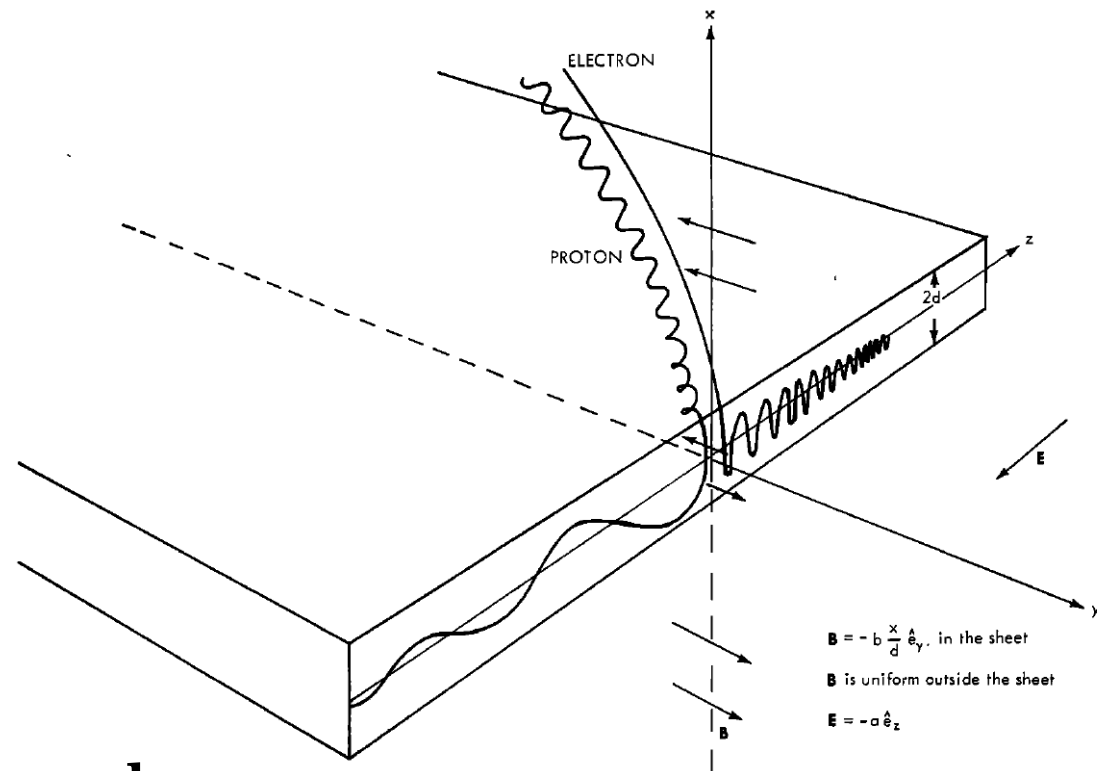
Particles' beam
 $E > B$, non-ideal
MHD!

y



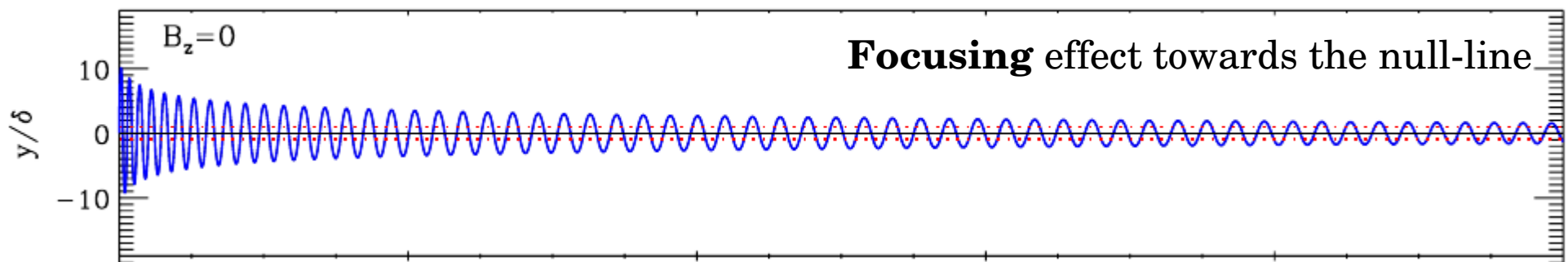
Particle energy **increases linearly with time**
Fast compared with diffusive time $\sqrt{\text{time}}$ as in DSA.

Speiser orbit



[Speiser 1965]

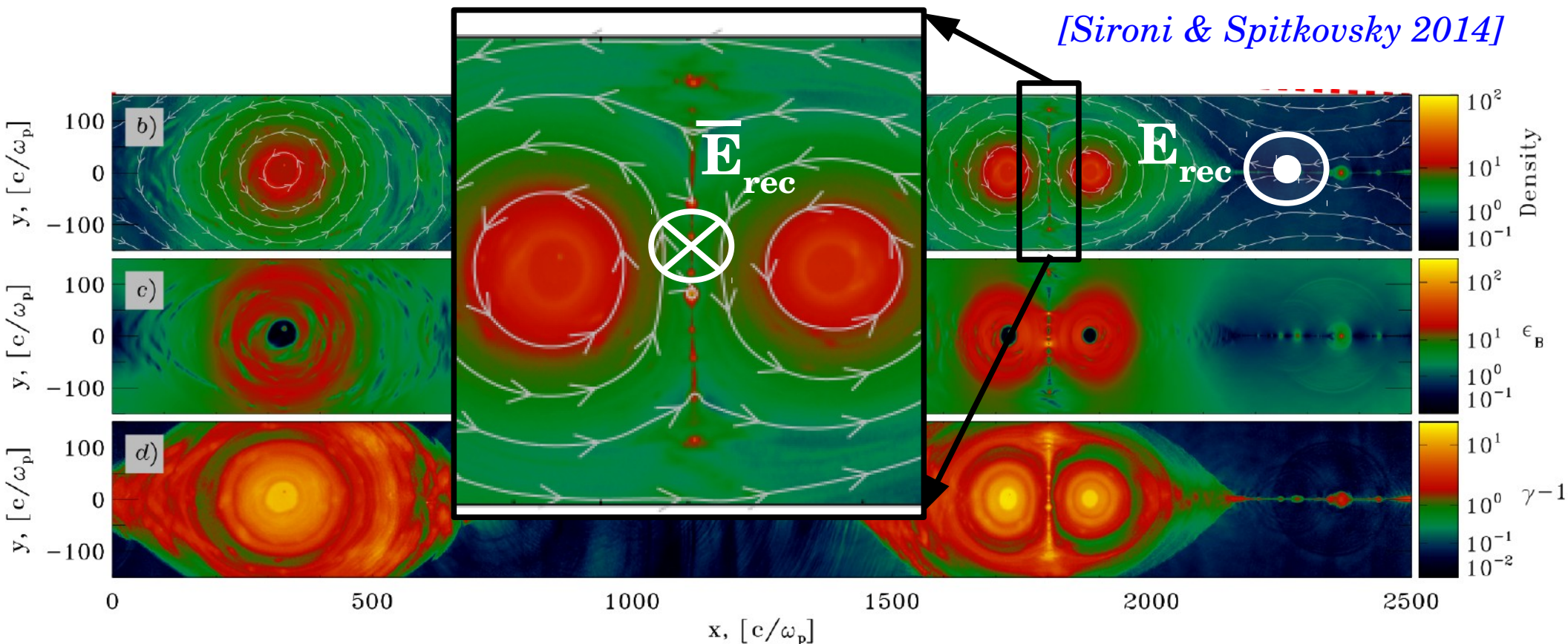
Relativistic analog



ExB drifting of the particle orbit guiding center towards the center of the sheet.

[Uzdensky+2011 ; Cerutti+2012]

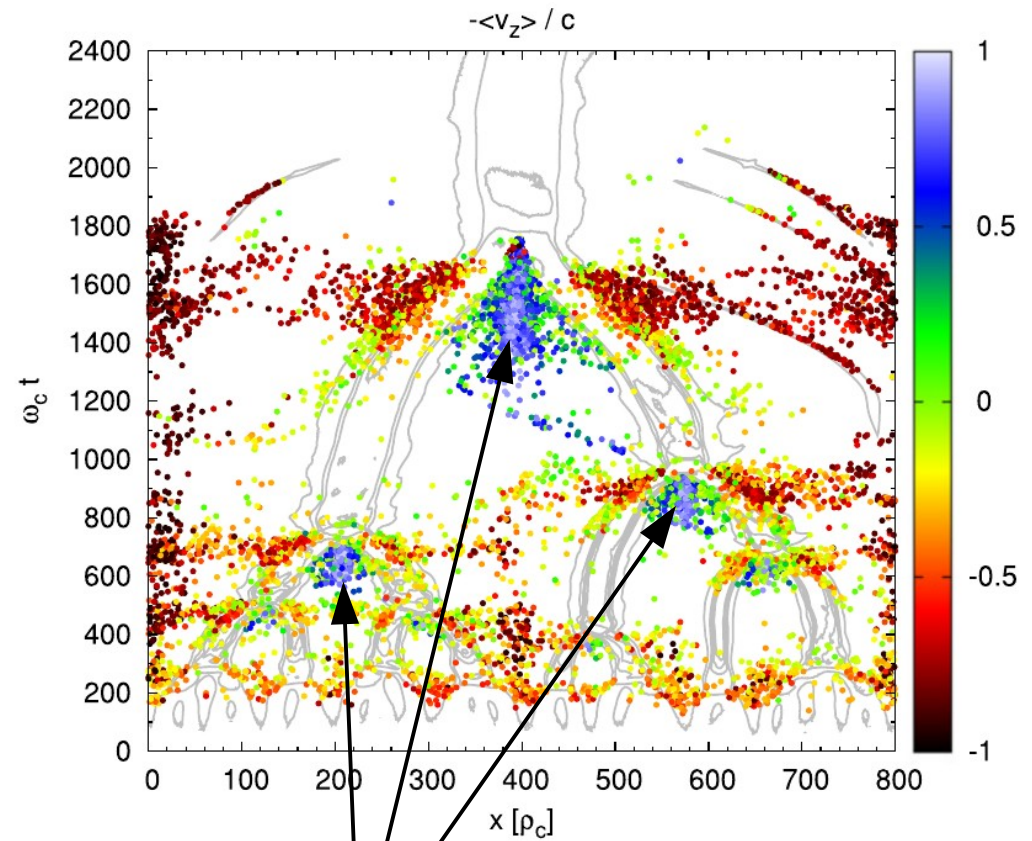
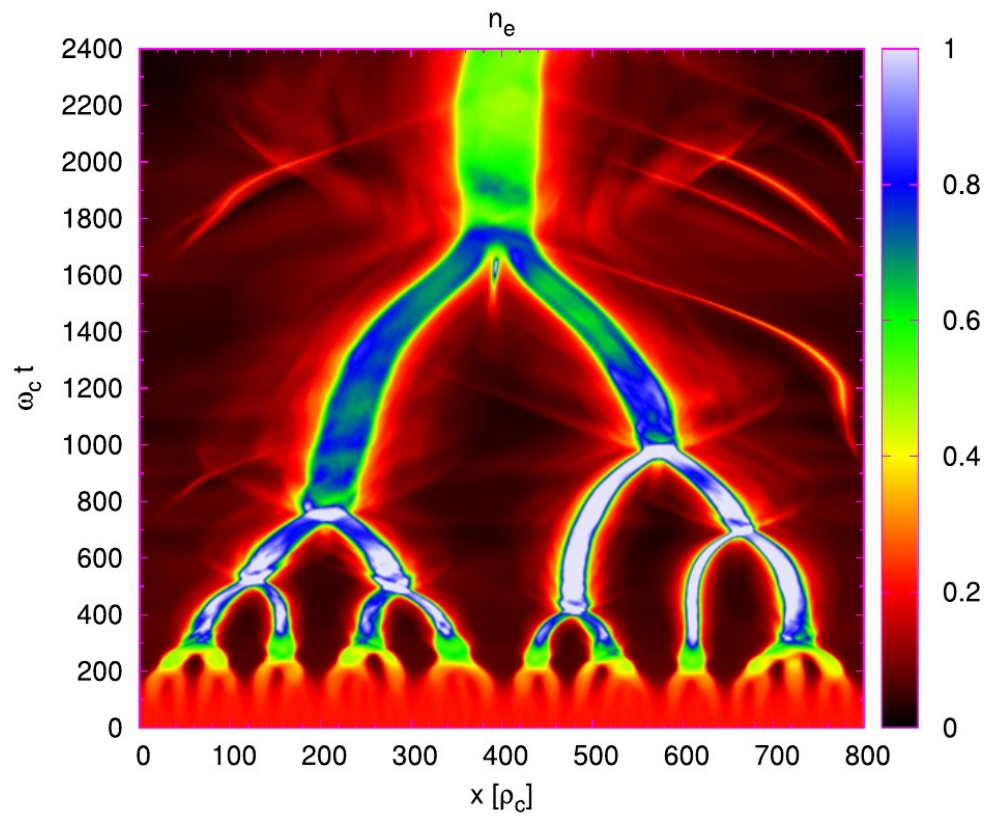
Acceleration mechanism #2 : Plasmoid mergers



[See also Oka et al. 2010 (non relativistic), Nalewajko et al. 2015]

Extra particle acceleration by the **anti-reconnection electric field** between merging islands: **increase the maximum energy!**

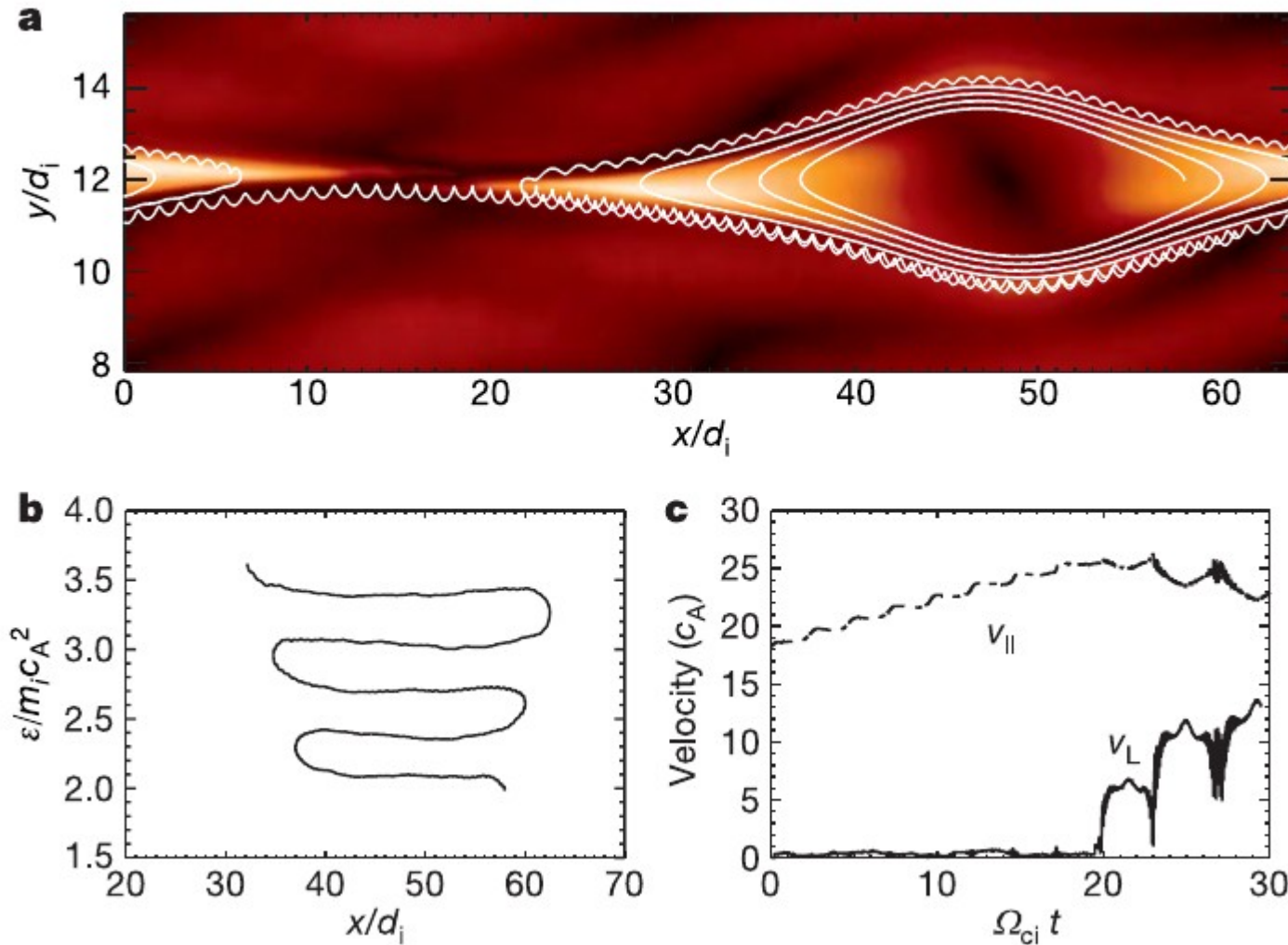
Space-time diagram: Merger tree



Nalewajko+2015

Anti-reconnection sites

Acceleration mechanism #3 : Contracting islands

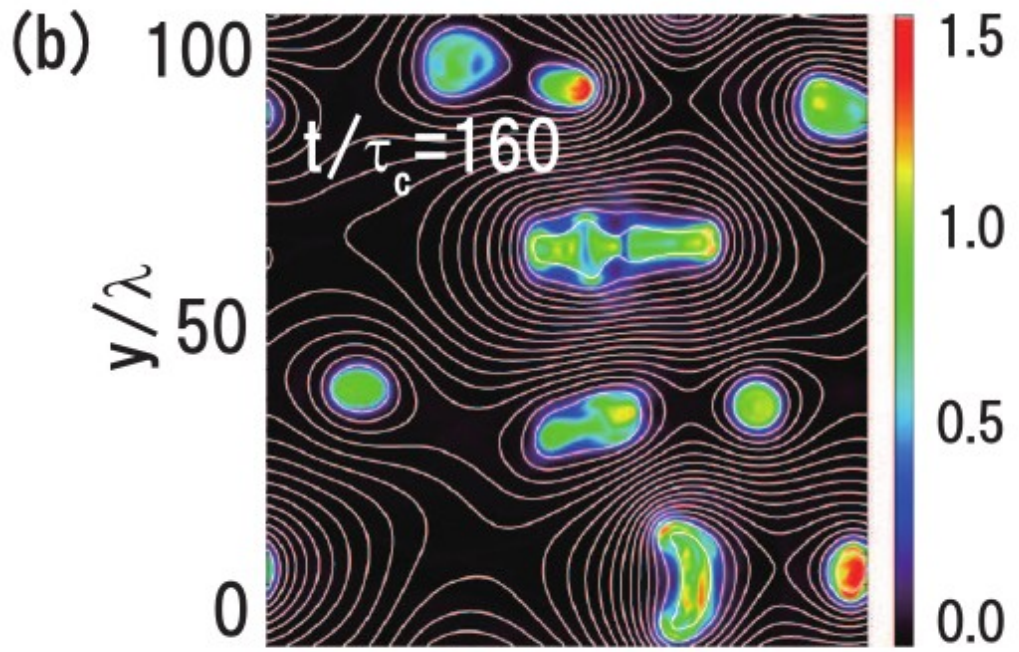
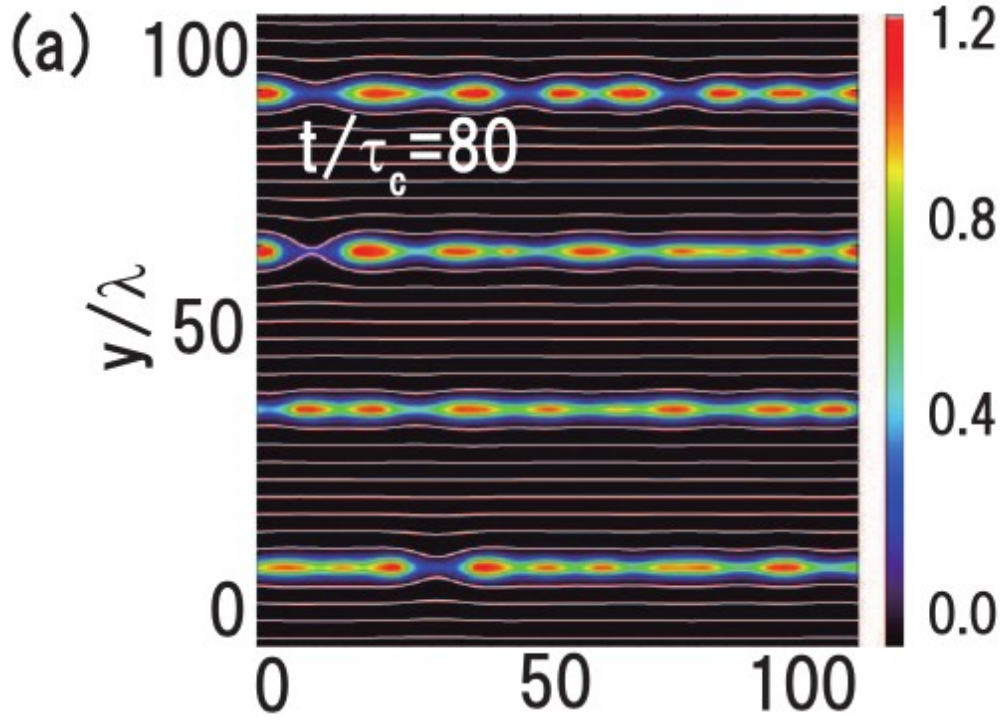


Drake+2006

Fermi-like process, particles bounce back and forth between the island edges as it circularizes.

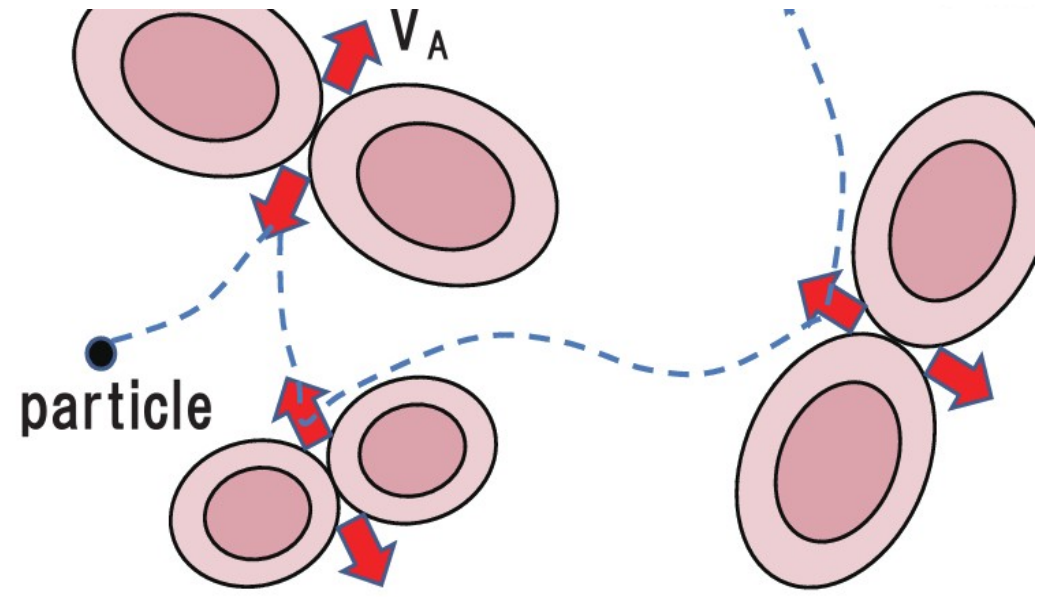
But, negligible in the relativistic regime, too fast ! [*Nalewajko+2015*]

Acceleration mechanism #4 : Multi-island scattering



Stochastic acceleration between plasmoids. Fermi-like acceleration.

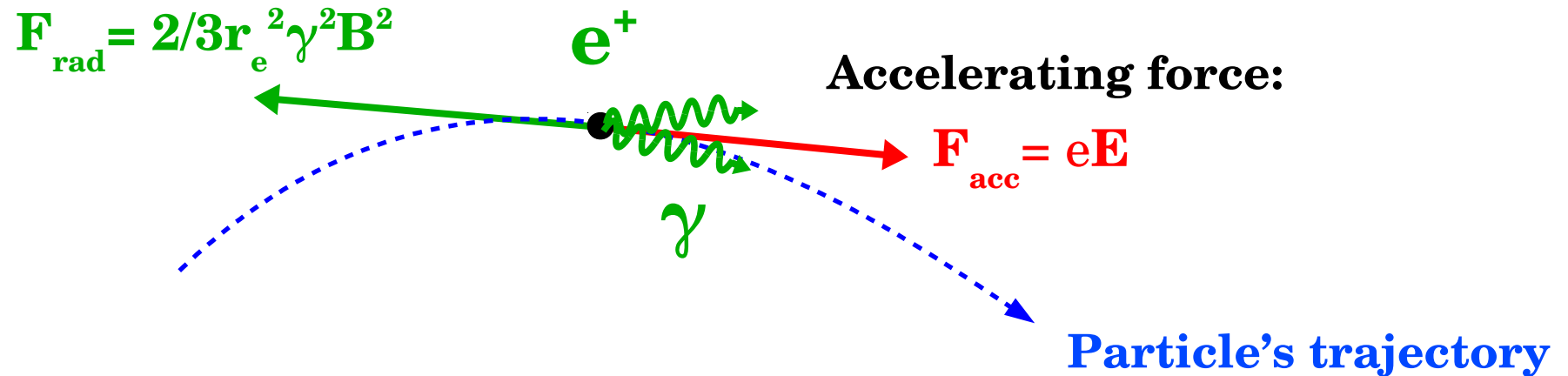
Special conditions needed : Needs a stack of current layers **interacting** with each other.



Hoshino+2012

Radiation-reaction-limited particle energy

Radiation reaction force:



Radiation reaction limit: $\mathbf{F}_{\text{acc}} = \mathbf{F}_{\text{rad}} \Rightarrow \gamma_{\text{rad}}$

Synchrotron photon energy: $\epsilon_{\text{max}} = \frac{3}{2} \gamma_{\text{rad}}^2 \hbar \omega_c = 160 \times (E/B) \text{ MeV}$

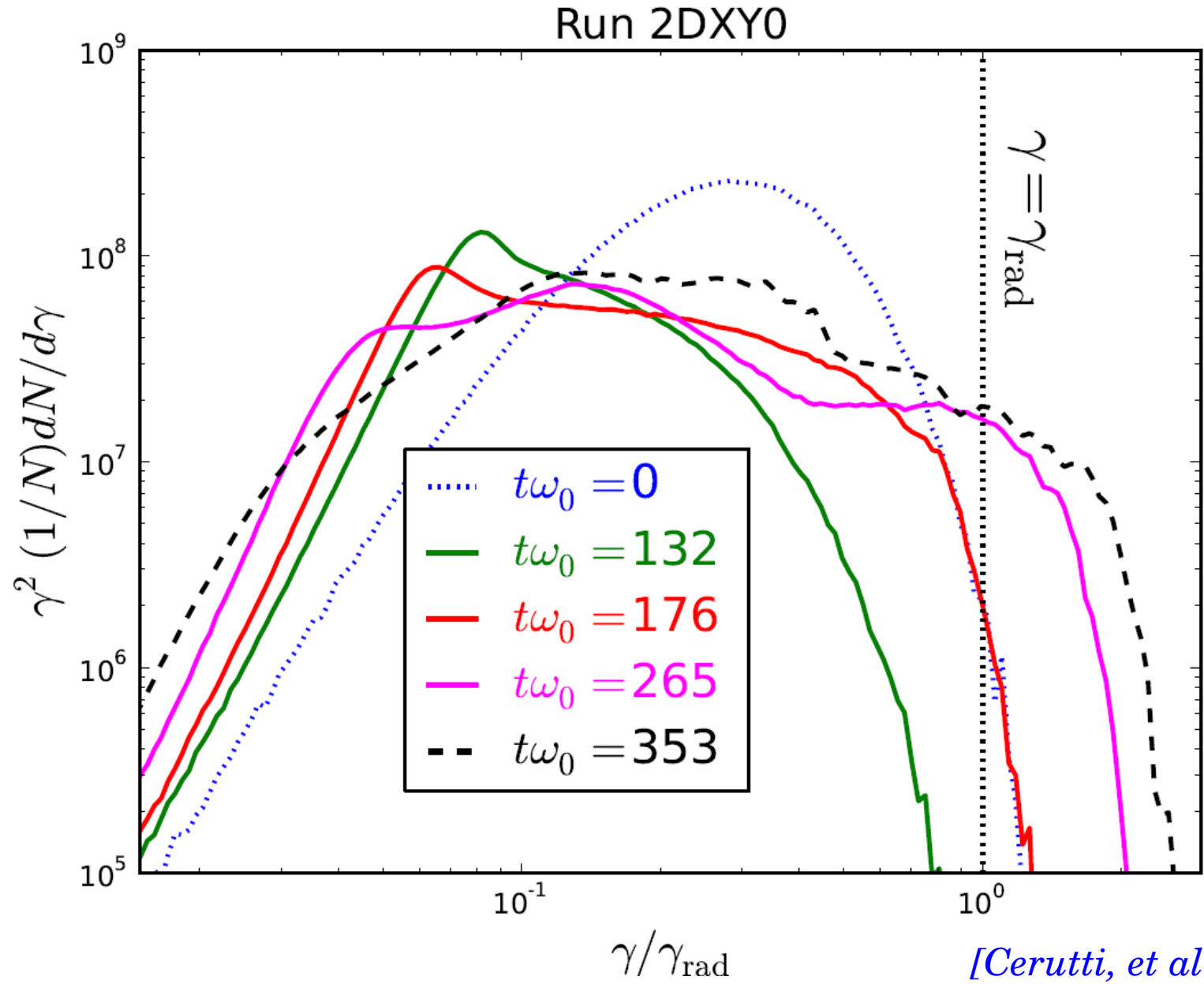
Under ideal MHD conditions: $E < B$ (ideal MHD) $\Rightarrow \epsilon_{\text{max}} < 160 \text{ MeV}$

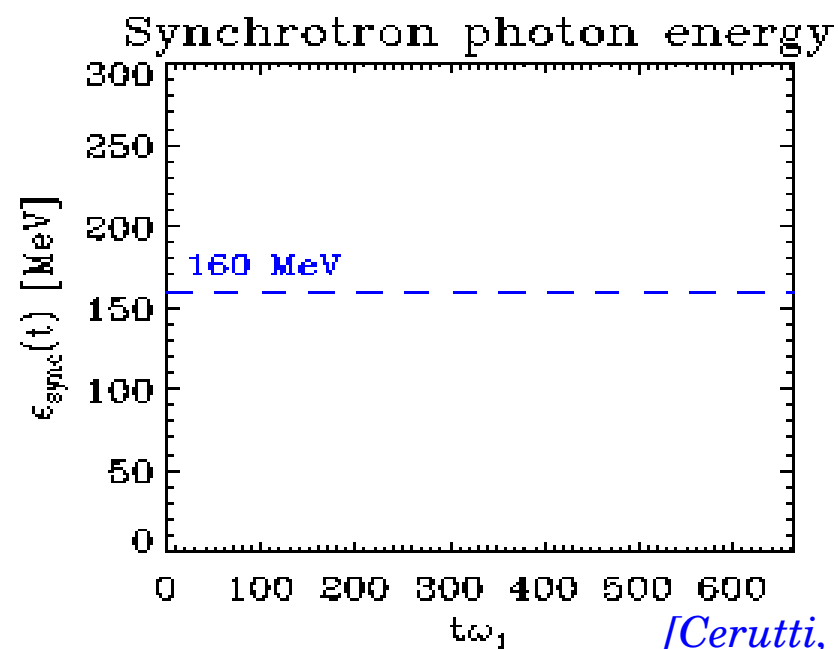
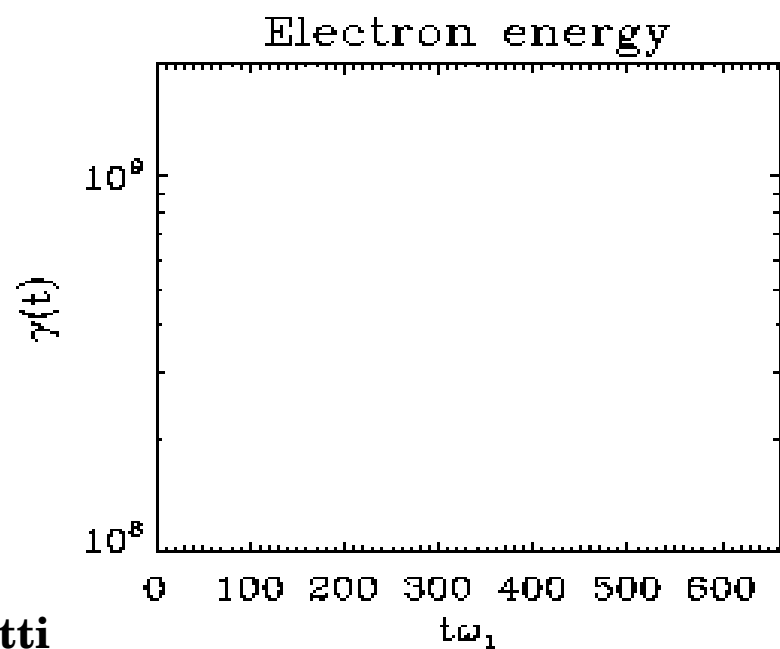
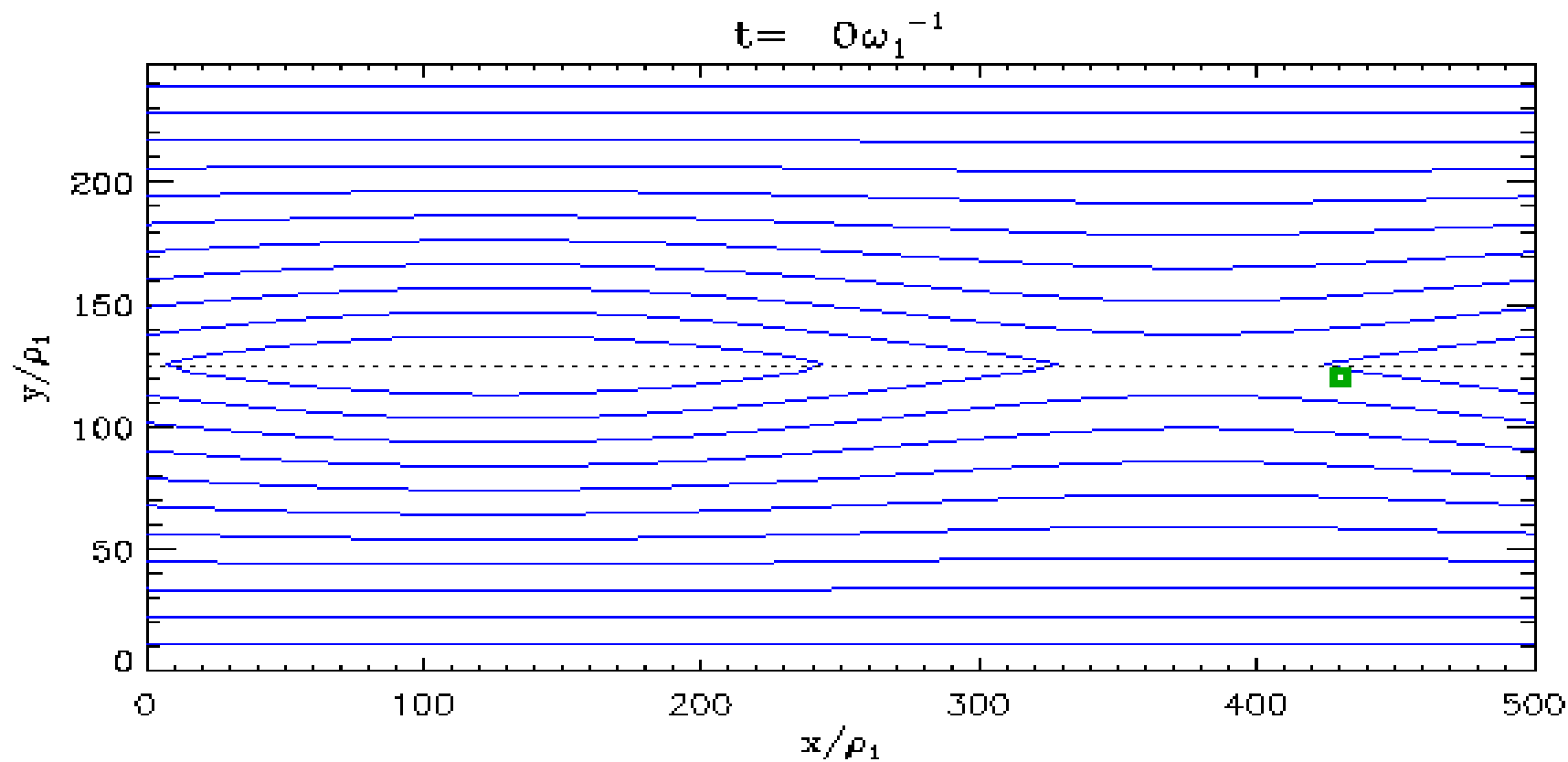
[e.g. Guilbert et al., 1983 ; de Jager et al., 1996 ; Uzdensky et al., 2011]

But this limit does not stricly apply to reconnection where $E > B$ in the layer

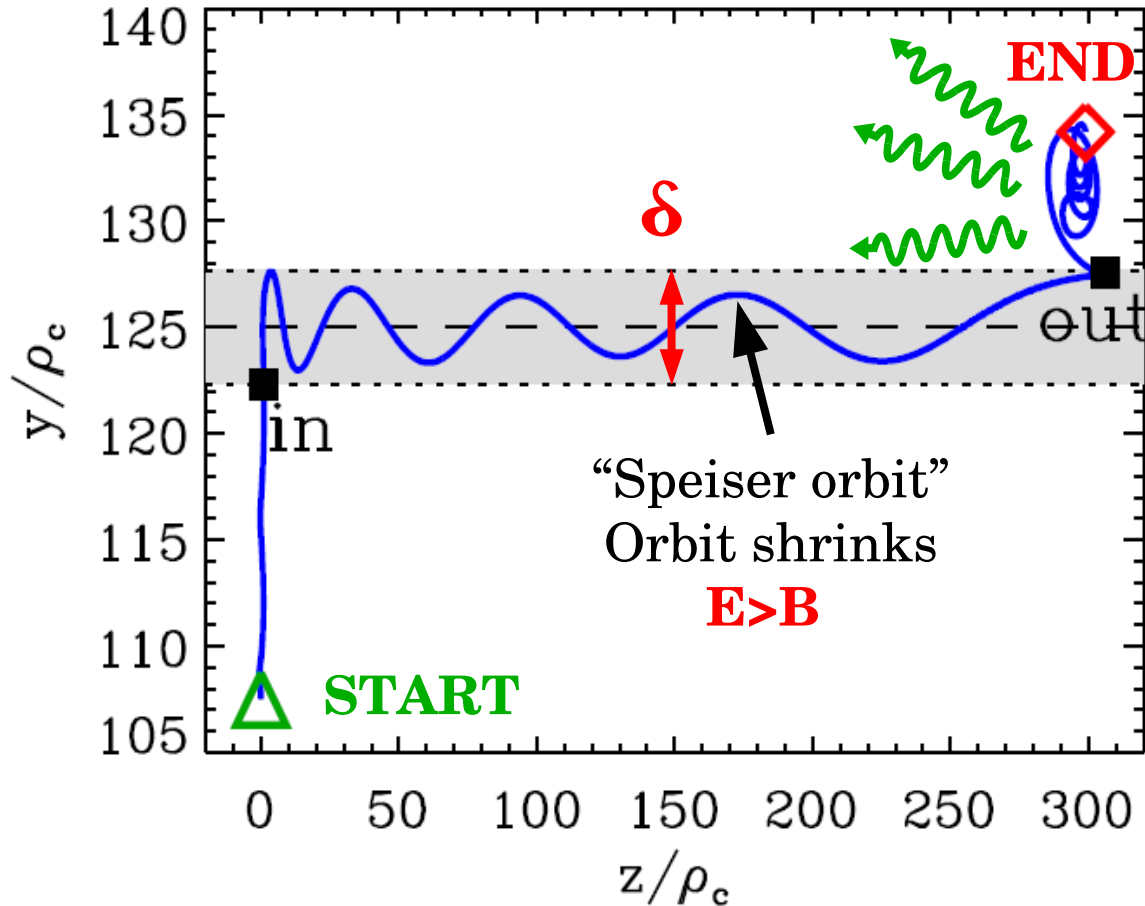
Particle acceleration beyond radiation reaction limit

2D PIC with radiation reaction force

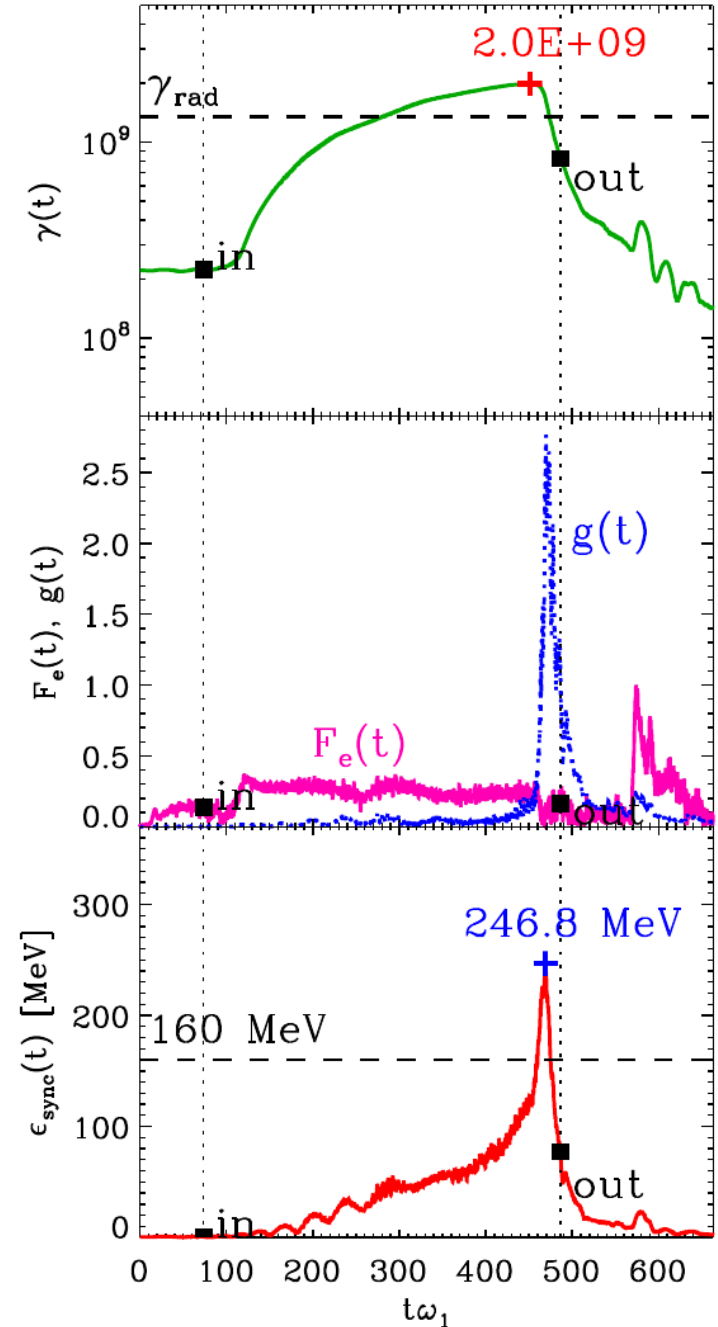




A typical high-energy particle orbit with $\gamma > \gamma_{\text{rad}}$



- Phase 1.** Drifting towards the layer
- Phase 2.** Linear acceleration, weak rad. losses, where $E > B$ (non-ideal MHD)
- Phase 3.** Ejection, fast cooling and emission of >160 MeV synchrotron.



Relativistic electron-ion reconnection

Regime relevant to **hot collisionless accretion flows** (e.g., Sgr A*), black hole coronae, **AGN jets**.

$$\sigma_e = \frac{B_0^2}{4\pi n_e m_e c^2}$$

$$\sigma_i = \frac{B_0^2}{4\pi n_i m_i c^2} = \sigma_e \left(\frac{m_e}{m_i} \right)$$

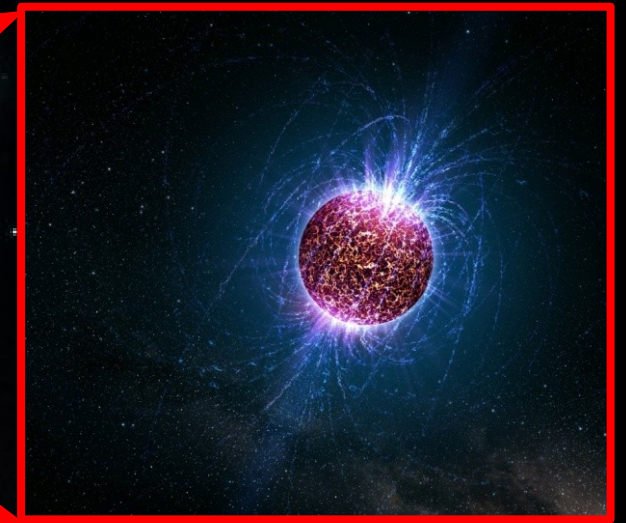
- **Non-relativistic regime** : $\sigma_i \ll 1$, $\sigma_e \ll 1$: Poor particle acceleration for both species, steep power laws.
- **Semi-relativistic regime** : $\sigma_i \ll 1$, $\sigma_e \gg 1$: Mildly relativistic ions, ultra-relativistic electrons. *Melzani+2014* (**see Journal Club**), *Werner+2017*.
- **Ultra-relativistic regime** : $\sigma_i \gg 1$, $\sigma_e \gg 1$: Identical to pair plasmas since γmc^2 enters in plasma scales (skin depth, Larmor radius & frequencies...), if ignoring radiative effects.

II. Astrophysical applications

Astrophysical flares

What is the Crab Nebula ?

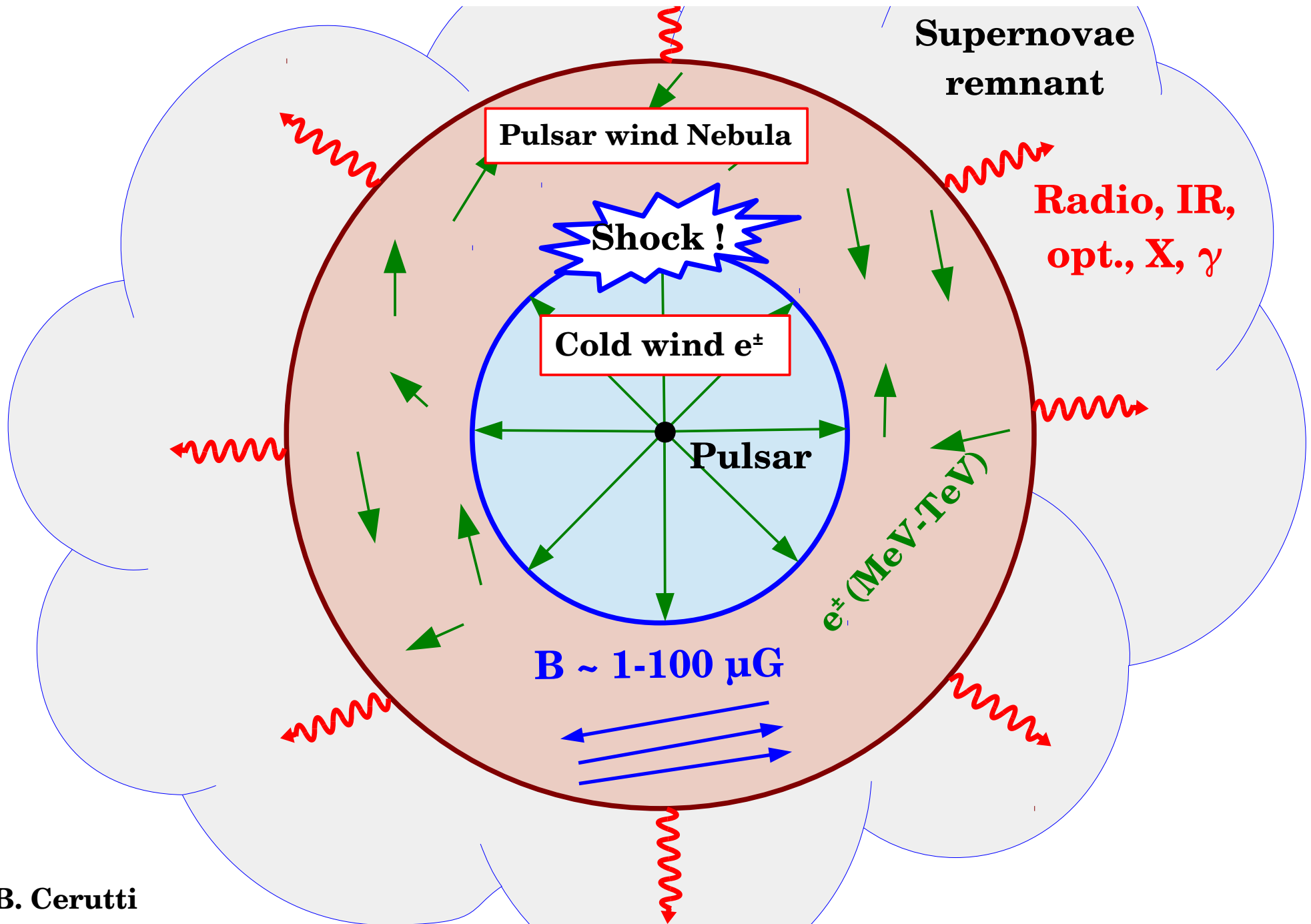
- Born after a supernova explosion
- Birth date: **1054 AD**
- Distance: **2 - 2.5 kpc**
- Size: **~ 1 pc**



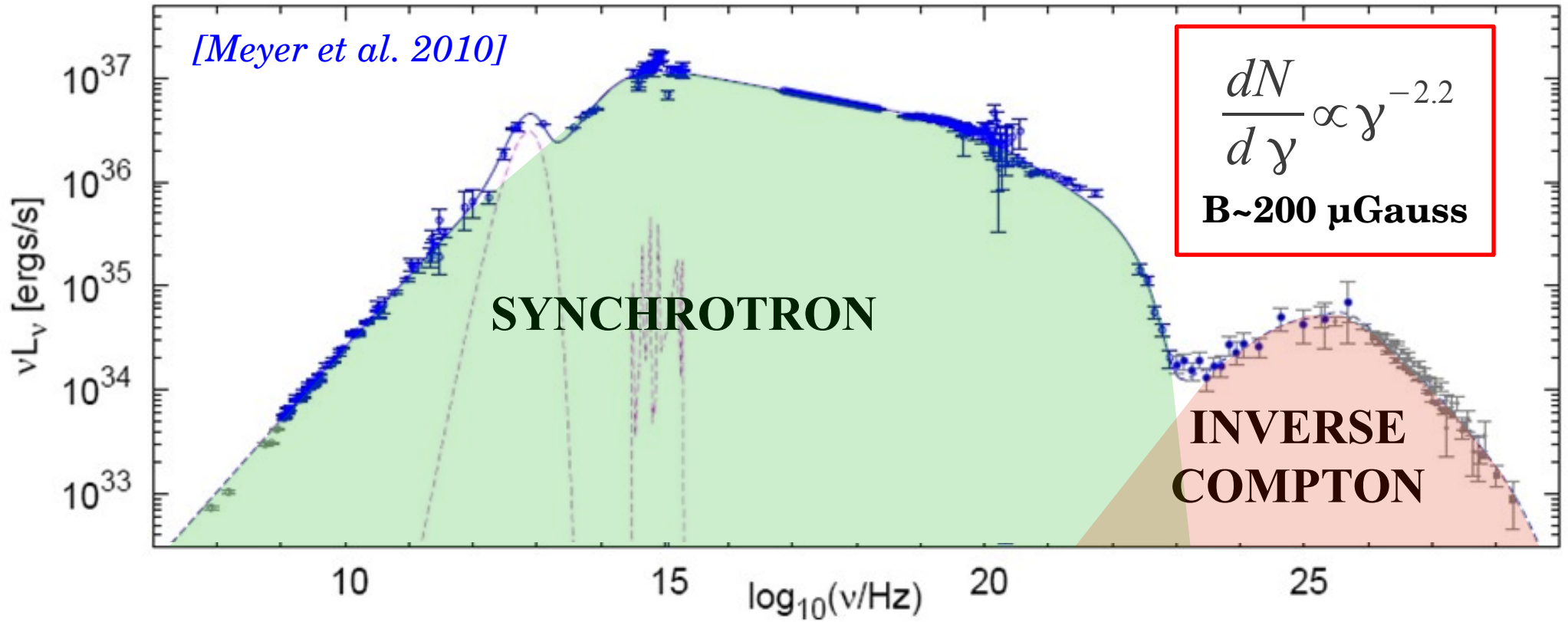
Crab Pulsar:

- Spin period: **33 ms**
- Magnetic field: **$\sim 4 \times 10^{12}$ G**
- Radius: **~ 10 km**

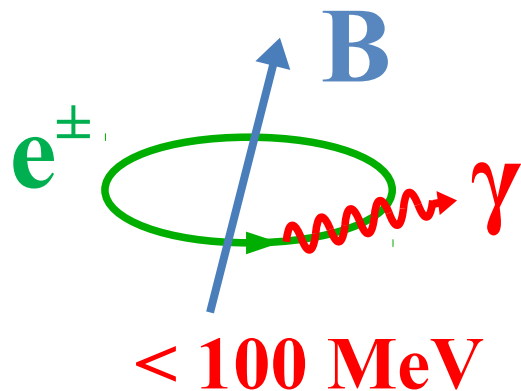
Pulsar wind nebulae



Crab nebula broad band emission



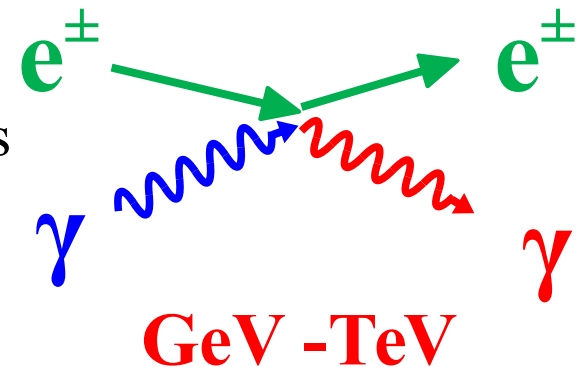
Synchrotron radiation



Inverse Compton scattering

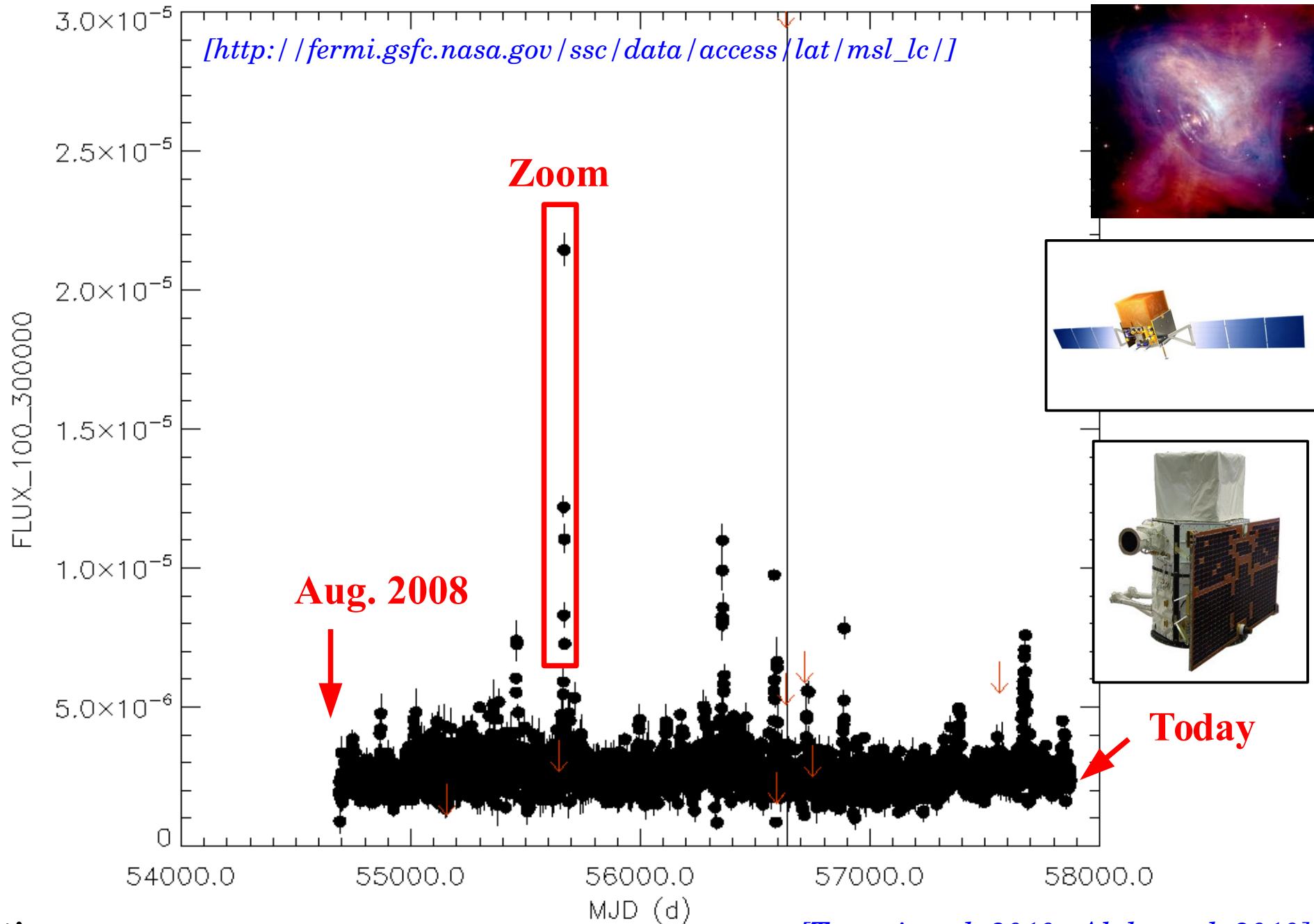
Target photons:

- Synchrotron photons
- Dust (93 K)
- CMB (2.7 K)



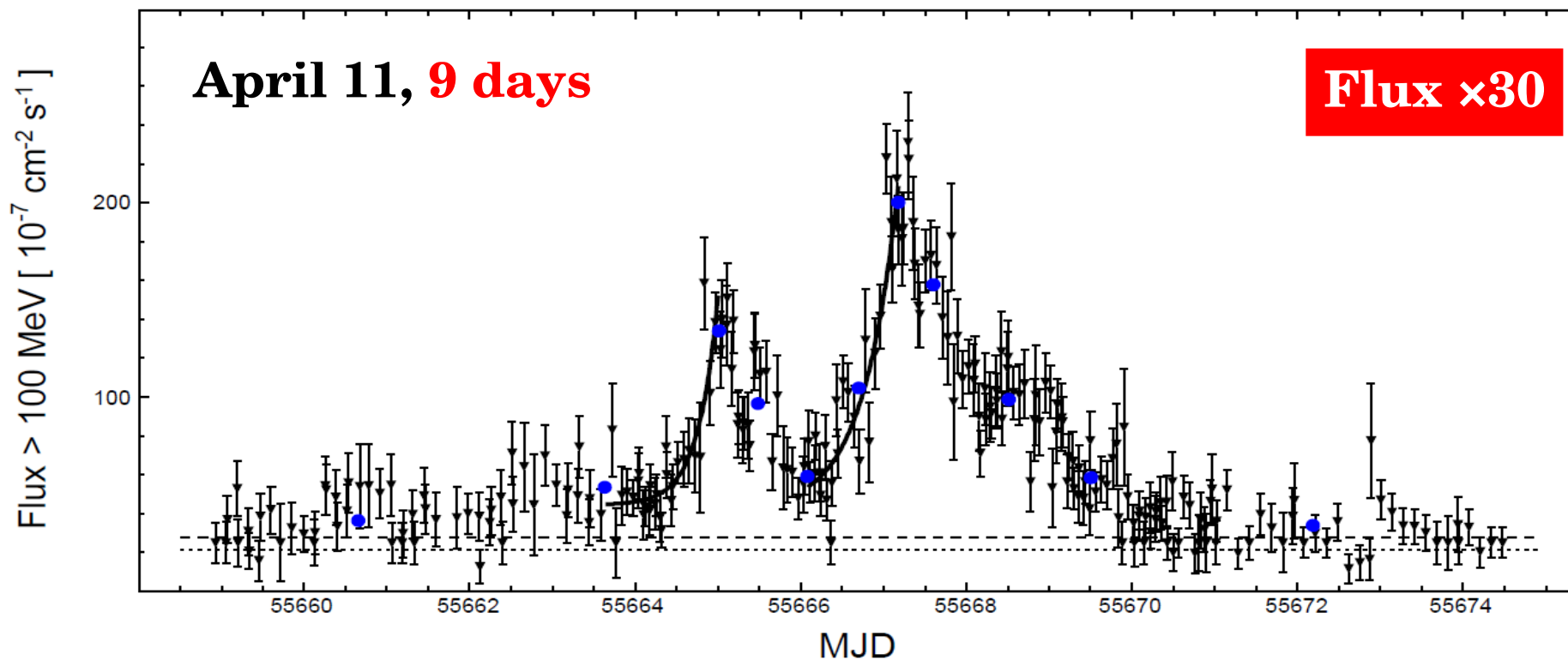
Time variations in the >100 MeV flux

Source = Crab Pulsar Duration = 86400.0



Ultra-short time variability

[Buehler et al., 2012]



(1) Size : $ct_{\text{flare}} \sim 10^{16}$ cm, **~1% radiates 30 times** the whole Nebula flux

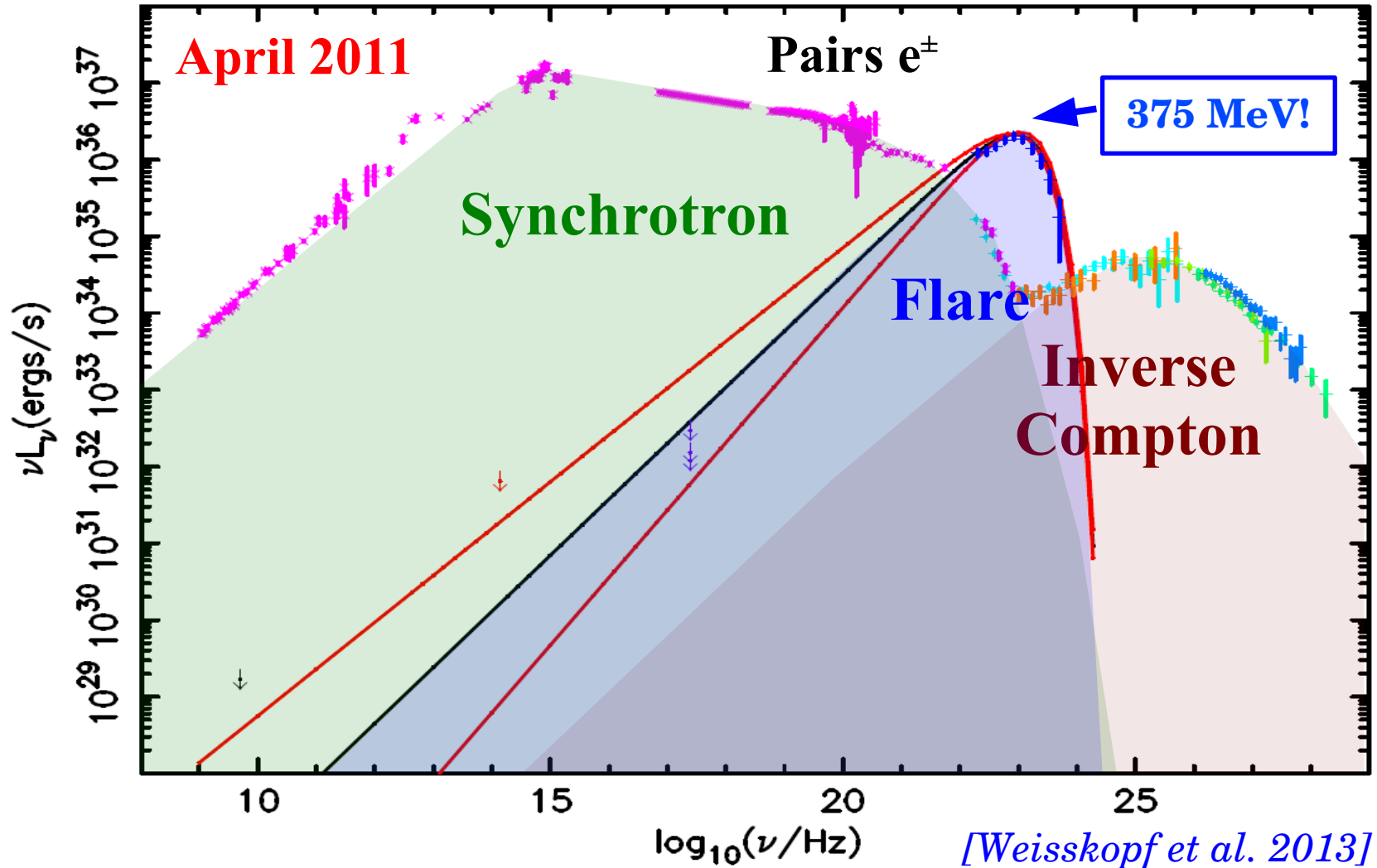
(2) If $t_{\text{flare}} = t_{\text{sync}} \Rightarrow$ **B ~ few mG \gg 200 μ G, and PeV pairs**

(3) PeV pairs $\Rightarrow t_{\text{gyration}} \sim t_{\text{flare}}$, pairs accelerated within ~ 1 Cyclotron orbit

Inconsistent with diffuse shock acceleration

The acceleration is **turned ON** during the gamma-ray flare

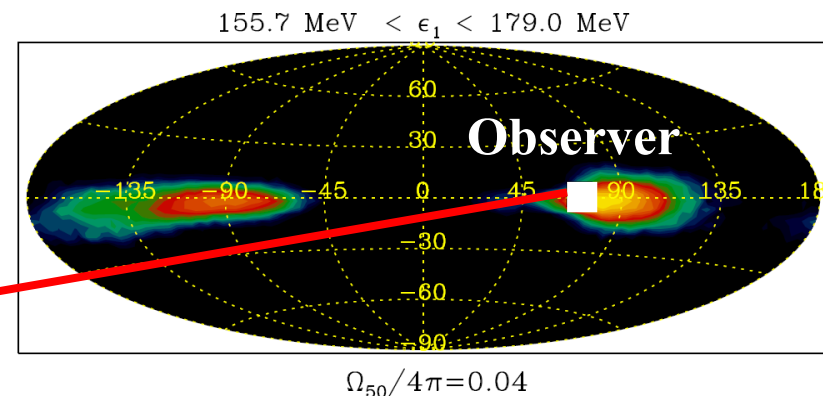
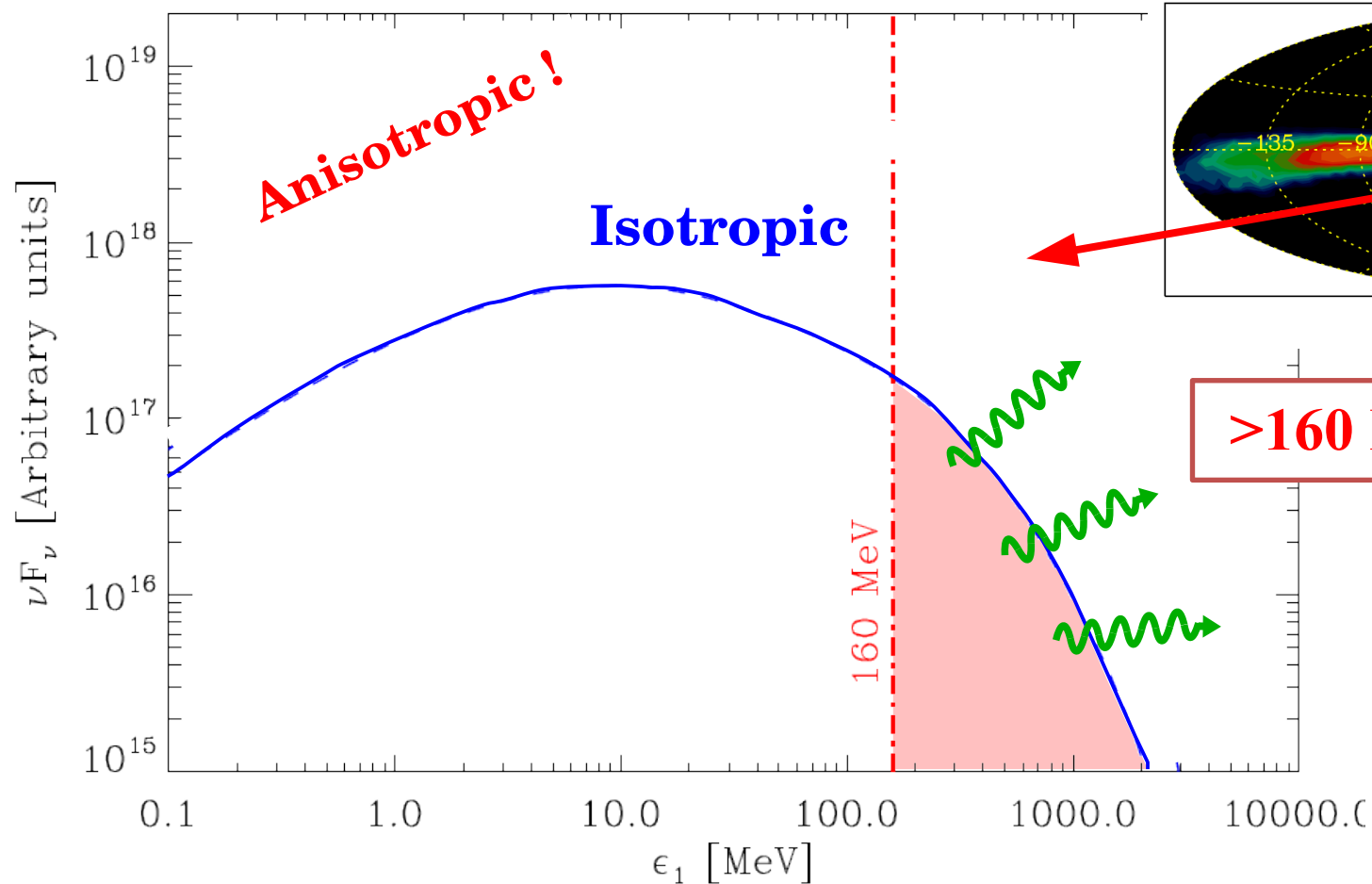
Flare spectral component



- (1) **No obvious counter-parts** (radio, IR, opt., X, TeV)
- (2) **Spectrum ~mono-energetic**, inconsistent with shock-acceleration
- (3) **Synchrotron photons > 100 MeV**

PIC simulations of radiative relativistic reconnection

Total synchrotron flux (optically thin)



>160 MeV!

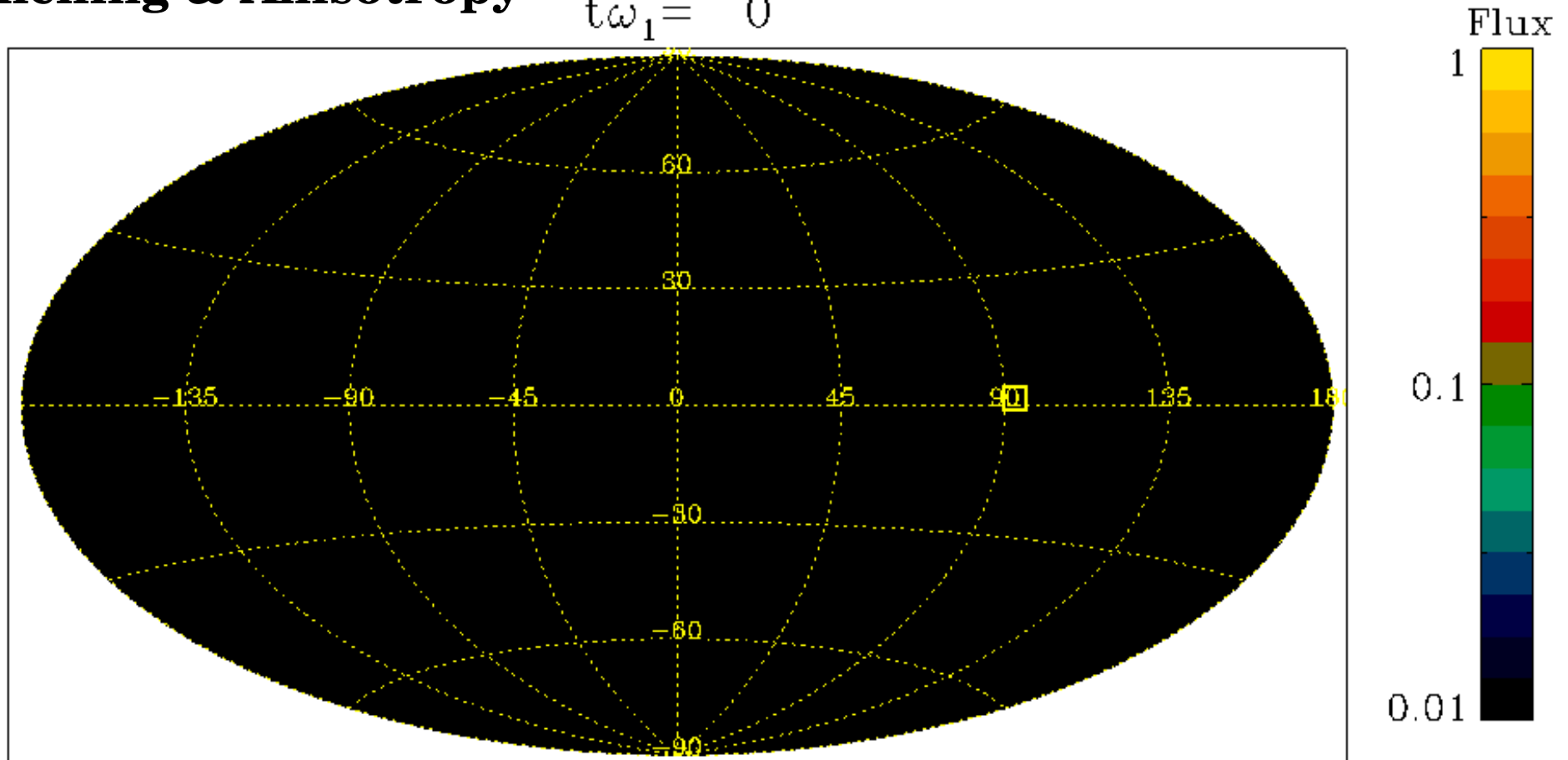
[Cerutti et al. 2013]

Apparent high-energy flux INCREASED!
Good to reduce the energetic constraints !

Time variations of the >100 MeV emission

Bunching & Anisotropy

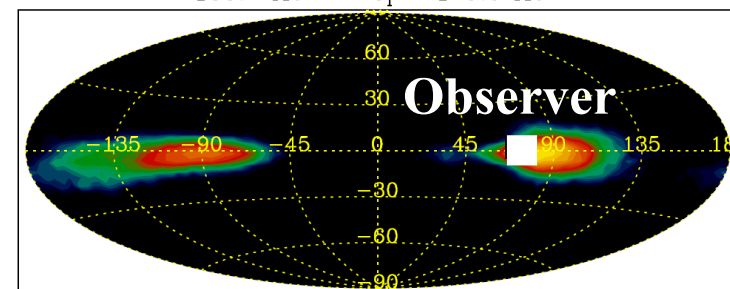
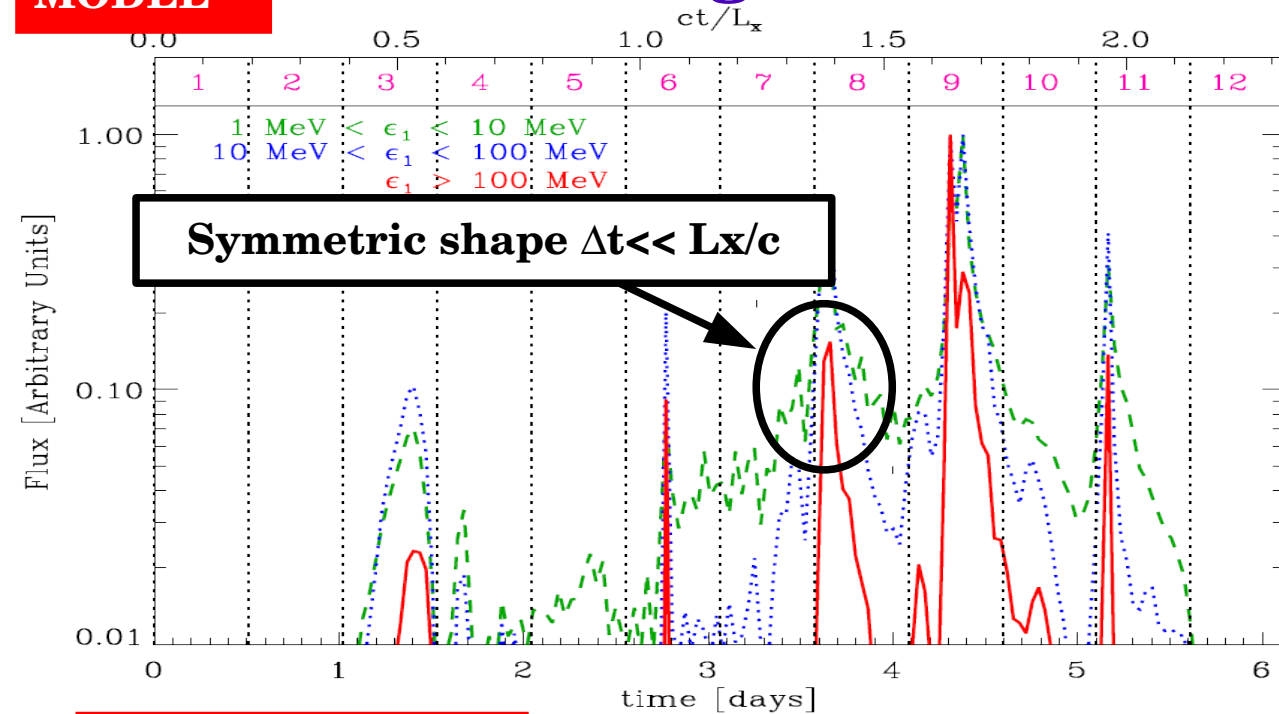
$$t\omega_1 = 0$$



The beam of high-energy radiation **sweeps across the line of sight** intermittently \Rightarrow **bright symmetric flares.**

Light curve modelling

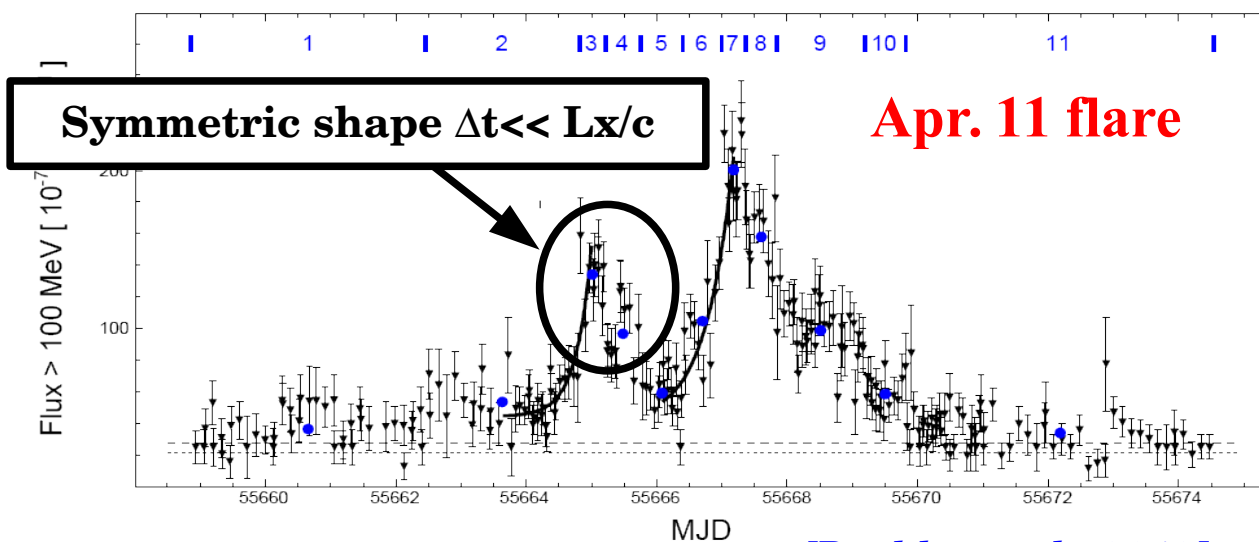
MODEL



Ultra-short time variability
< 6 hours due to particle
bunching and anisotropy

OBSERVATIONS

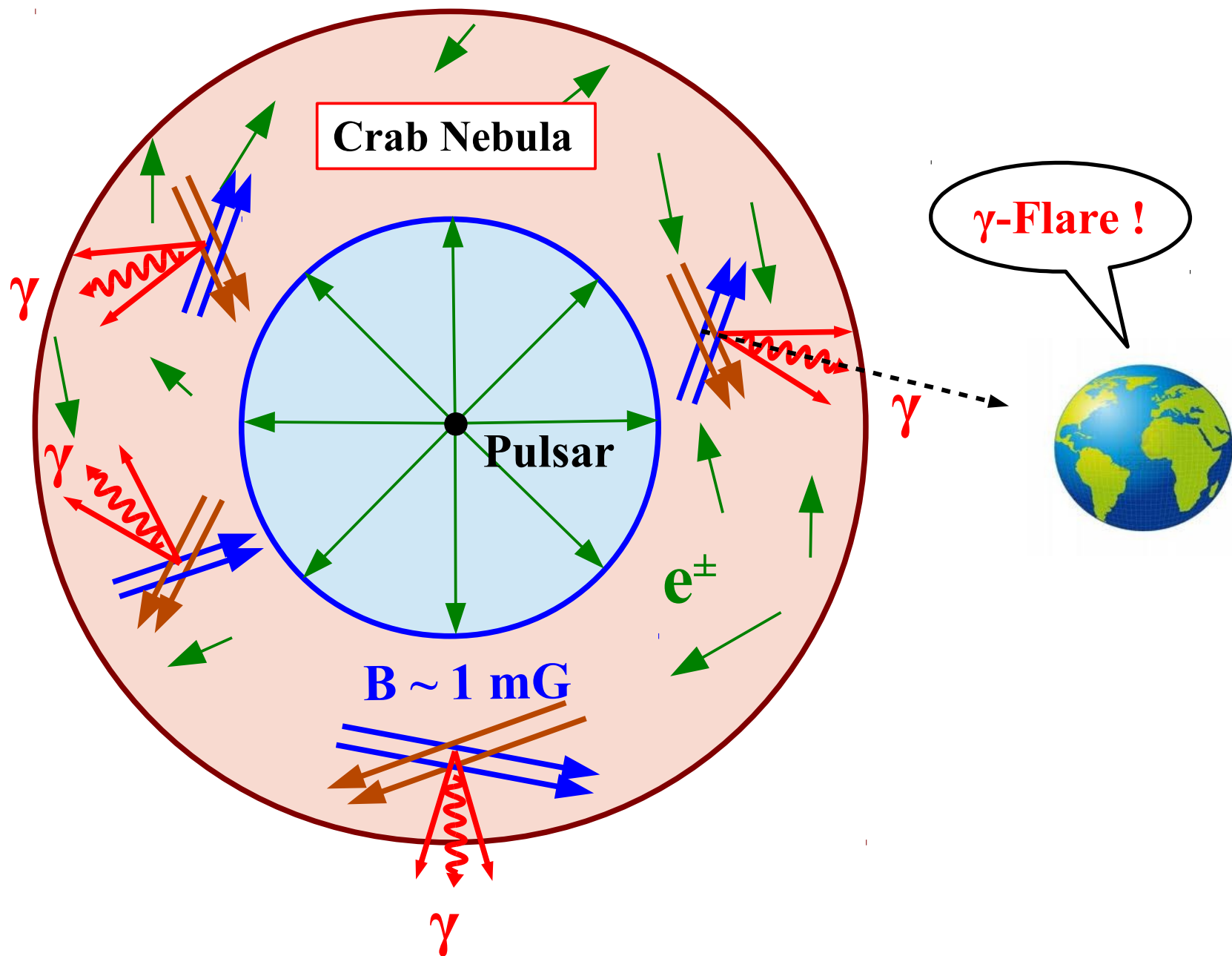
[Cerutti et al. 2013]



Observed ultra-short time
 variability **< 8 hours**

[Buehler et al., 2012]

Flare when reconnection site oriented towards Earth

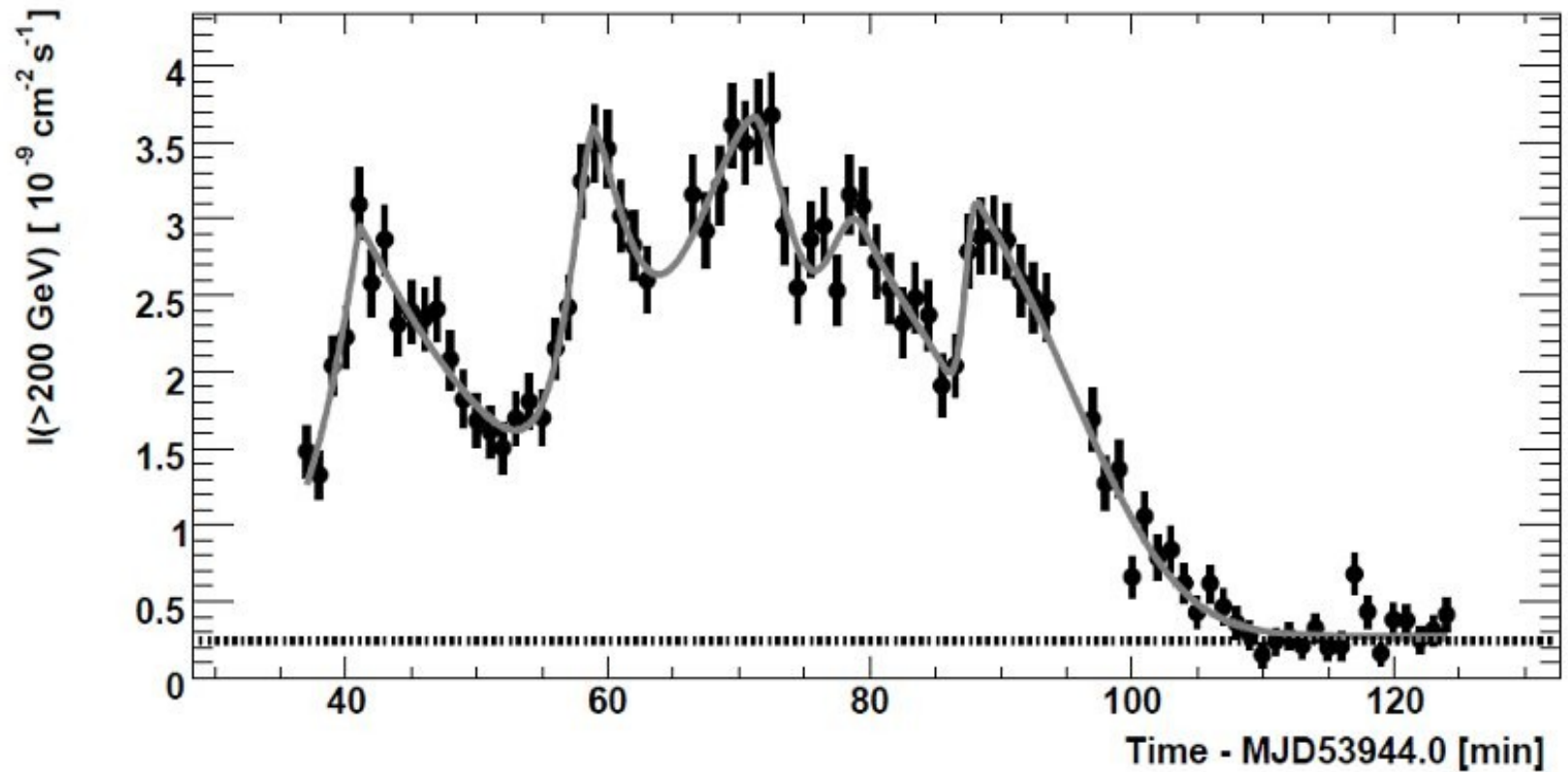


Application to TeV blazar flares ?

PKS 2155-304 (and others !): short-time variability $\sim 200 \text{ s} \ll R_{\text{BH}}/c$.

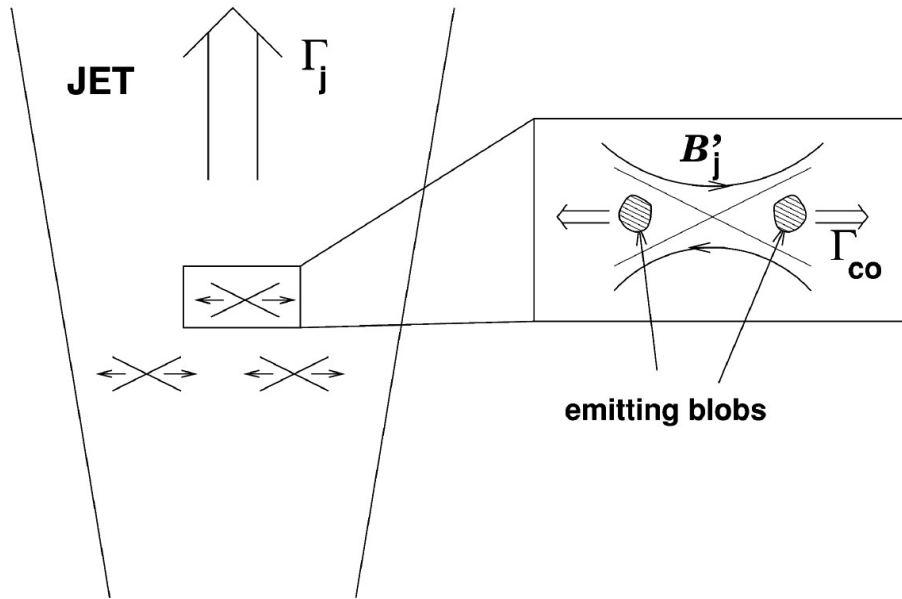
=> Requires high Doppler factor $\delta_{\text{dopp}} > 50$.

[Saugé & Henri 2006 ; Begelman et al. 2008]



Aharonian et al. 2007 (H.E.S.S. Collaboration)

Jet-in-a-jet scenario for blazar super-fast flares

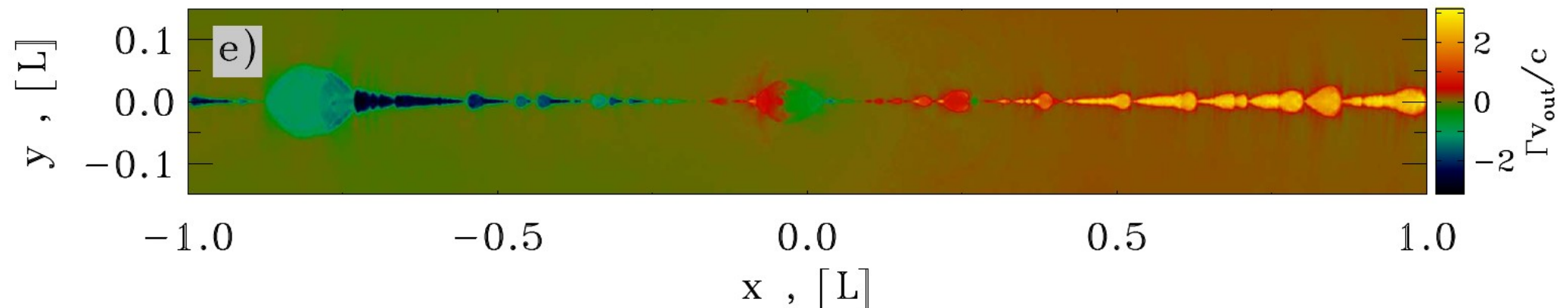


Jet bulk Lorentz factor combined with ultra-relativistic reconnection outflows :

$$\Gamma_{em} = \Gamma_j \Gamma_{co} (1 + v_j v_{co} \cos \theta')$$

$$\Gamma_{co} \approx \sqrt{\sigma} \quad [Lyubarsky 2005]$$

[Giannios et al. 2009 ; Giannios 2013]



Application to pulsars

Some of the big questions

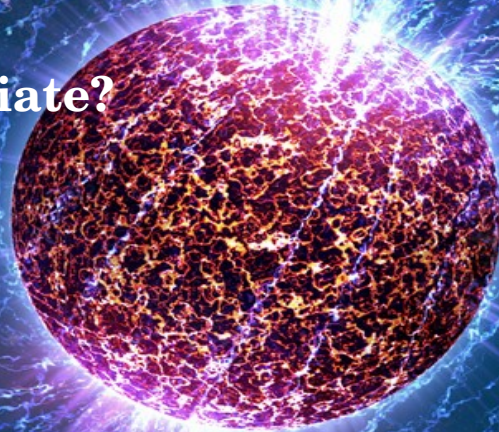
How does the star spin-down?

How is this energy channeled to particles and radiation?

How is the plasma generated?

How are particles accelerated and radiate?

Where is the emission coming from ?



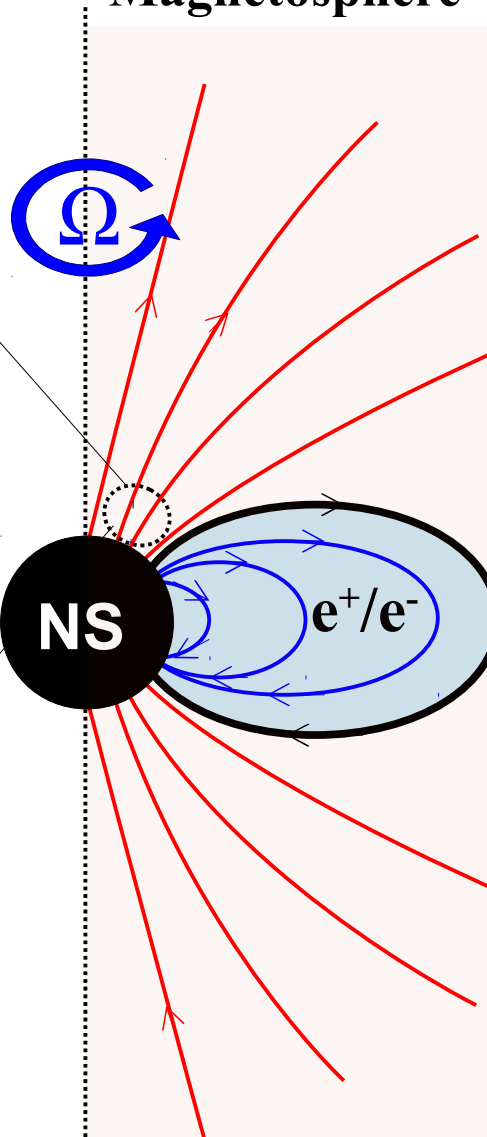
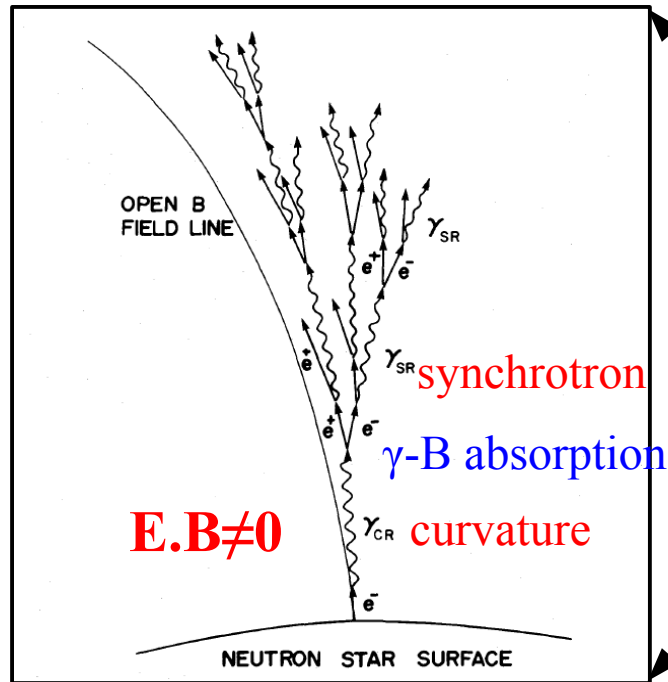
=> To address these questions, we need a model of the magnetosphere!

Elements of a pulsar magnetosphere

Magnetosphere

Wind region

Copious pair creation in the polar caps

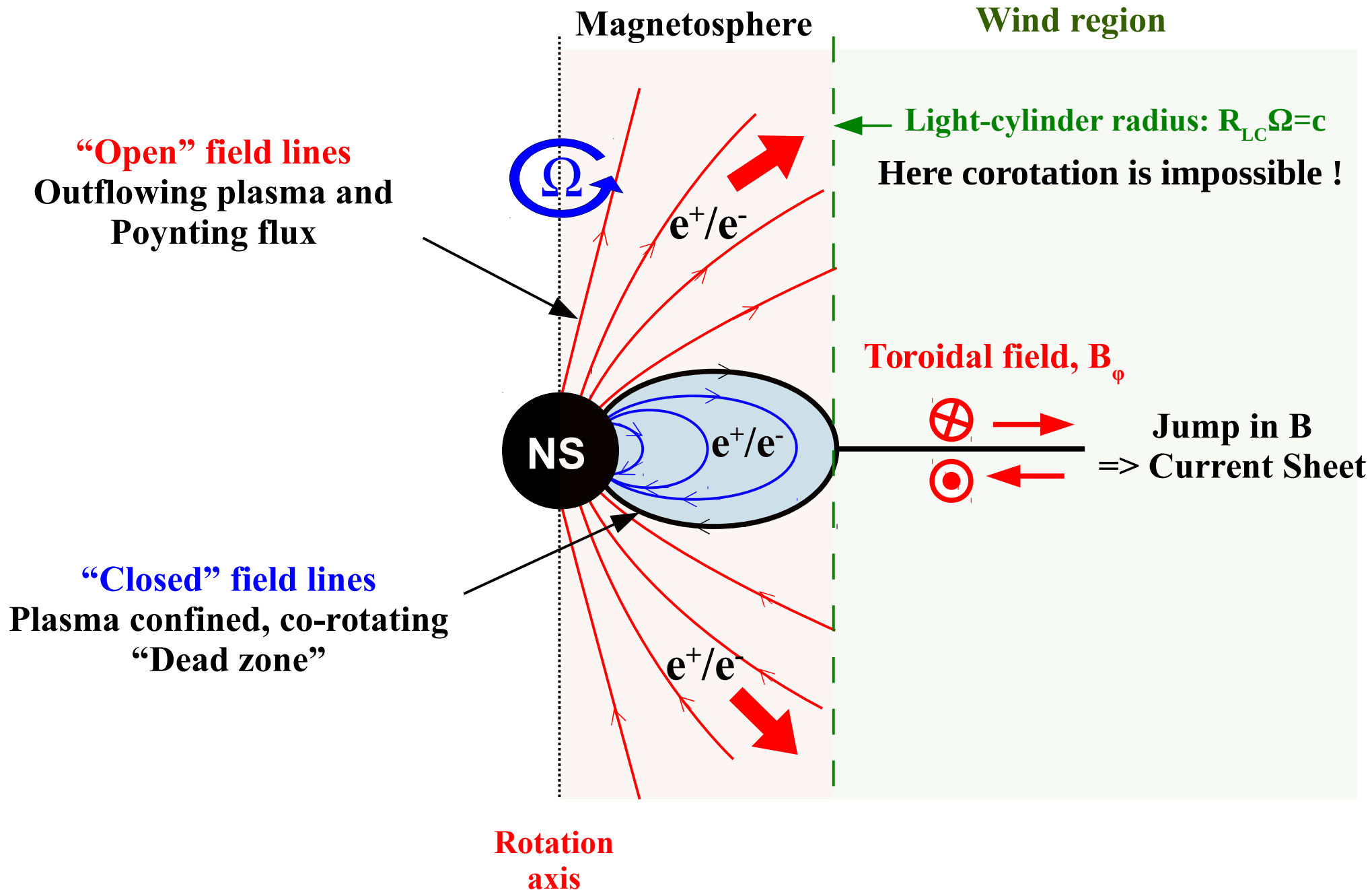


← Light-cylinder radius: $R_{LC} \Omega = c$
 Here corotation is impossible !

Rotation axis

Daugherty & Harding 1982 ; Timokhin & Arons 2013 ; Chen & Beloborodov 2014 ; Philippov et al., 2015

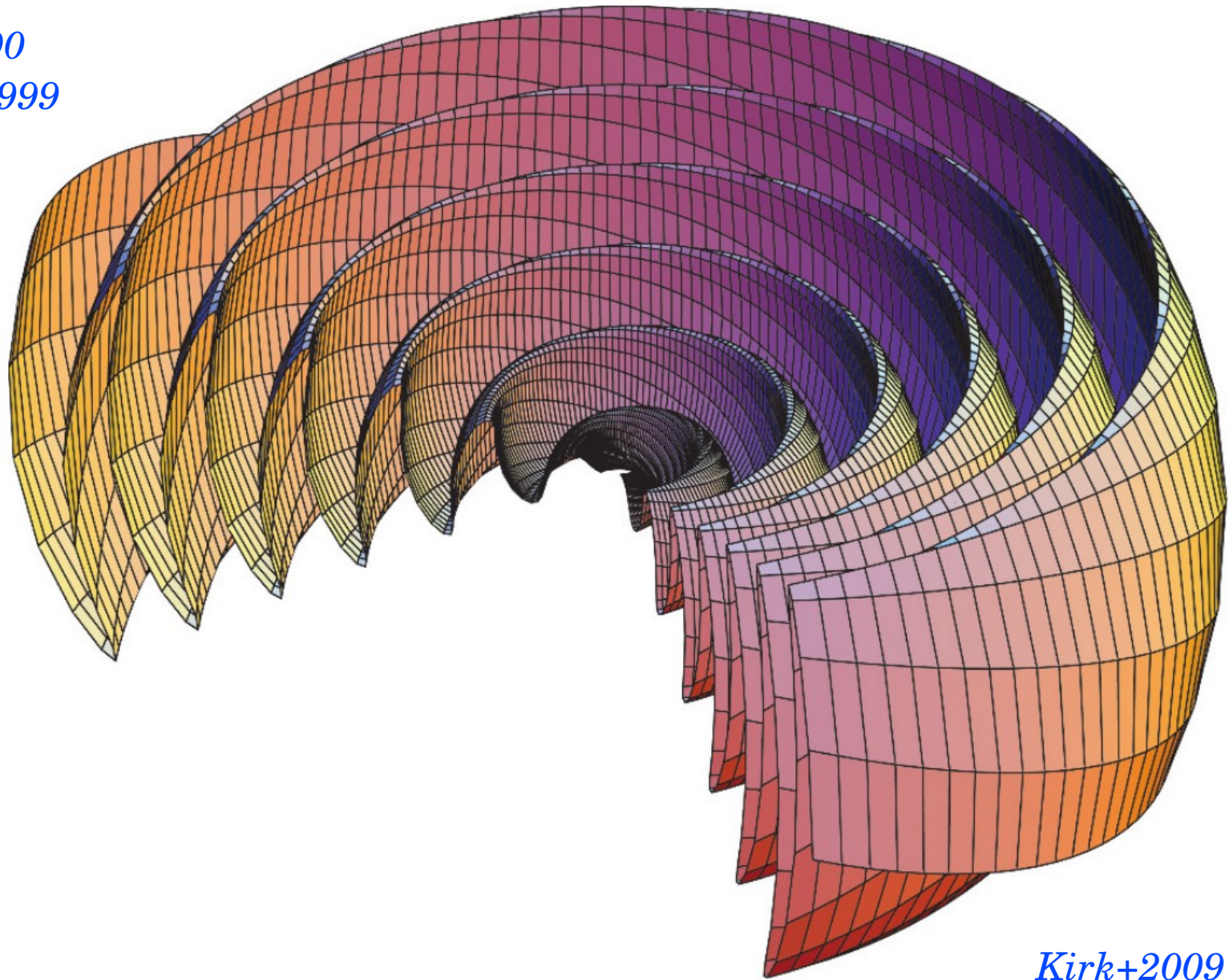
Elements of a pulsar magnetosphere



Inclined rotation : « Striped » pulsar current sheet

Coroniti 1990

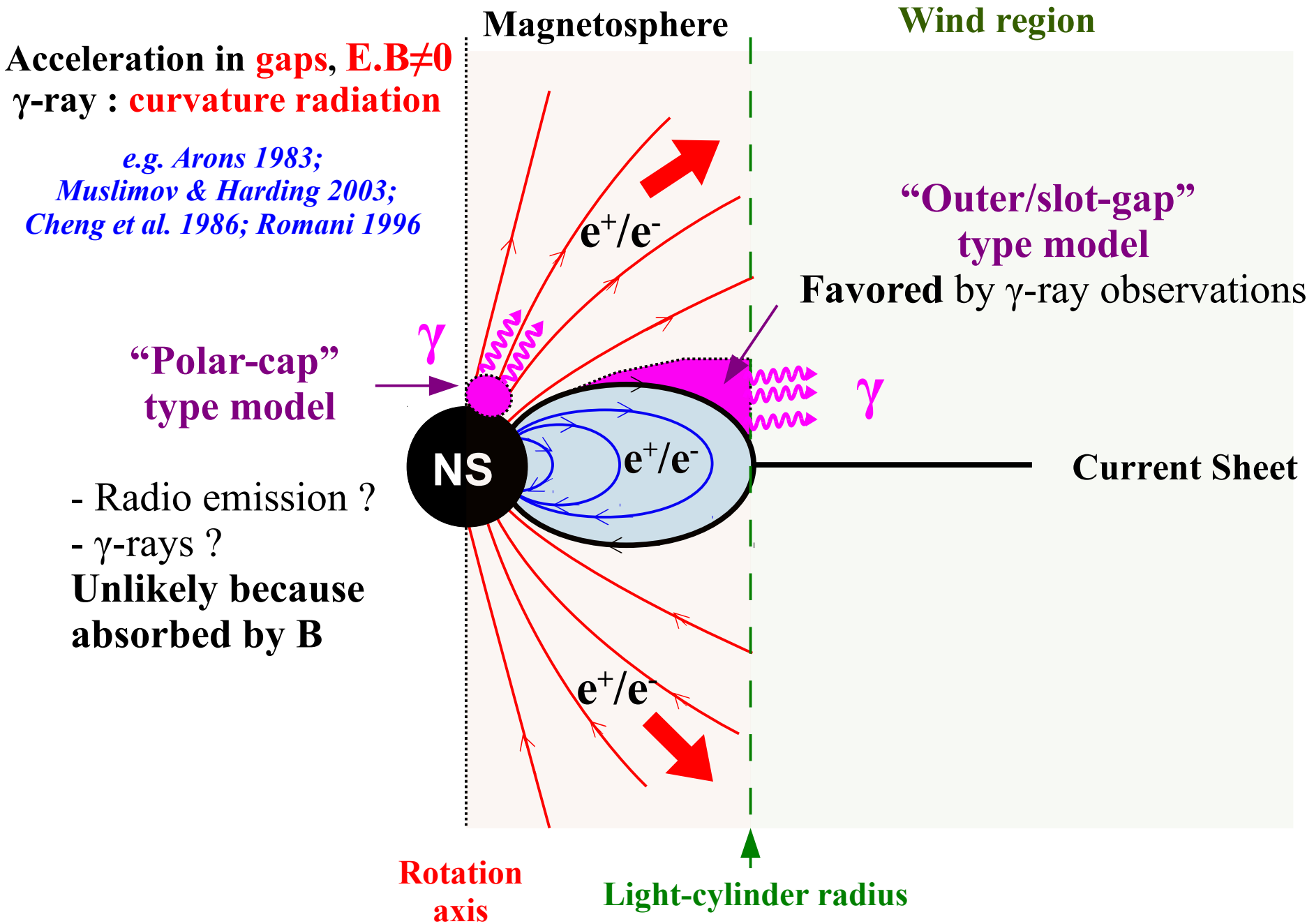
Bogovalov 1999



Kirk+2009

Relativistic analog of the heliospheric current sheet

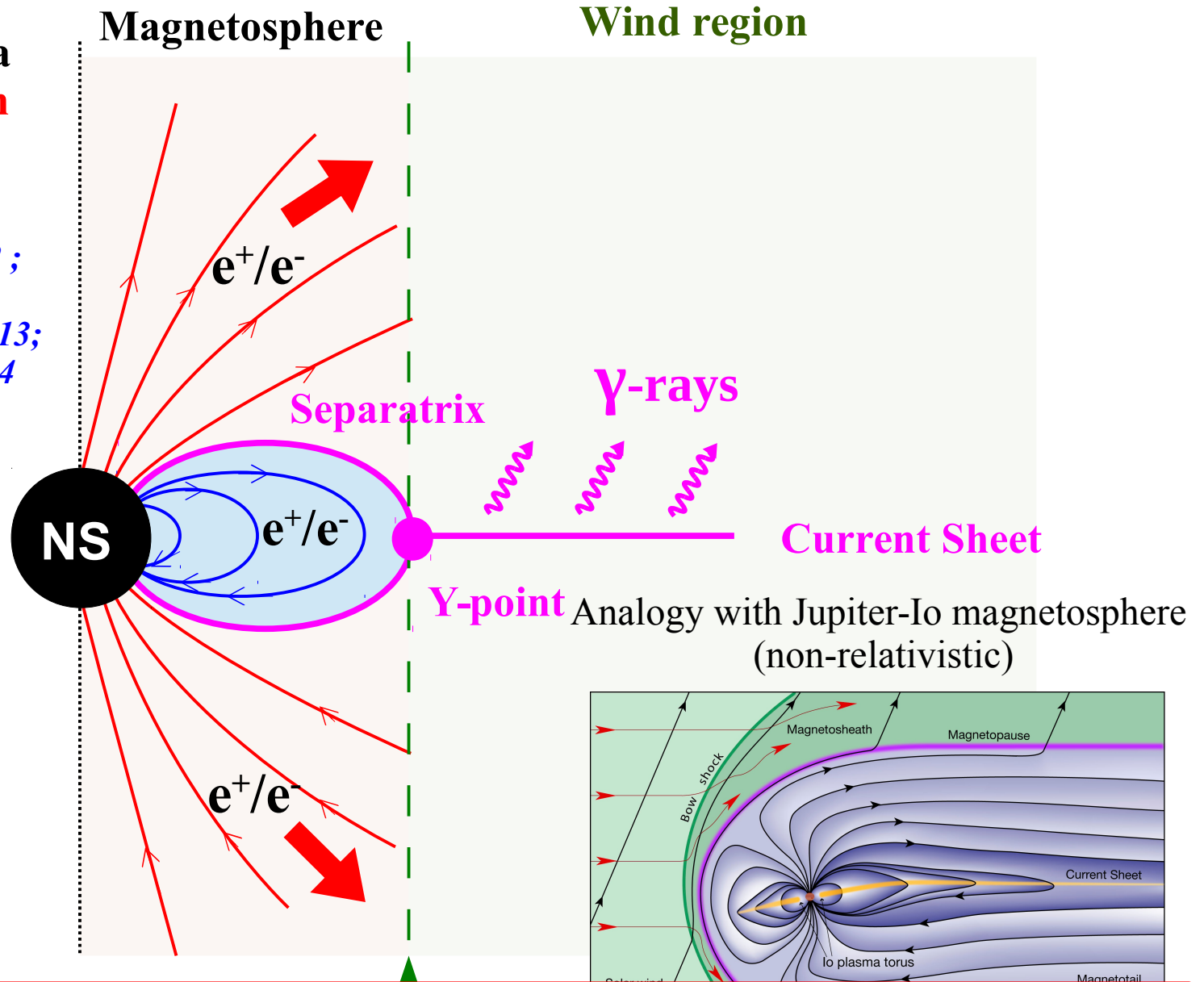
Proposed sites for particle acceleration



Proposed sites for particle acceleration

Particle acceleration via
relativistic reconnection
 γ -ray: **Synchrotron**

*Coroniti 1990 ;
Lyubarskii 1996 ; Kirk+2002 ;
Bai & Spitkovsky 2010 ;
Pétri 2012 ; Arka & Dubus 2013 ;
Uzdensky & Spitkovsky 2014
Mochol & Pétri 2015*



Models dependent on the geometry of the magnetosphere

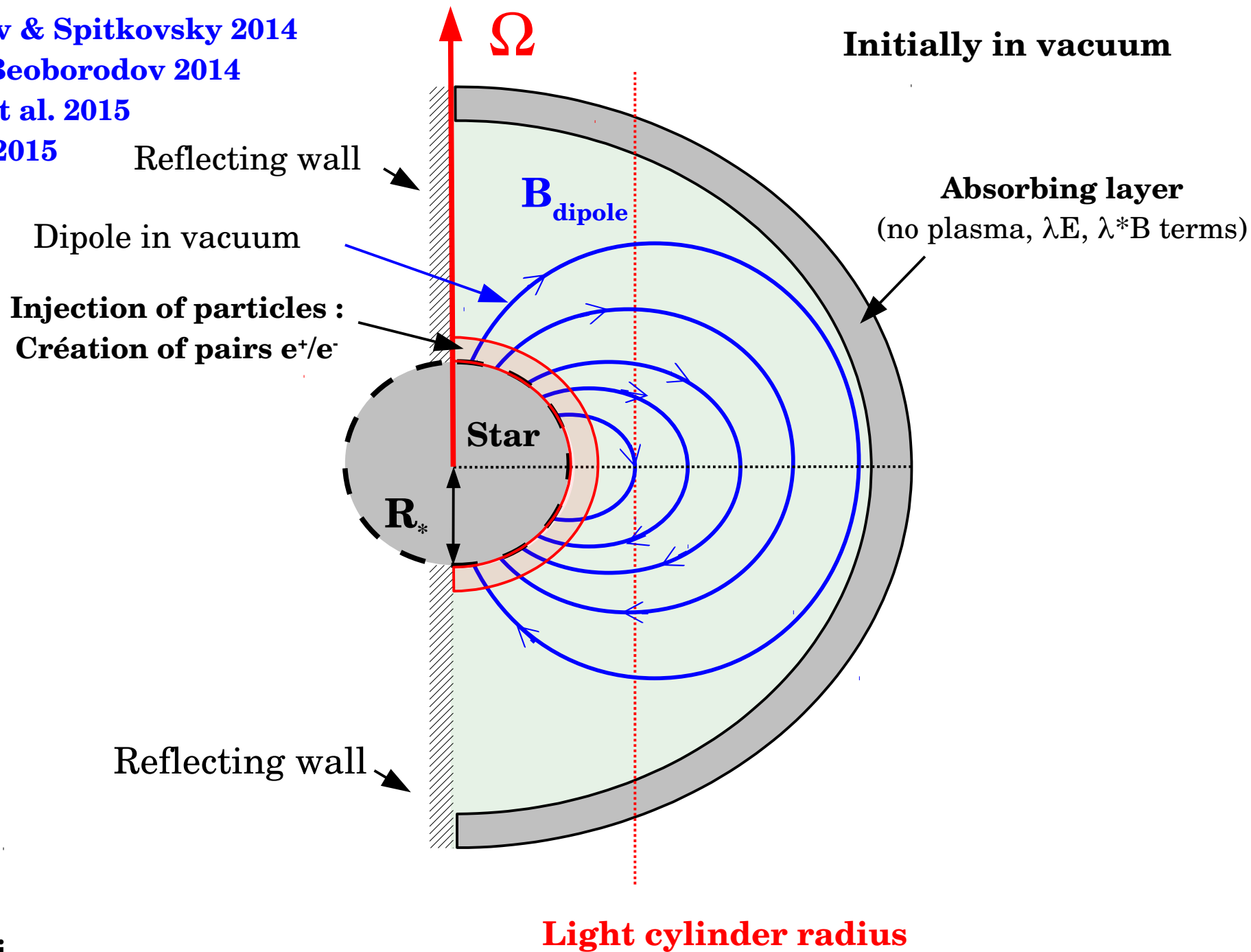
PIC setup : Aligned rotator (2D)

Philippov & Spitkovsky 2014

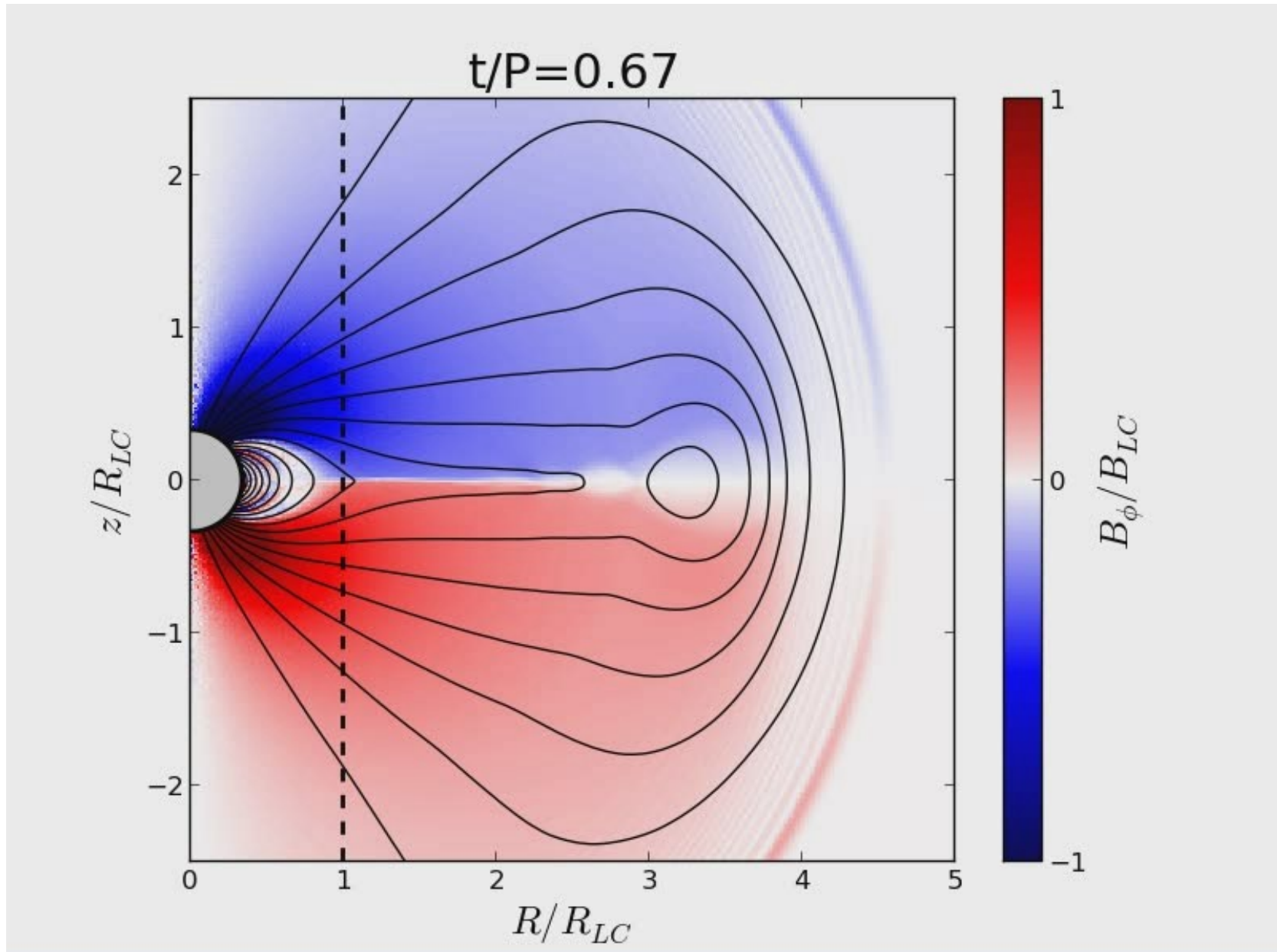
Chen & Beorodov 2014

Cerutti et al. 2015

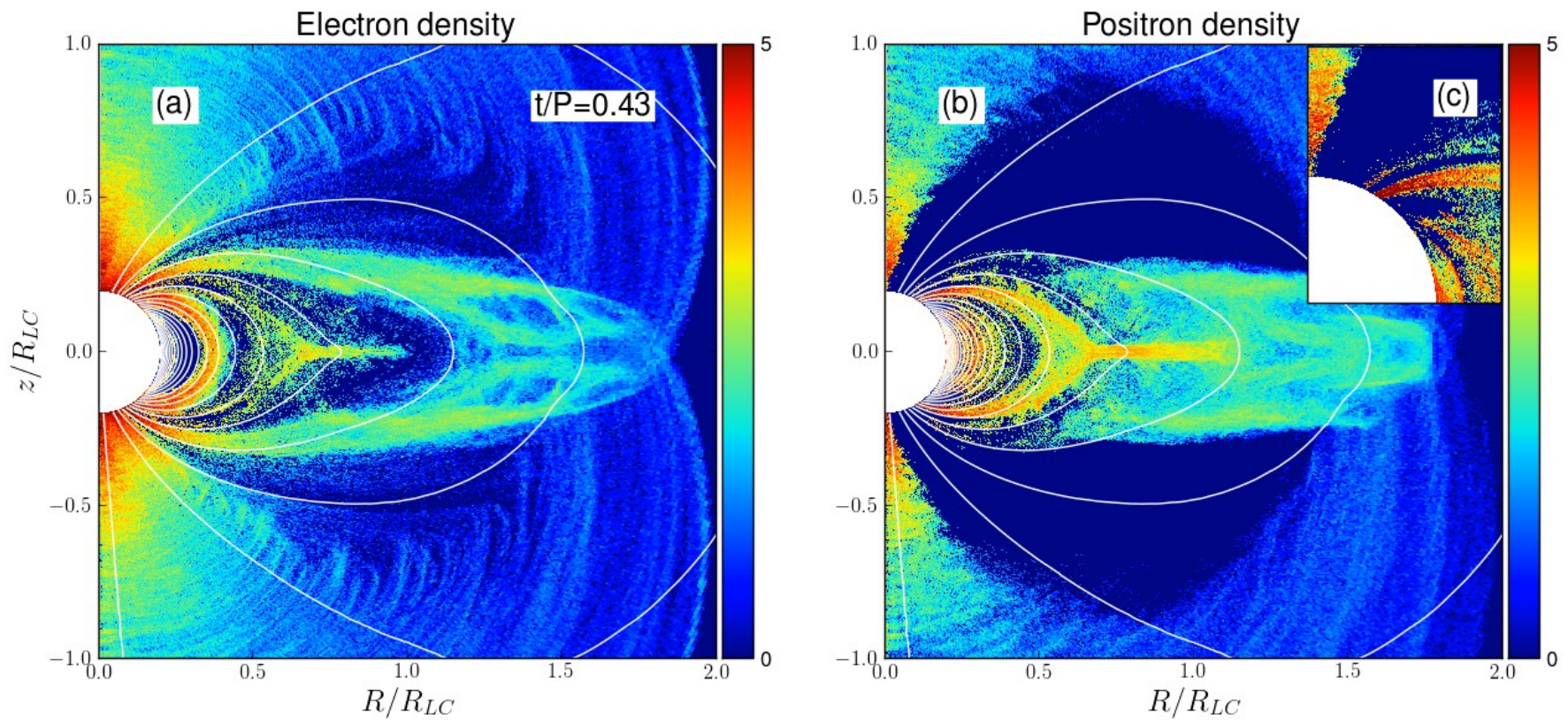
Belyaev 2015



Toroidal magnetic field



Plasma density, pair creation

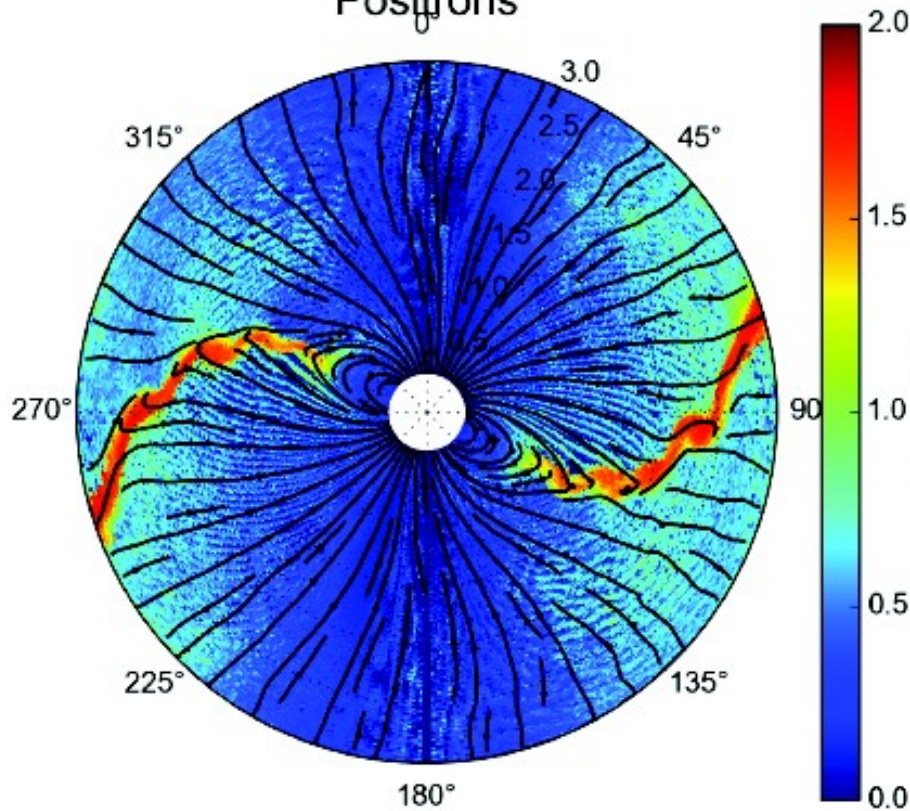


Philippov et al. 2015

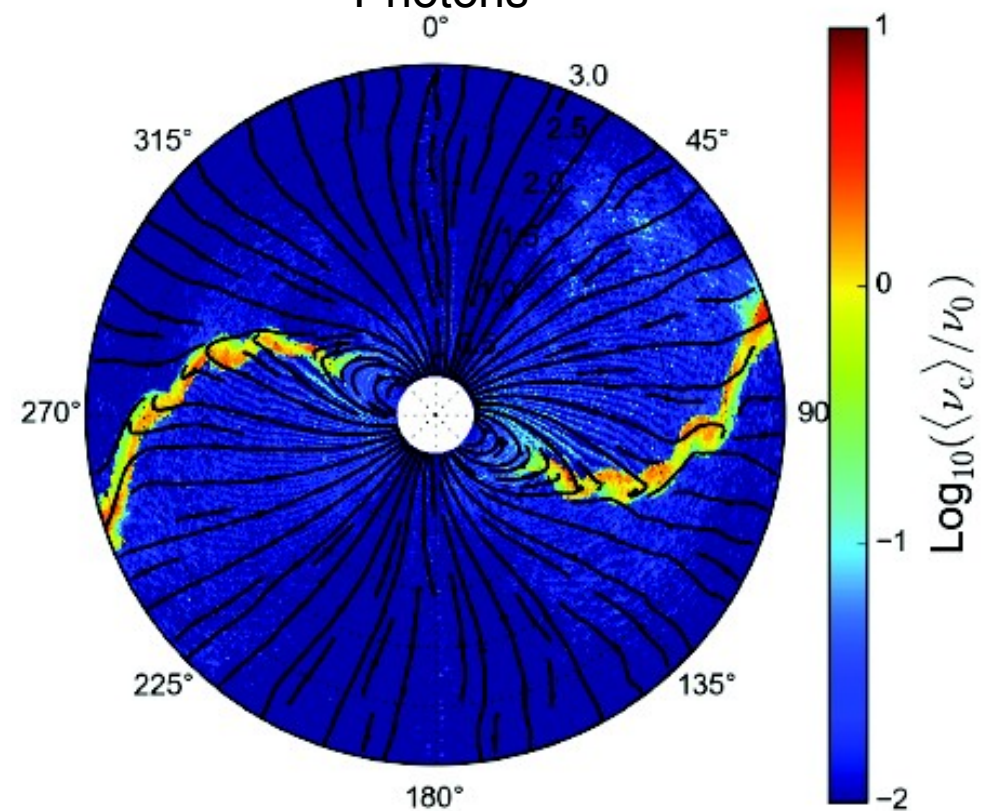
3D PIC : Particle / radiation mean energy ($\chi=30^\circ$)

Cerutti et al. 2016

Positrons



Photons



Particle acceleration via relativistic **reconnection** in the **current sheet**
High-energy radiation is **synchrotron radiation**

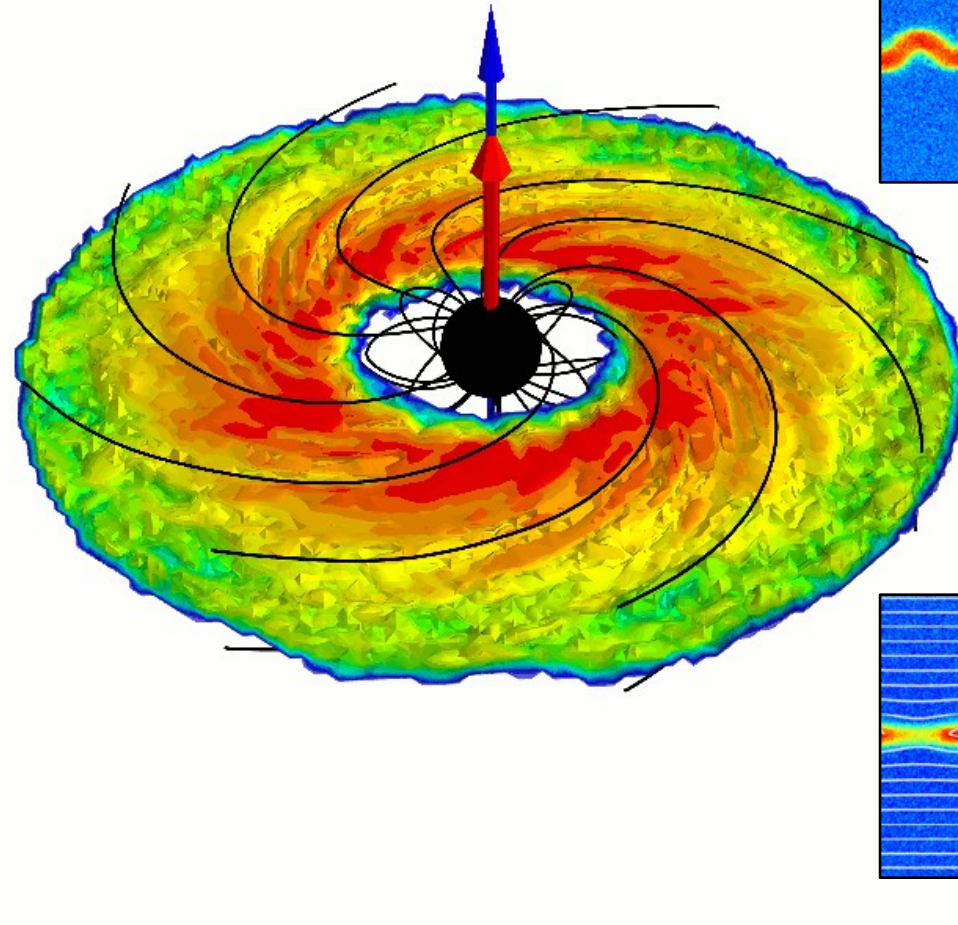
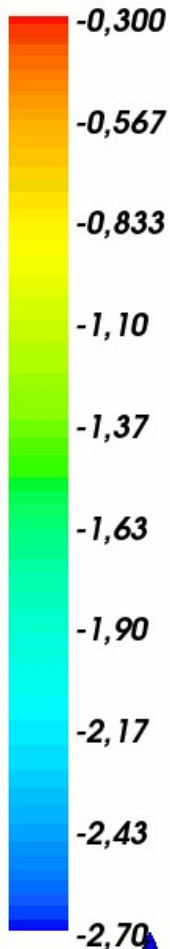
Particle energy in the sheet given by :

$$\sigma_{LC} = \frac{B_{LC}^2}{4\pi\Gamma n_{LC} m_e c^2} \approx 50 \quad (\text{here})$$

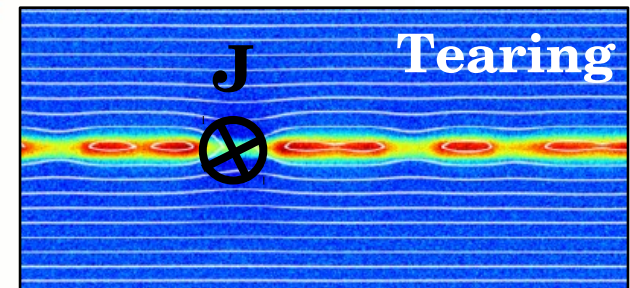
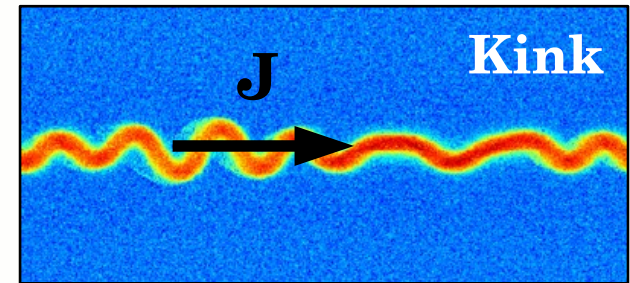
High-energy radiation flux ($\nu > \nu_0$, $\chi = 0^\circ$)

$i=0$ - Phase=0.00 - Positrons -

Log(Flux)



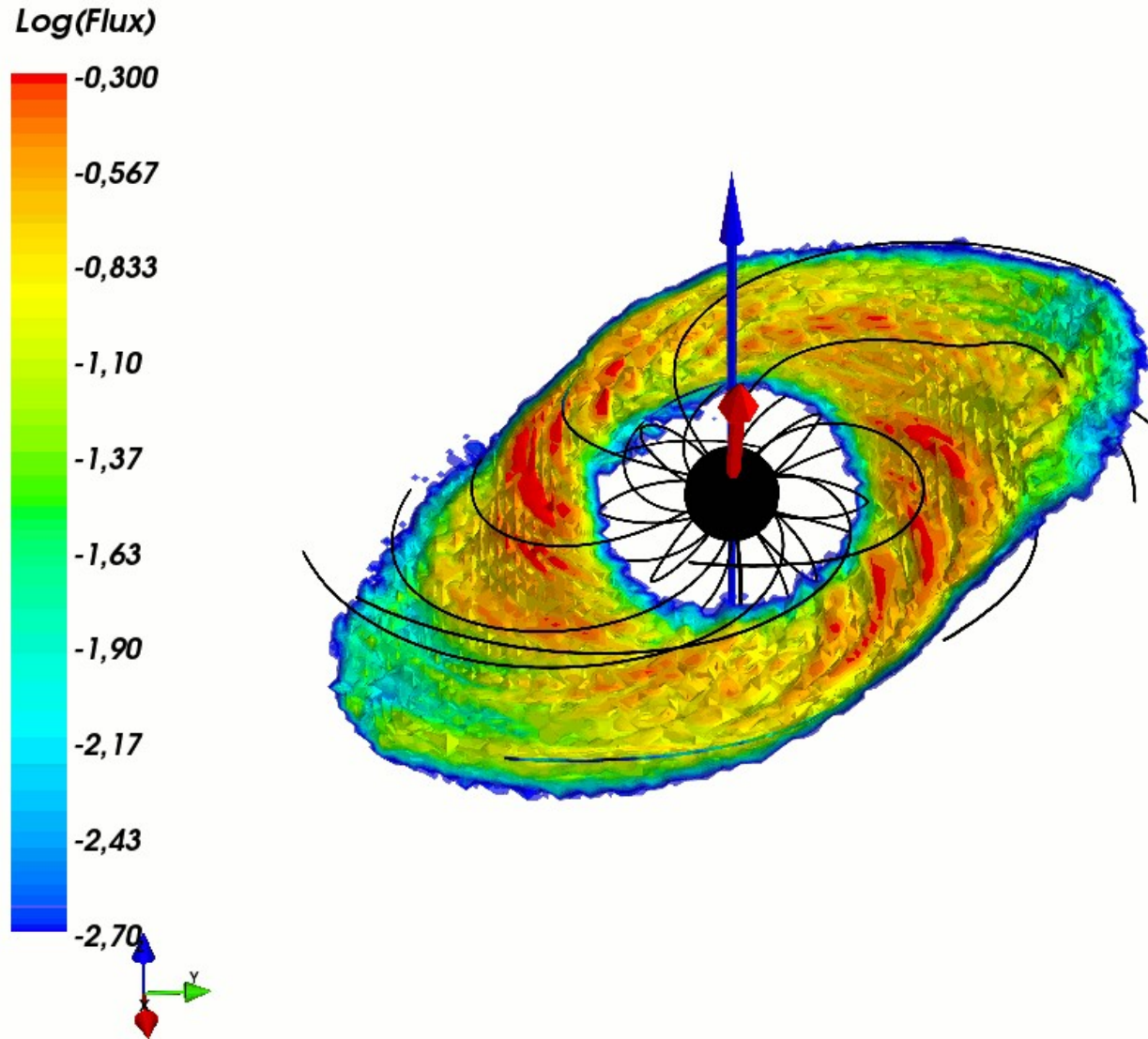
(from Local reconnection simulations)



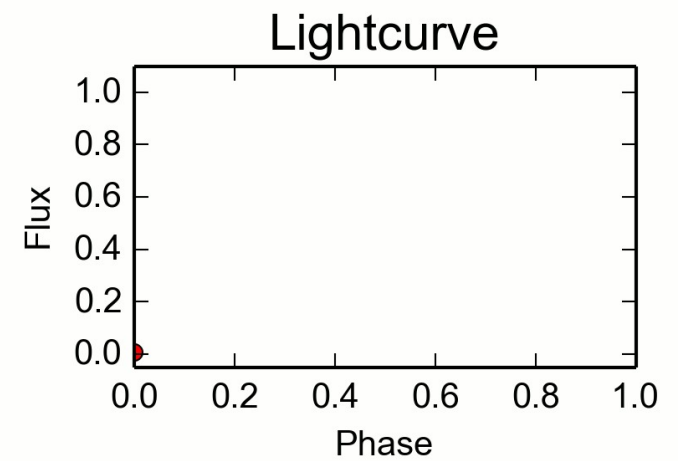
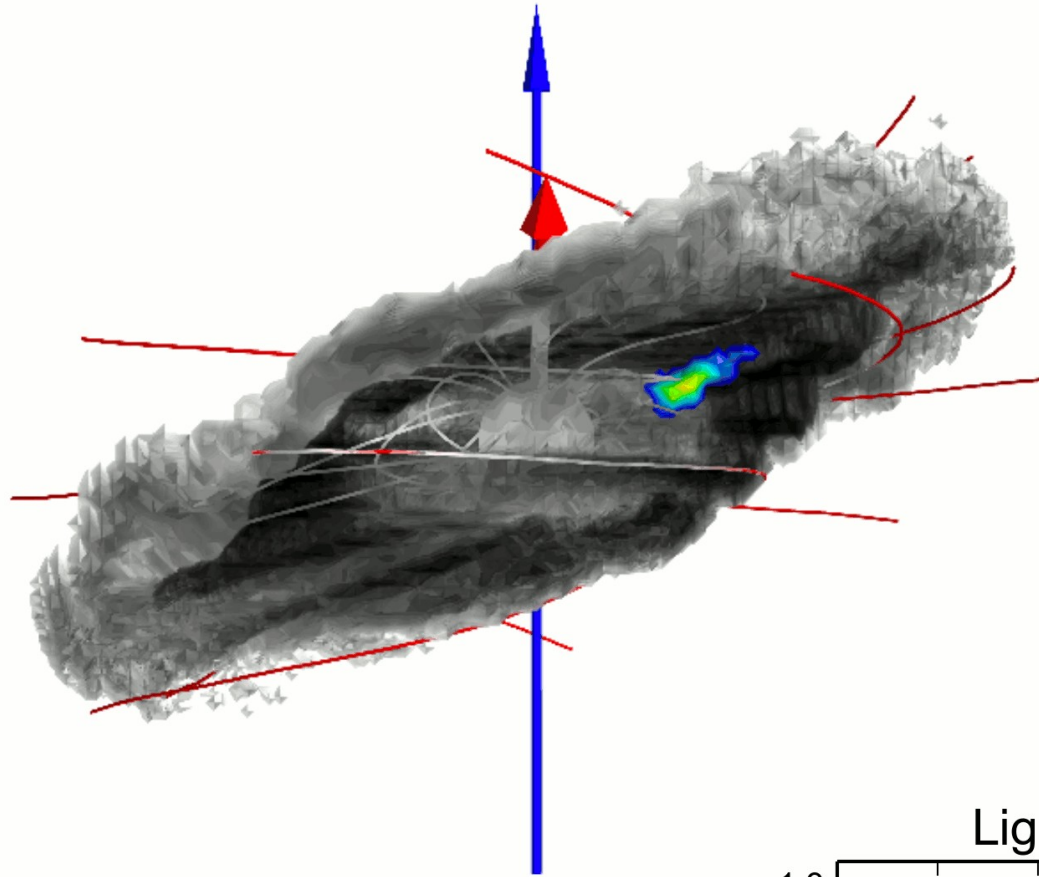
Presence of spatial irregularities due to **kinetic instabilities** in the sheet
e.g., **kink** and **tearing** modes

High-energy radiation flux ($\nu > \nu_0$, $\chi = 30^\circ$)

$i = 30$ - Phase = 0.00 - Positrons -

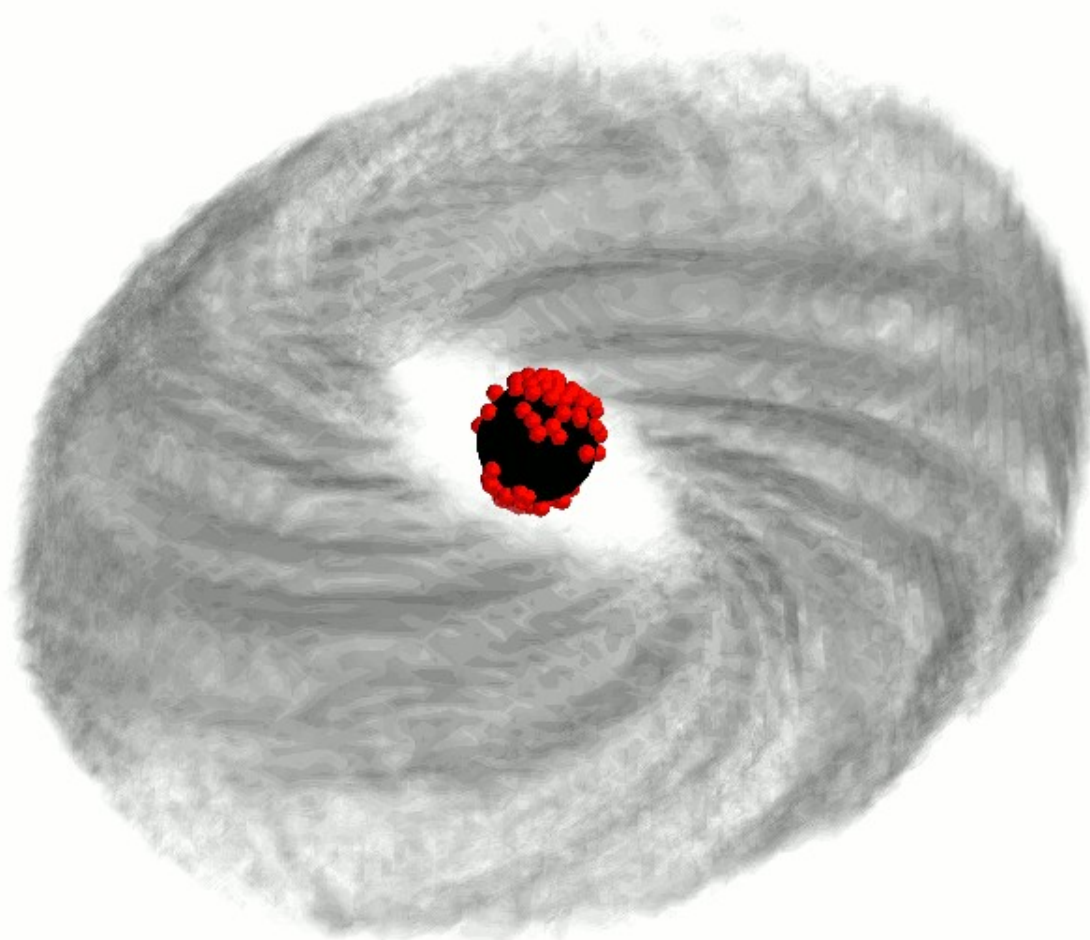
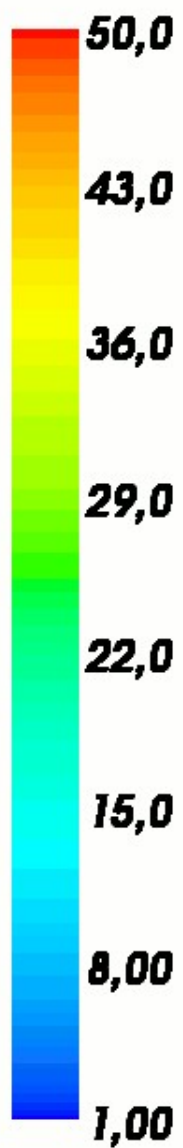


$i=30$ - Phase=0.00 - Positrons -



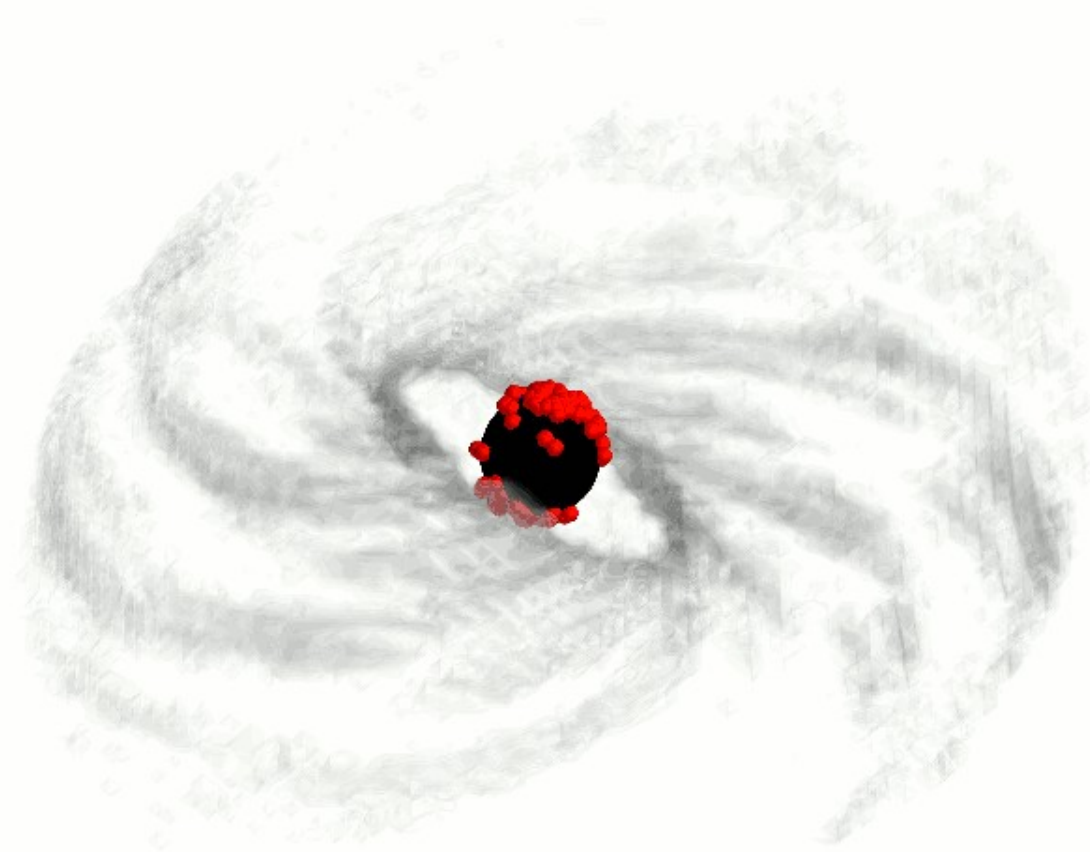
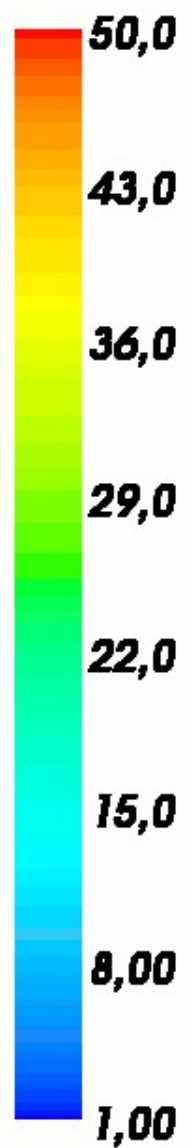
Tracked positrons

gamma



Tracked electrons

gamma



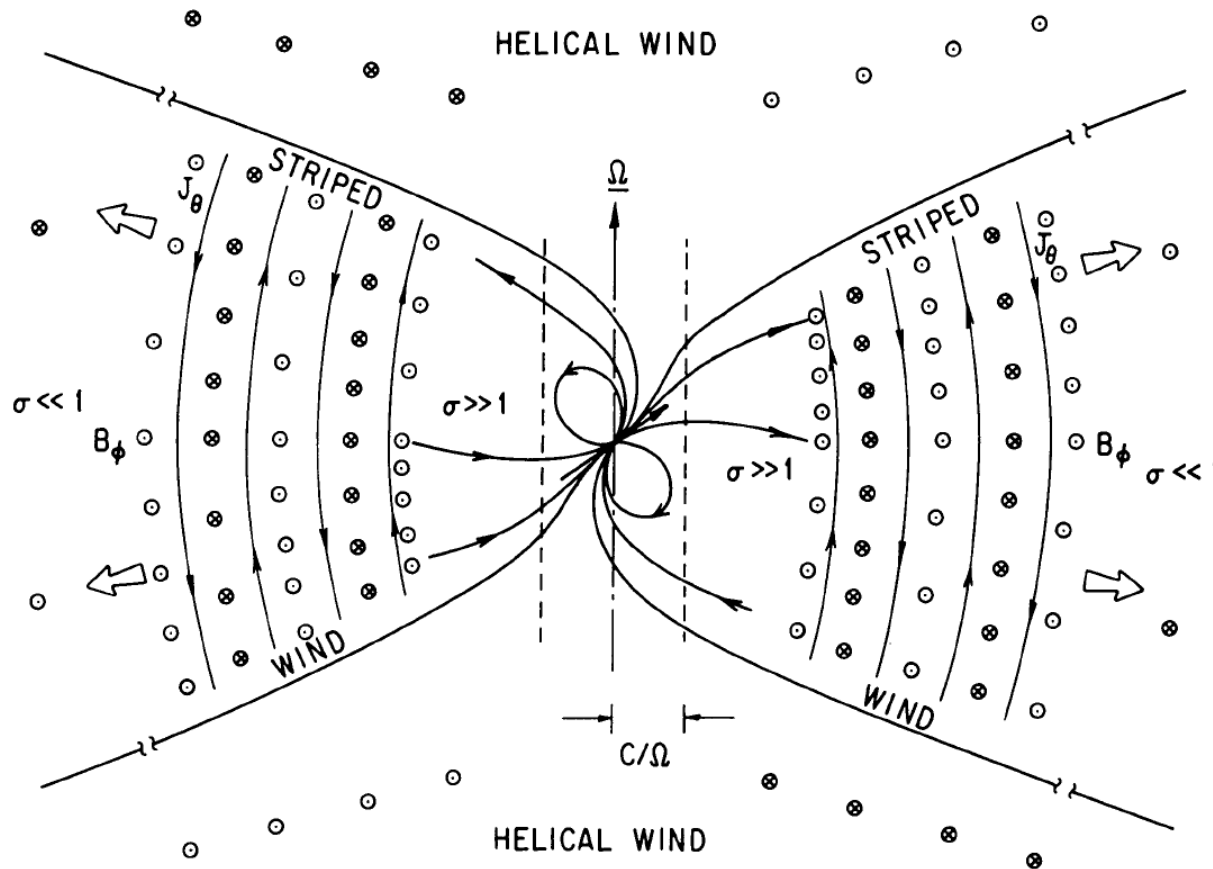
Dissipation in the wind

Close to light cylinder: ultra-magnetized plasma ($\sigma \gg 1$)

Nebula: Particle dominated ($\sigma \ll 1$)

=> Dissipation somewhere in between needed!

Coroniti 1990 ; Michel 1994 ; Lyubarsky & Kirk 2001 ; Kirk & Skjæraasen 2003 ; Zrake & Arons 2016



How far does magnetic reconnection proceed in the wind?

Shock-driven reconnection

If dissipation does not happen before the shock :

=> Shock-driven reconnection at the termination shock

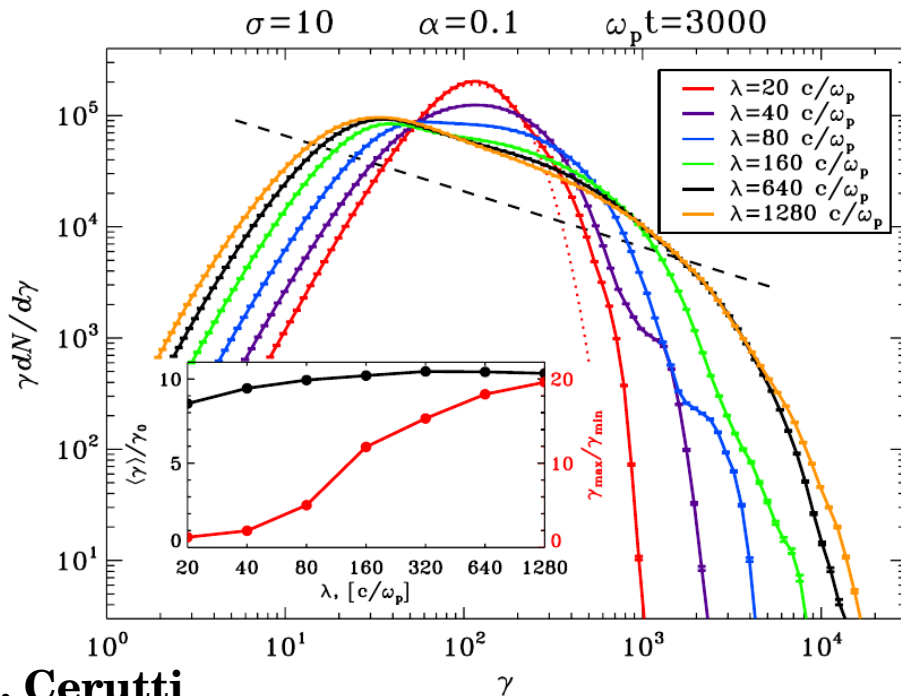
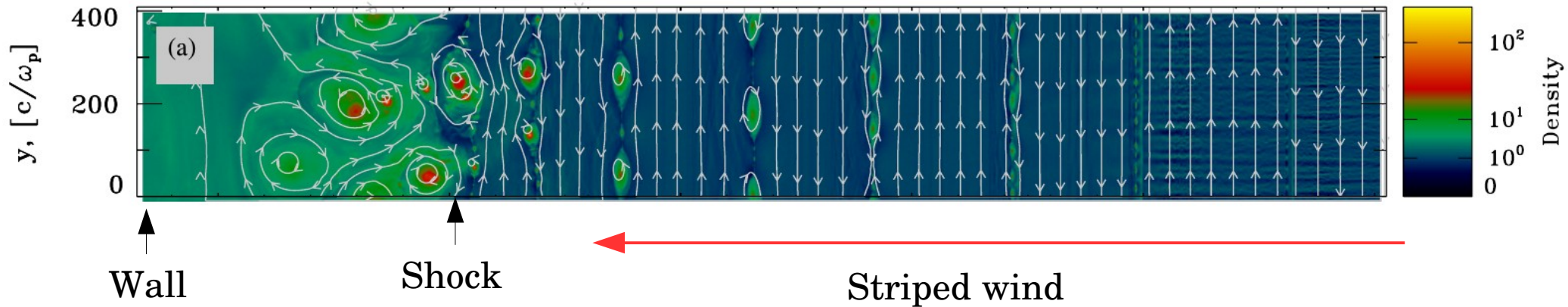
[Lyubarsky 2003 ; Pétri & Lyubarsky 2007 ; Sironi & Spitkovsky 2011]

$\sigma=10$

$\lambda=640 \text{ c}/\omega_p$

$\alpha=0.1$

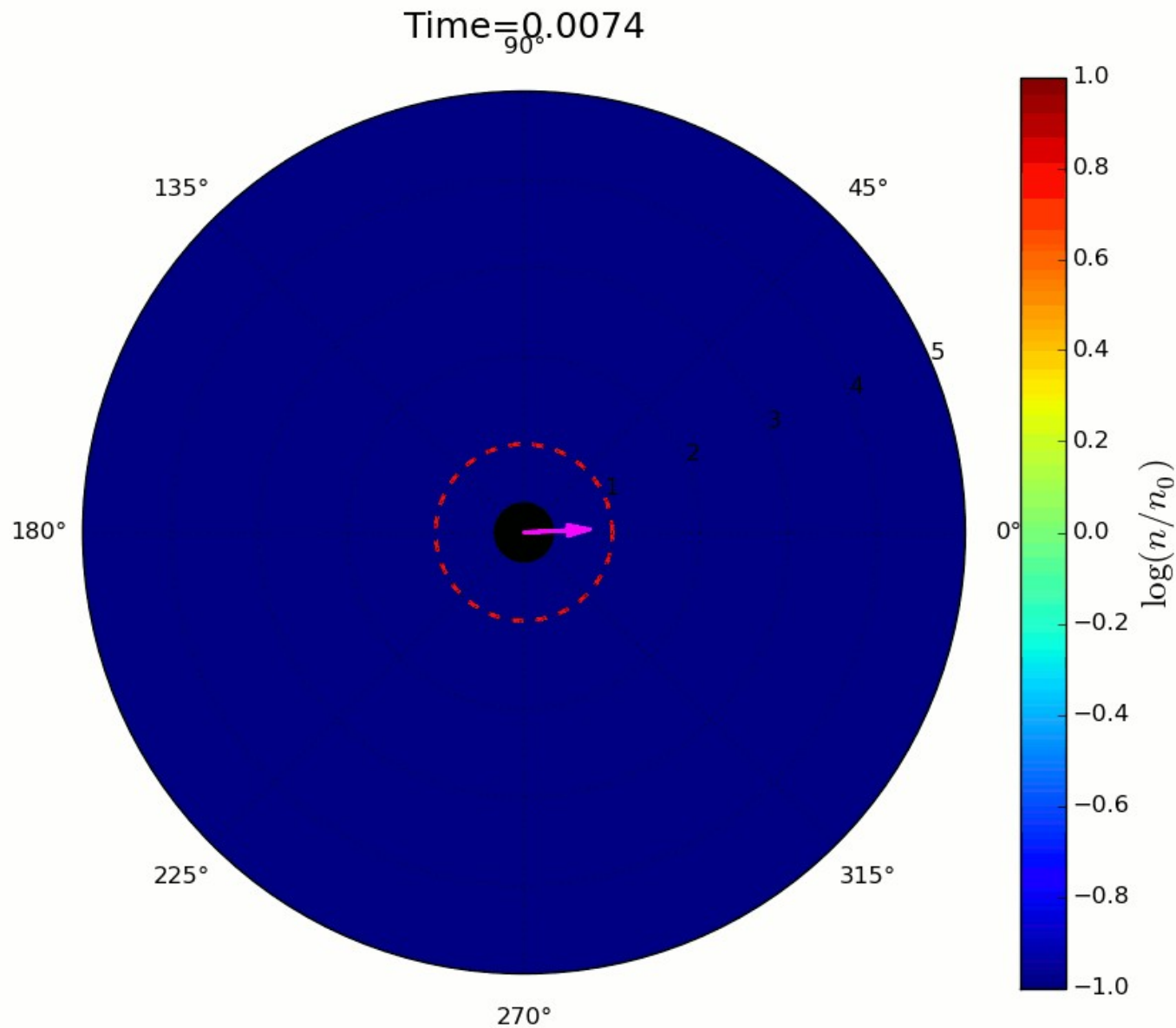
$\omega_p t=3750$



Non-thermal distribution if

$$\frac{\lambda}{r_L \sigma} \gg 1$$

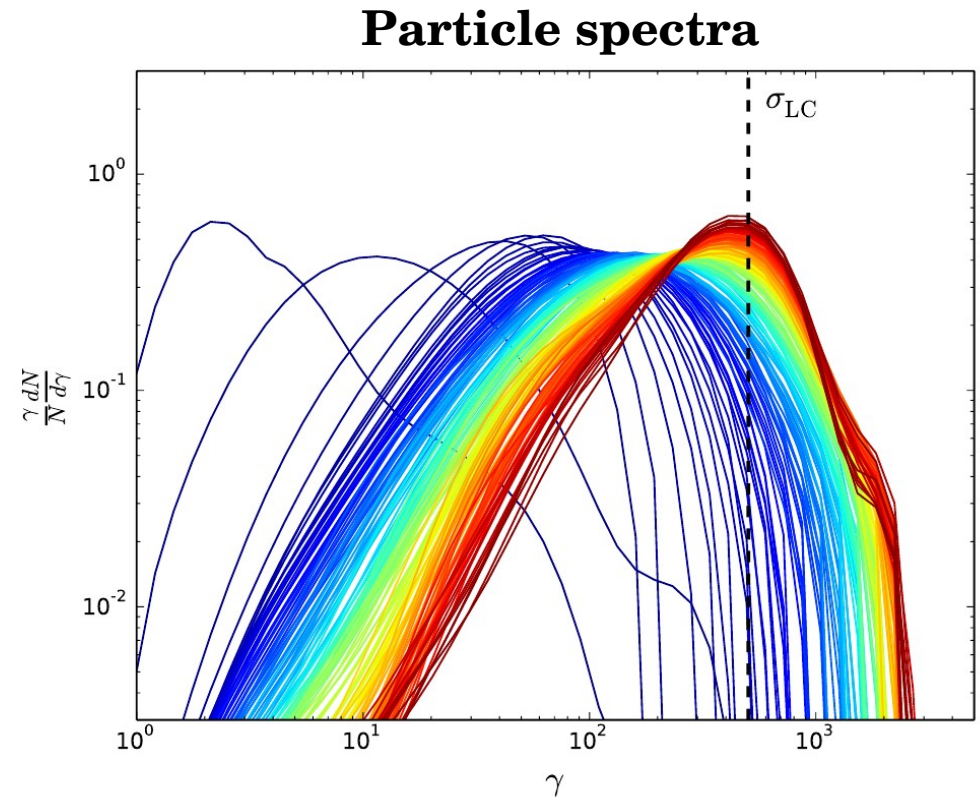
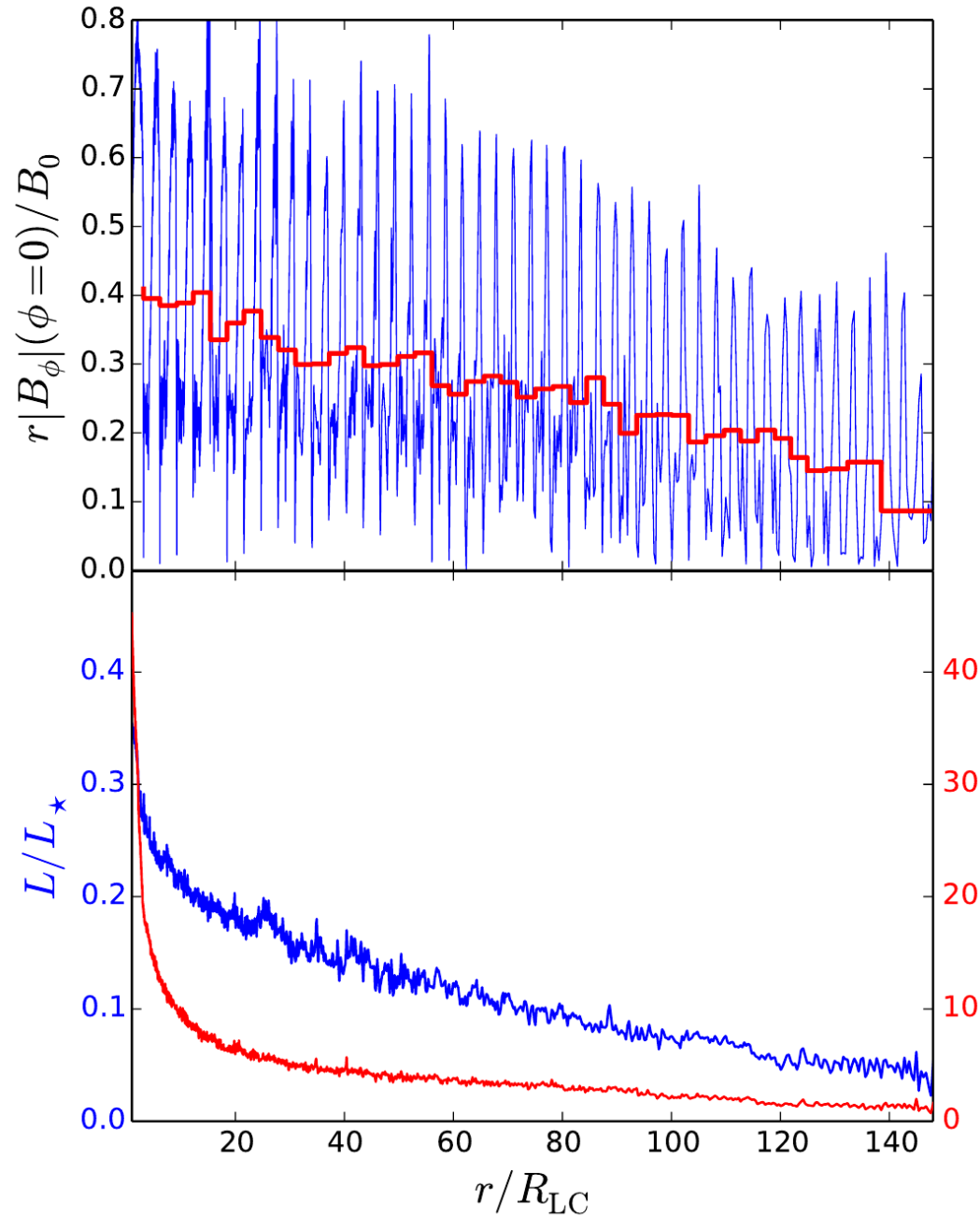
Otherwise : **Maxwellian** distribution



Small radii : *Turbulent-like* reconnection : Chain of plasmoids, mergers

Large radii ($r \gg RLC$) : *Smooth, laminar dissipation*, islands frozen by adiabatic expansion

Dissipation in the wind



Complete annihilation of the striped wind after a couple **$100 R_{\text{LC}}$** .

Narrow particle energy distribution set by **σ_{LC}** .

Cerutti & Philippov in prep

Conclusions

- Relativistic reconnection is **fast**, and **efficient at accelerating particles** !
Viable **alternative to shock** acceleration in **magnetized environments**
- **Hard** particle energy spectra are expected.
- **Several acceleration sites identified, mostly X-points and mergers.**
- Maximum particle energy in the sheet not limited by **synchrotron losses**
=> Emission **>100 MeV synchrotron** radiation possible
- Crab gamma-ray flare most convincing case for relativistic reconnection.
- Application to **blazars flares**, **pulsar** magnetospheres, magnetized wind...

...and most likely more applications to come, stay tuned!